

Congratulations!

John

On Fifty Remarkable Years of
Research at Argonne





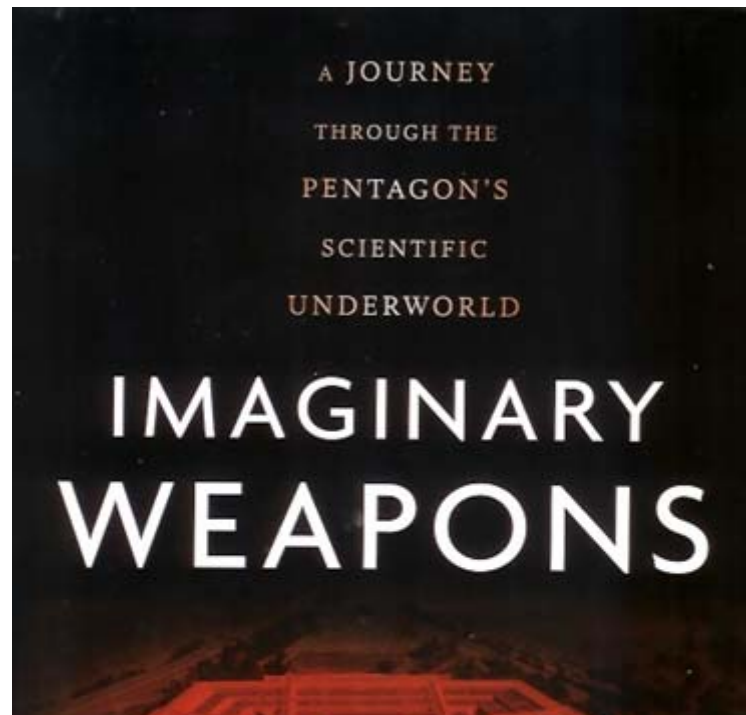
me

Me again
John P. Schiffer

art

BUT WAIT,
THESE
ARE THE
BEST TO FIND

CAN STRAIGHT



their students. Among colleagues, Schiffer had a reputation as a world-class scientist, albeit a tad blunt, with a habit of calling people fools to their face. But Schiffer was not the

Me

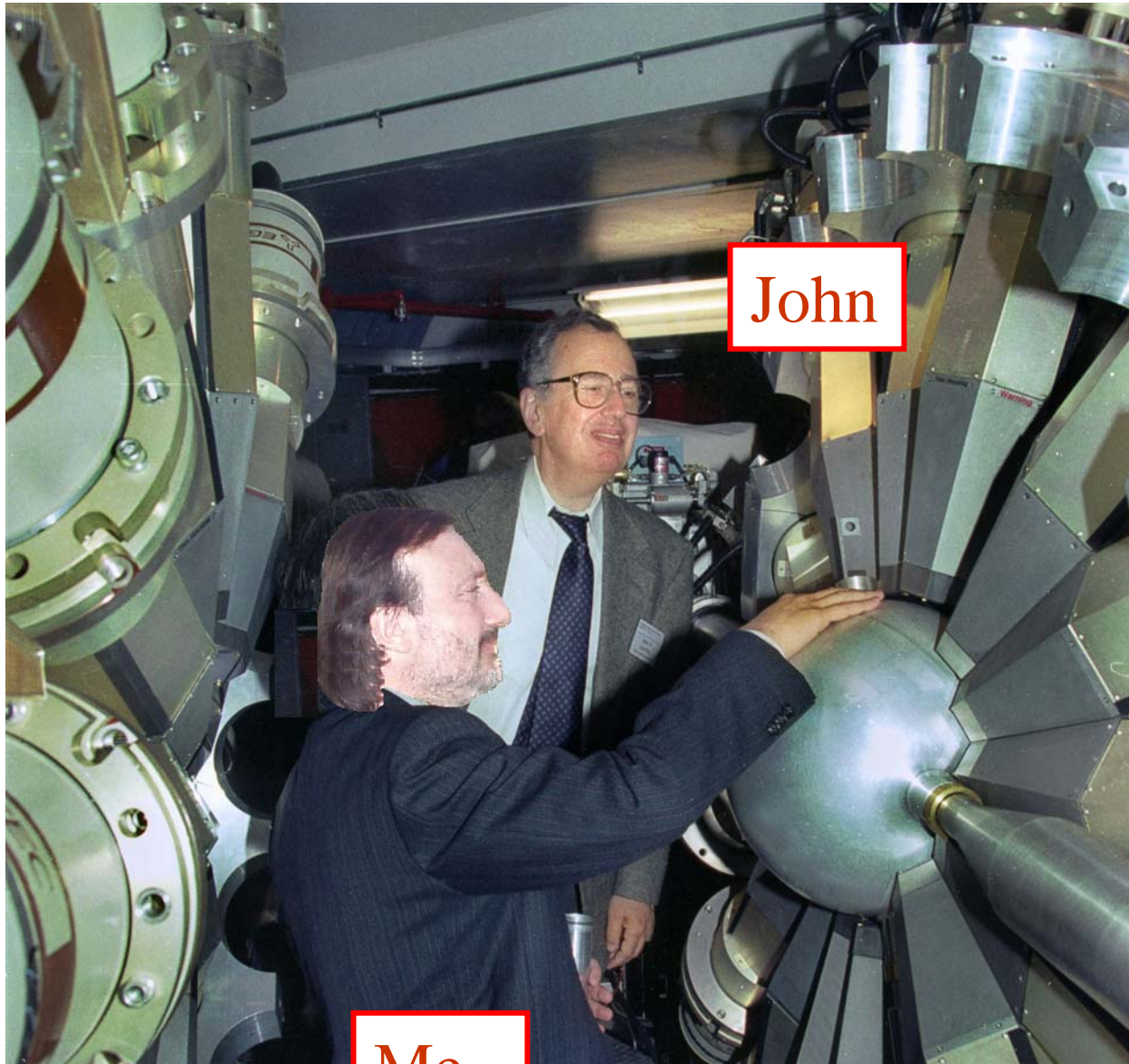
Me

John



GORDON RESEARCH CONFERENCE
Colby Sawyer College (1)
NUCLEAR CHEMISTRY

Checking a setup in GAMMASPHERE !

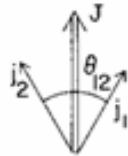


John

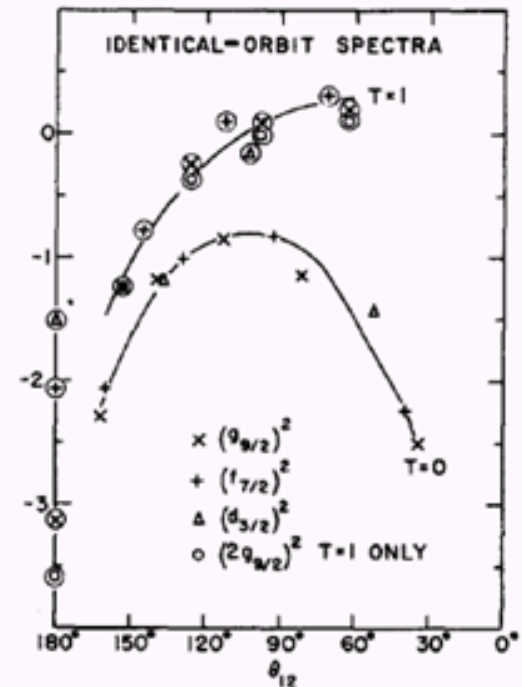
Me Lee

The effective interaction between nucleons deduced from nuclear spectra

John P. Schiffer, William W. True



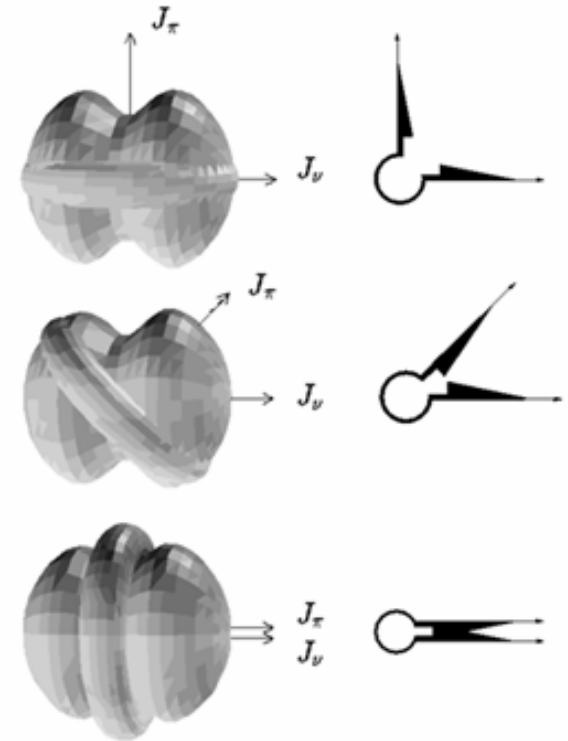
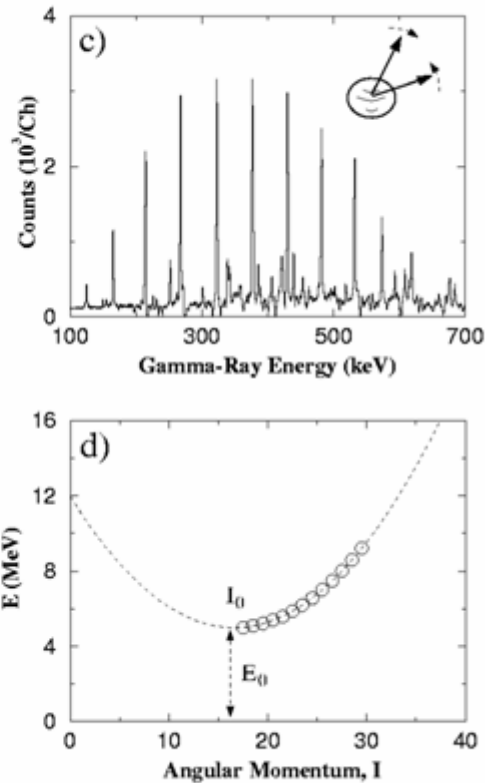
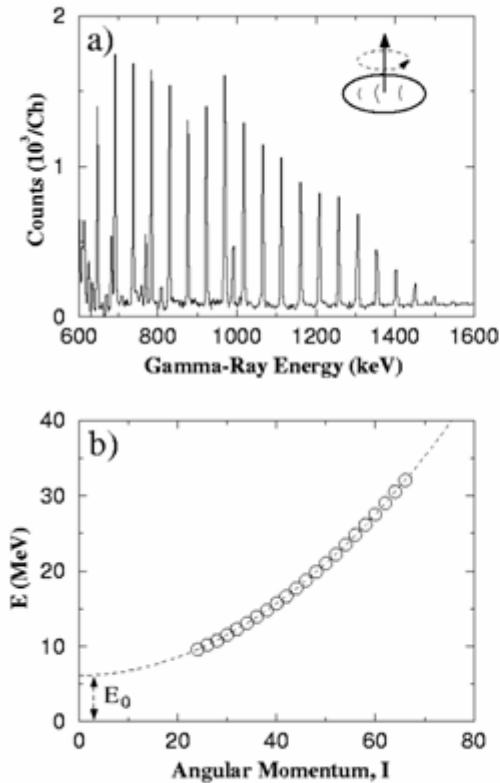
$$\theta_{12} = \cos^{-1} \frac{j_1(j_1+1) + j_2(j_2+1) - J(J+1)}{[2j_1j_2(j_1+1)(j_2+1)]^{1/2}}$$



Rotational-like (Shears) bands in near-spherical nuclei

SD Band in ^{152}Dy

M1 Band in ^{199}Pb

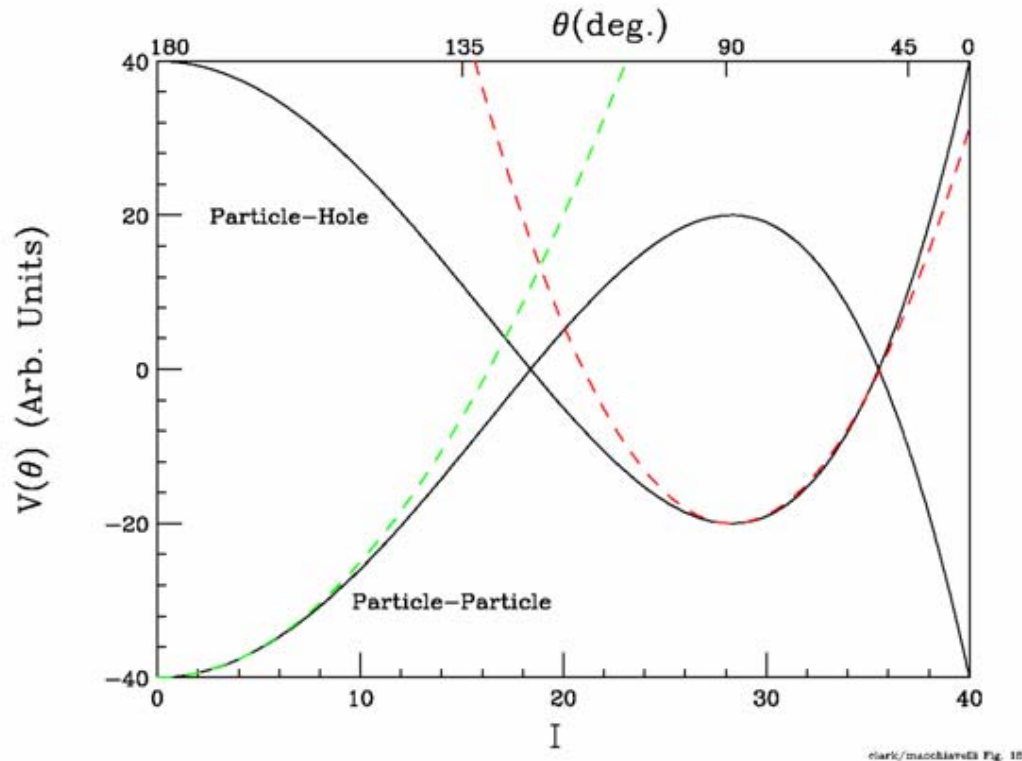


darkmacchiwell Fig. 1

Angular momentum generated by the recoupling of the protons and neutrons components (shears)

$B(M1)$, proportional to $\mu_\perp^2 \Rightarrow$ should decrease with I

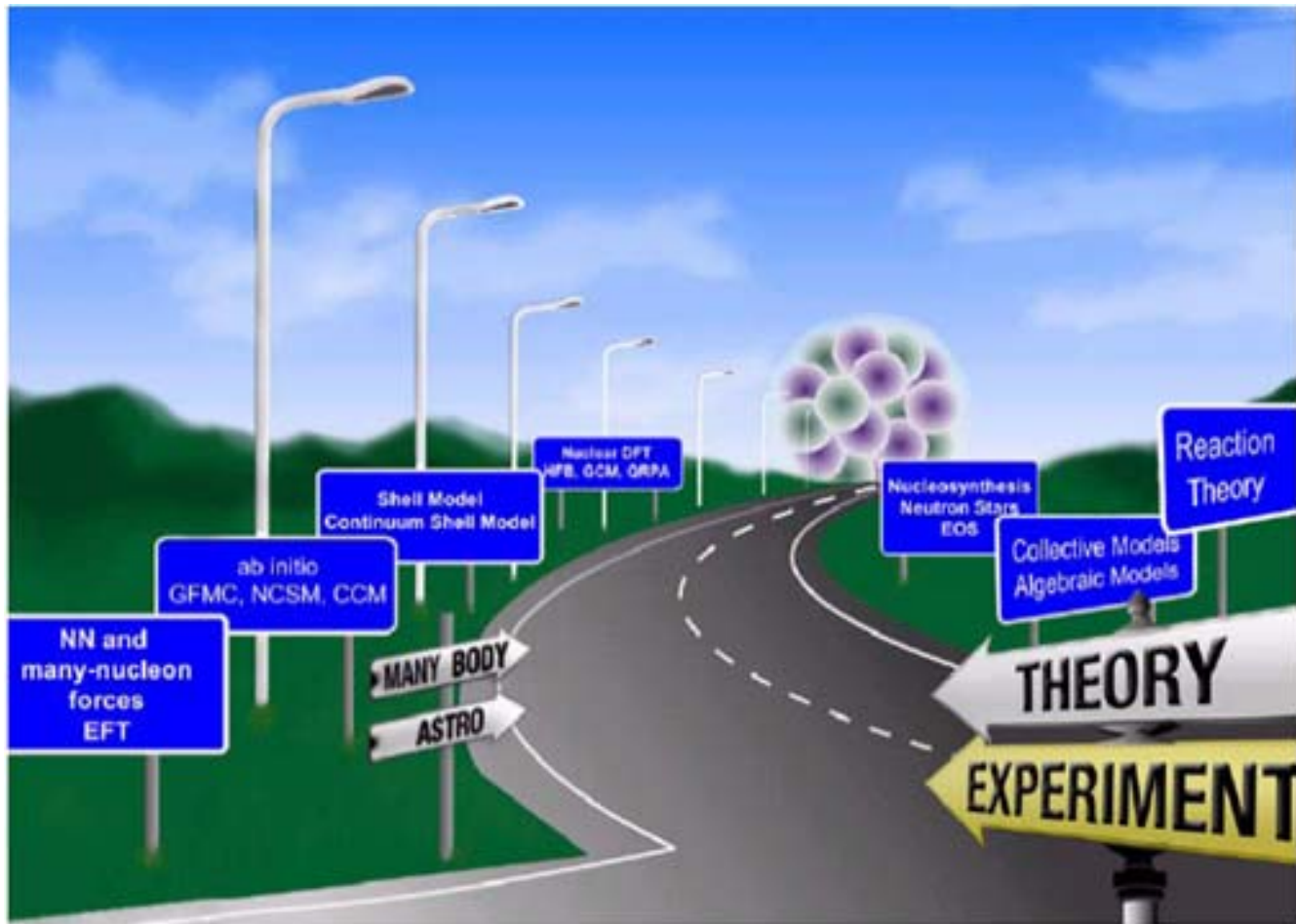
Effective P_2 Interaction ?



Effective P_2 interaction can give rise to the rotational-like spectrum
Particle-particle case will correspond to Anti-magnetic rotation

\Rightarrow Although the main ingredients and the general properties of the shears bands are understood, the regularity of these bands is still an important theoretical question from a microscopic point of view.

Inspired by the RIA Theory Blue Book and the recent RIA brochure ..



Transfer Reactions and Nuclear Structure: The Road Ahead

A.O. Macchiavelli
Lawrence Berkeley National Laboratory

Many Thanks to Rod Clark, Paul Fallon, I. Yang Lee,
David Radford and Alan Wuosmaa

Nuclear Physics, The Core of Matter, the Fuel of Stars
Argonne National Laboratory, IL, September 21-22, 2006

Outline

- Short introduction
- One-nucleon transfer reactions
 - Single-particle degrees of freedom
 - Spin-orbit splitting
 - Ab initio calculations
- Two-nucleon transfer reactions
 - Correlations
 - NP pairing
 - Pairing phase transition
- Heavy-ion transfer reactions
- Summary and conclusions

A central theme of study in nuclear structure has been the understanding of the elementary modes of excitation of the atomic nucleus, and their evolution with A (size), I (rotational frequency), $(N-Z)$ (Isospin) and E^* (Temperature)

Transfer reactions have played a major role in this endeavor, particularly in the characterization of the single particle degrees of freedom and their correlations.

With the development of exotic beams and new instruments to handle the “reverse kinematics” nature of these reactions, there has been a *renaissance* of transfer reaction studies.

On Angular Distributions from (d, p) and (d, n)
Nuclear Reactions

S. T. BUTLER*

Department of Mathematical Physics, University of Birmingham,
Birmingham, England

October 30, 1950

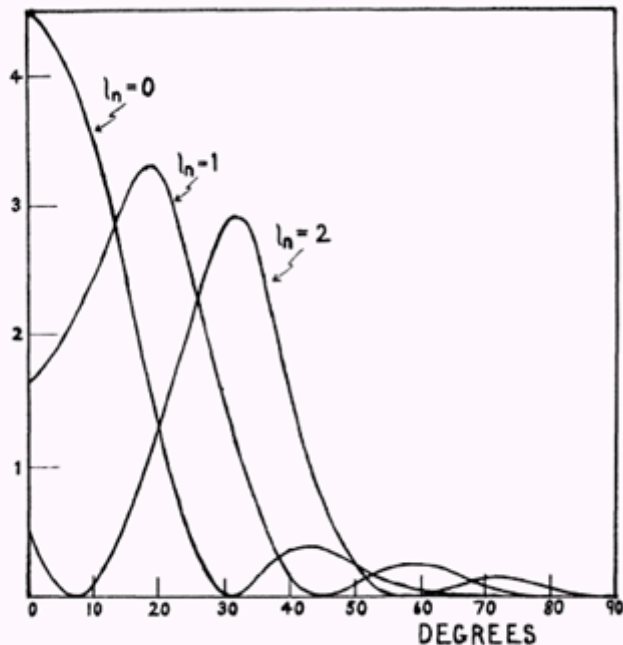


FIG. 1. Theoretical angular distributions for (d, p) and (d, n) reactions for different angular momentum transfers to the initial nucleus.

Angular distribution of the outgoing particles reflect the transferred angular momentum \Rightarrow l-value

Spectroscopic factors. Overlap between initial and final state \Rightarrow Test wave functions

Ideal tool to study the single-particle degree of freedom. Complementary to other processes that probe collective aspects.

One of the important goals in the development of exotic beams is to study the evolution of the nuclear shell structure with neutron excess.

In fact, it is already known that in light nuclei the magic numbers are not as “robust” as originally conceived.

Of particular interest here is the spin-orbit force.

Is the Nuclear Spin-Orbit Interaction Changing with Neutron Excess?

J. P. Schiffer,¹ S. J. Freeman,^{1,2} J. A. Caggiano,³ C. Deibel,³ A. Heinz,³ C.-L. Jiang,¹ R. Lewis,³ A. Parikh,³ P. D. Parker,³
K. E. Rehm,¹ S. Sinha,¹ and J. S. Thomas⁴

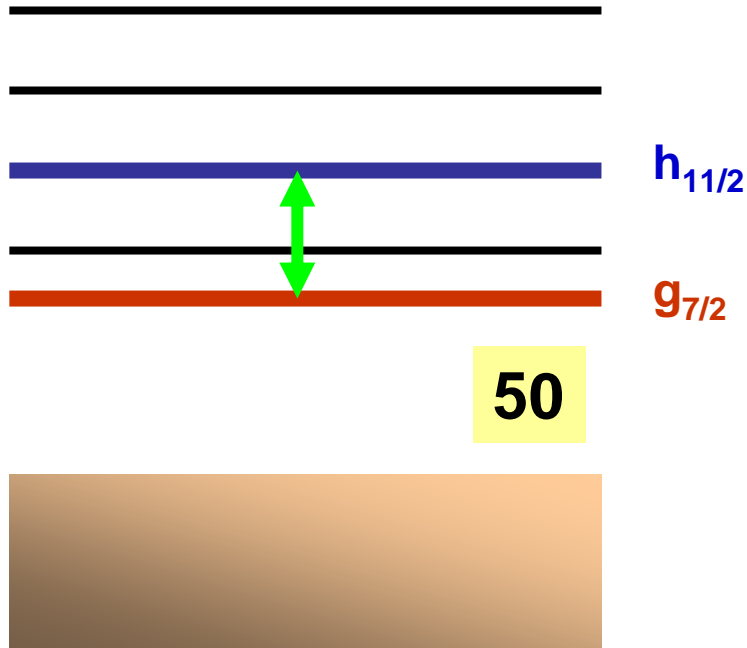
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(Received 17 December 2003; published 20 April 2004)



$$\delta E \approx \hbar \omega_0 - 4.5V_{ls} / A^{2/3}$$

Use (α,t) reaction to populate high- l orbitals

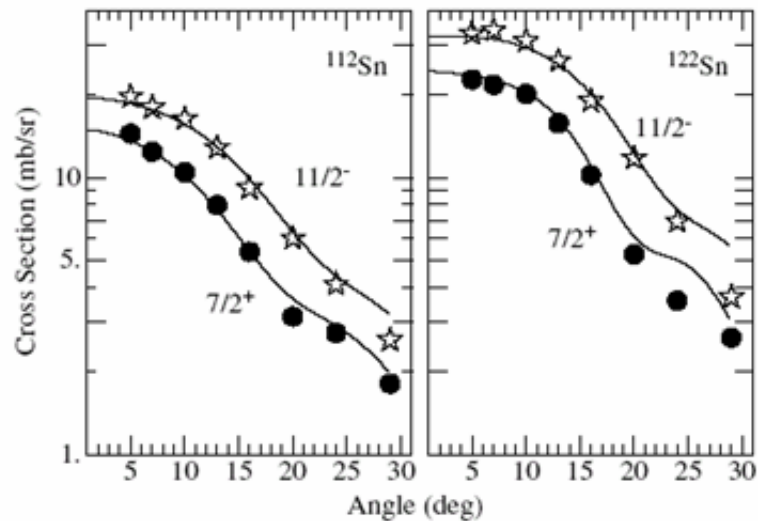
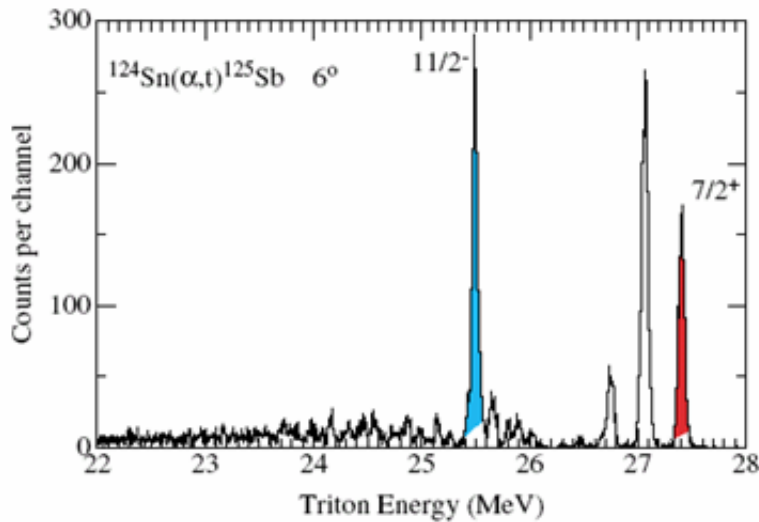
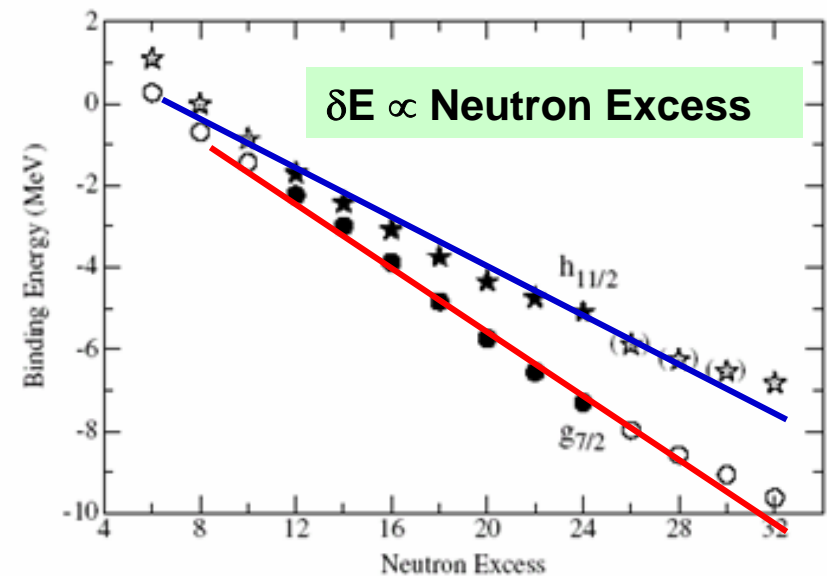


TABLE I. Cross sections (mb/sr) at 6° for the lowest $7/2^+$ and $11/2^-$ states, their ratios, and spectroscopic factors. The uncertainties in the cross sections are estimated at 10% and those in the ratio, at about 5%. The accuracy of the relative spectroscopic factors are estimated at 5%.

| Target | $7/2^+$ | $11/2^-$ | Ratio | $C^2S_{7/2}$ | $C^2S_{11/2}$ |
|-------------------|---------|----------|-------|--------------|---------------|
| ^{112}Sn | 14.6 | 21.4 | 1.47 | 0.99 | 0.84 |
| ^{114}Sn | 19.6 | 27.3 | 1.39 | 1.10 | 0.93 |
| ^{116}Sn | 19.7 | 30.9 | 1.57 | 0.95 | 0.97 |
| ^{118}Sn | 20.4 | 33.5 | 1.64 | 0.88 | 0.99 |
| ^{120}Sn | 27.9 | 39.4 | 1.41 | 1.13 | 1.12 |
| ^{122}Sn | 24.6 | 35.5 | 1.45 | 0.98 | 1.00 |
| ^{124}Sn | 24.7 | 39.2 | 1.59 | 1.00 | 1.12 |



QUENCHING OF THE $2p_{1/2}-2p_{3/2}$ PROTON SPIN-ORBIT SPLITTING IN THE Sr-Zr REGION

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Department of Physics and Astronomy, San Francisco State University, San Francisco, CA 94132, USA

S. PITTEL ²

Barol Research Foundation of the Franklin Institute, University of Delaware, Newark, DE 19716, USA

and

A. ETCHEGOYEN

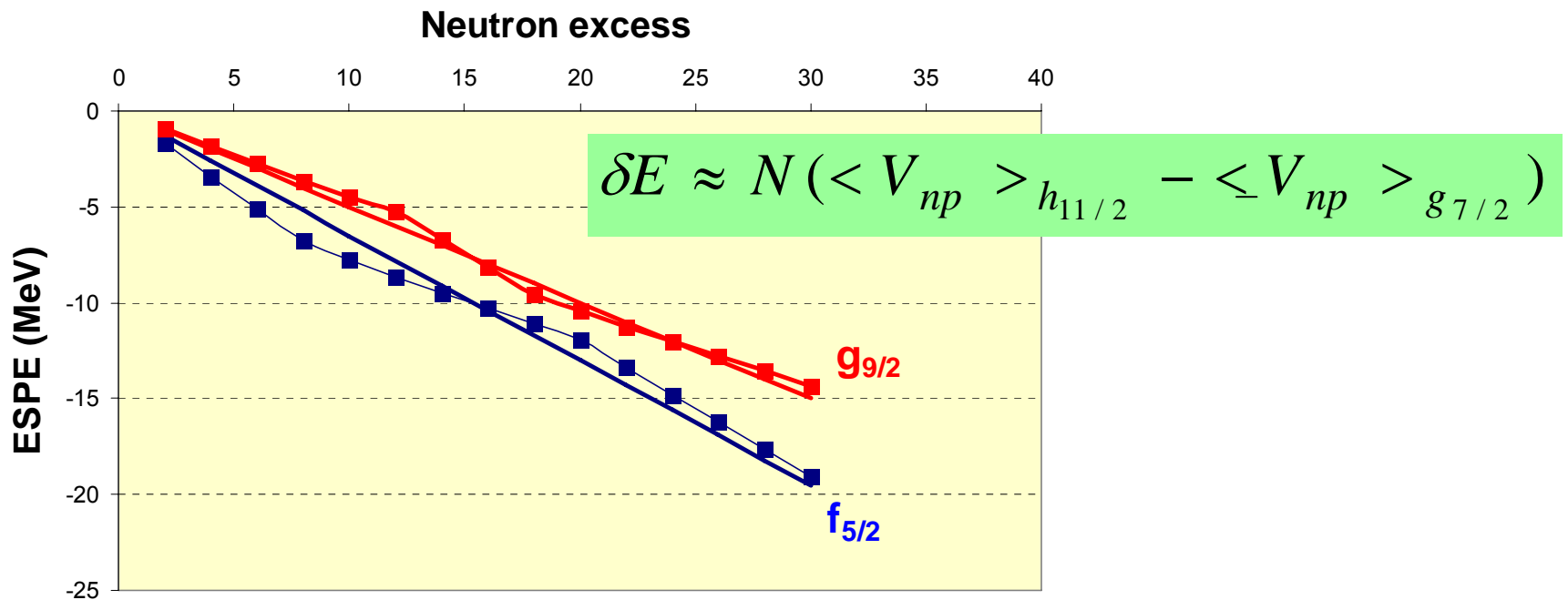
Departamento de Física, Comisión Nacional de Energía Atómica, Av. del Libertador 8250, 1429 Buenos Aires, Argentina

⁹⁵Y

$$\Delta\epsilon(N = 56) - \Delta\epsilon(N = 50)$$

⁸⁹Y

$$= 6 \left(\overline{V}_{p_{1/2}-d_{5/2}}^{n-p} - \overline{V}_{p_{3/2}-p_{5/2}}^{n-p} \right),$$



Reduction of the Spin-Orbit Splittings at the $N = 28$ Shell Closure

L. Gaudefroy,^{1,2} O. Sorlin,^{2,1} D. Beaumel,¹ Y. Blumenfeld,¹ Z. Dombrádi,³ S. Fortier,¹ S. Franchoo,¹ M. Gélín,² J. Gibelin,¹ S. Grévy,² F. Hammache,¹ F. Ibrahim,¹ K. W. Kemper,⁴ K.-L. Kratz,^{5,6} S. M. Lukyanov,⁷ C. Monrozeau,¹ L. Nalpas,⁸ F. Nowacki,⁹ A. N. Ostrowski,^{5,6} T. Otsuka,¹⁰ Yu.-E. Penionzhkevich,⁷ J. Piekarewicz,⁴ E. C. Pollacco,⁸ P. Roussel-Chomaz,² E. Rich,¹ J. A. Scarpaci,¹ M. G. St. Laurent,² D. Sohler,¹¹ M. Stanoiu,¹² T. Suzuki,¹³ E. Tryggestad,¹ and D. Verney¹

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$^{46}\text{Ar}(d,p)^{47}\text{Ar}$ reaction at SPIRAL and MUST detector array

In recent years substantial progress has been made in the development of *Ab Initio* theories of light nuclei.

The Quantum Monte Carlo method has been very successful in reproducing BE and excitation spectra of light nuclei starting with the basic NN force and 3-body force.

Neutron Spectroscopic Factors in ${}^9\text{Li}$ from ${}^2\text{H}({}^8\text{Li}, p){}^9\text{Li}$

A. H. Wuosmaa,¹ K. E. Rehm,² J. P. Greene,² D. J. Henderson,² R. V. F. Janssens,² C. L. Jiang,² L. Jisonna,³ E. F. Moore,²
 R. C. Pardo,² M. Paul,⁴ D. Peterson,² Steven C. Pieper,² G. Savard,² J. P. Schiffer,² R. E. Segel,³
 S. Sinha,² X. Tang,² and R. B. Wiringa²

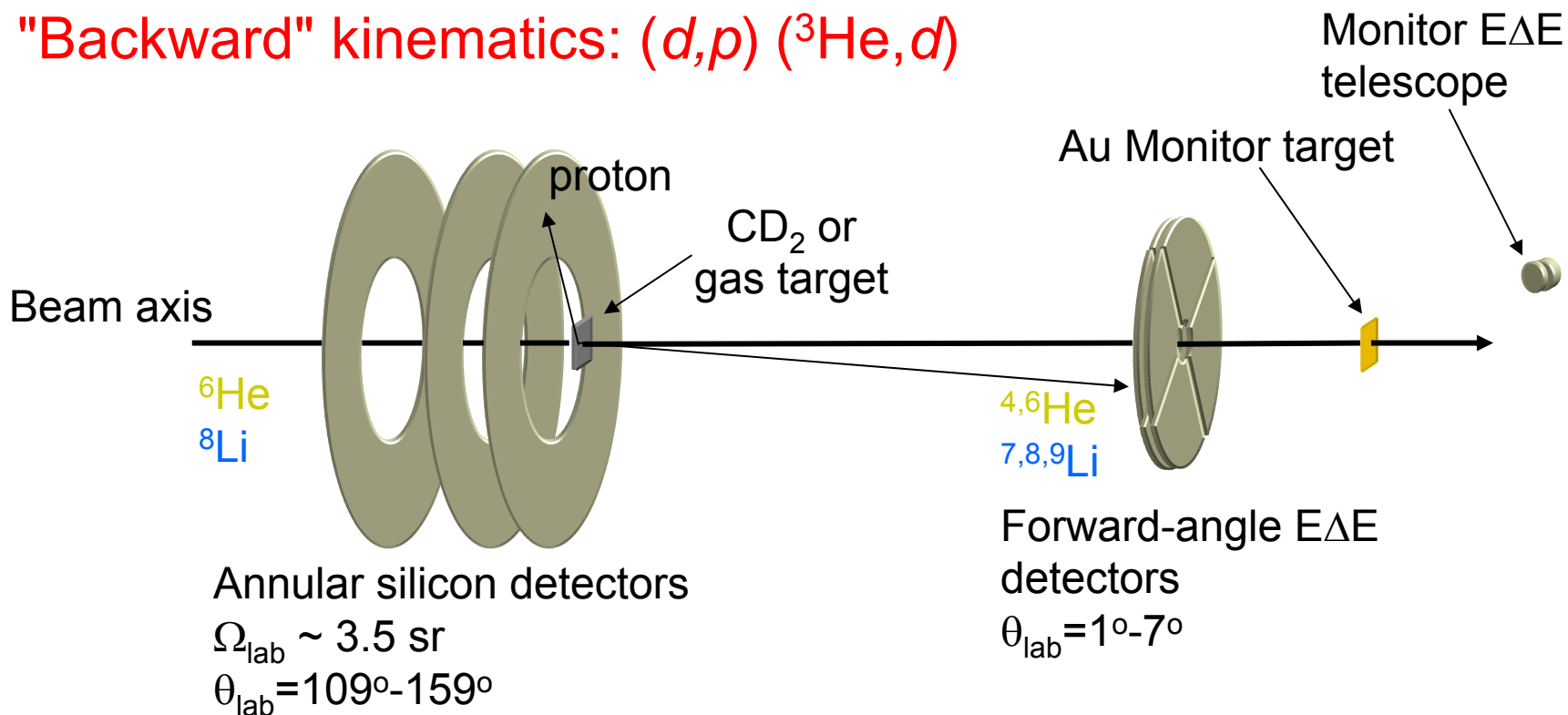
¹Physics Department, Western Michigan University, Kalamazoo, Michigan 49008-5252, USA

²Physics Division, Argonne National Laboratory, 9700 S. Cass Ave, Argonne Illinois 60439, USA

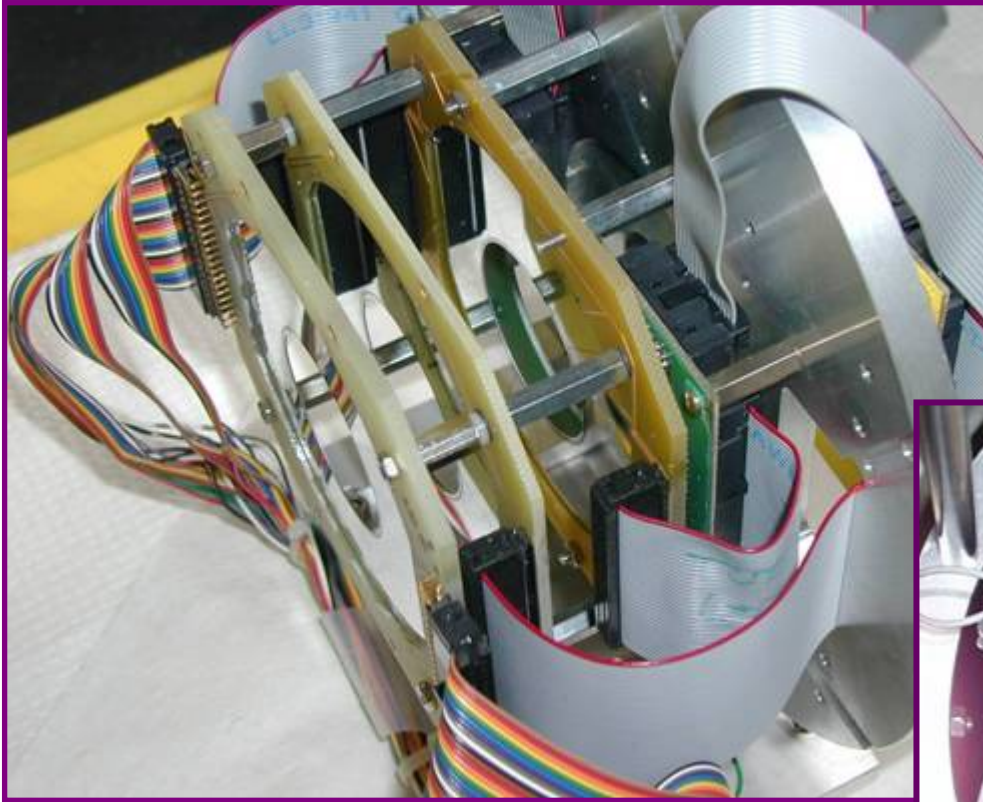
³Physics Department, Northwestern University, Evanston, Illinois 60208, USA

⁴Hebrew University, Jerusalem, Israel 91904

"Backward" kinematics: (d,p) (${}^3\text{He}, d$)

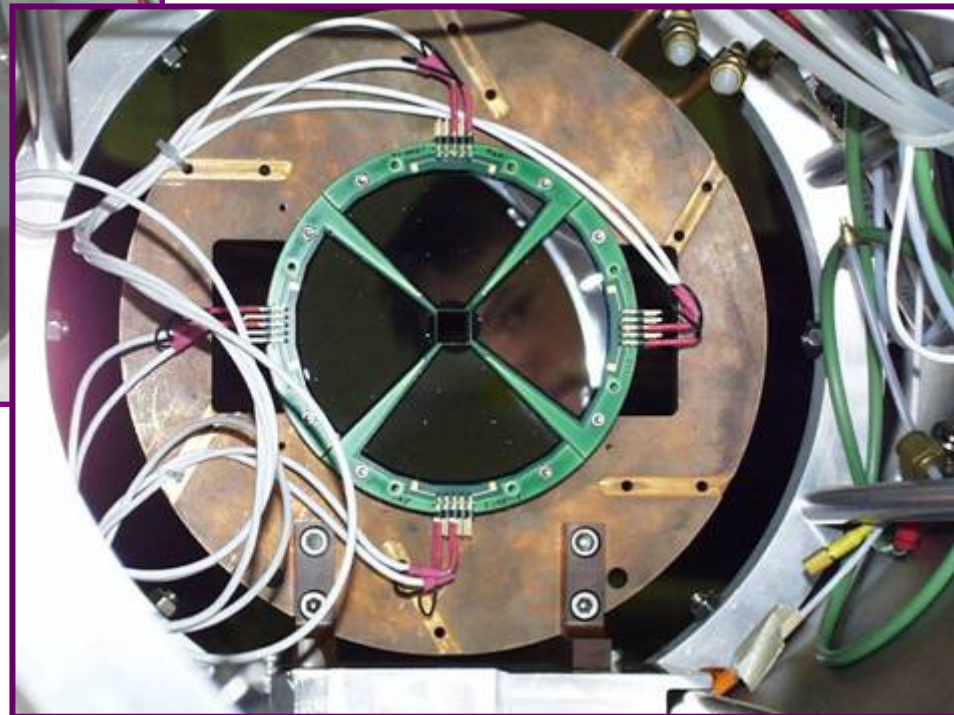
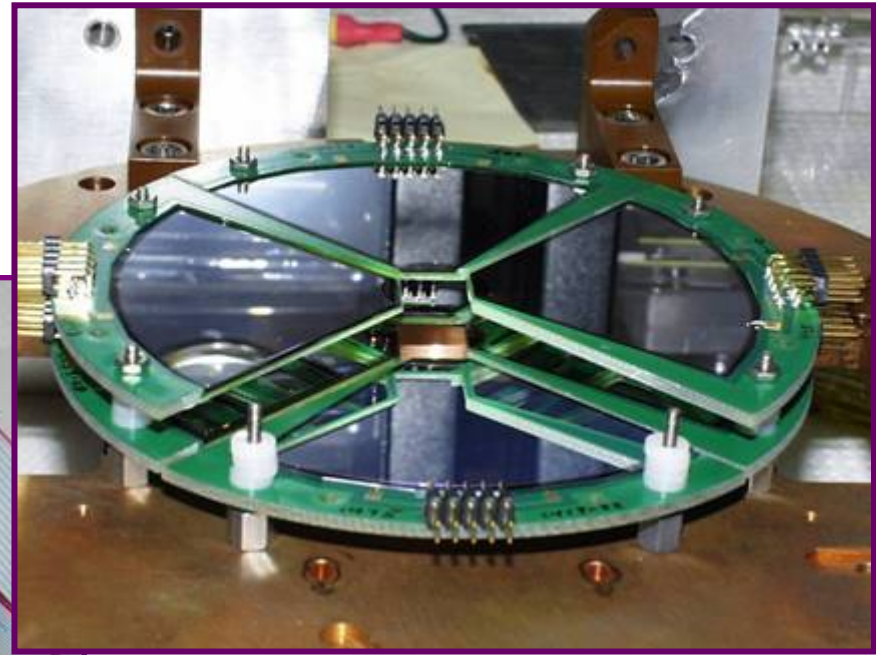


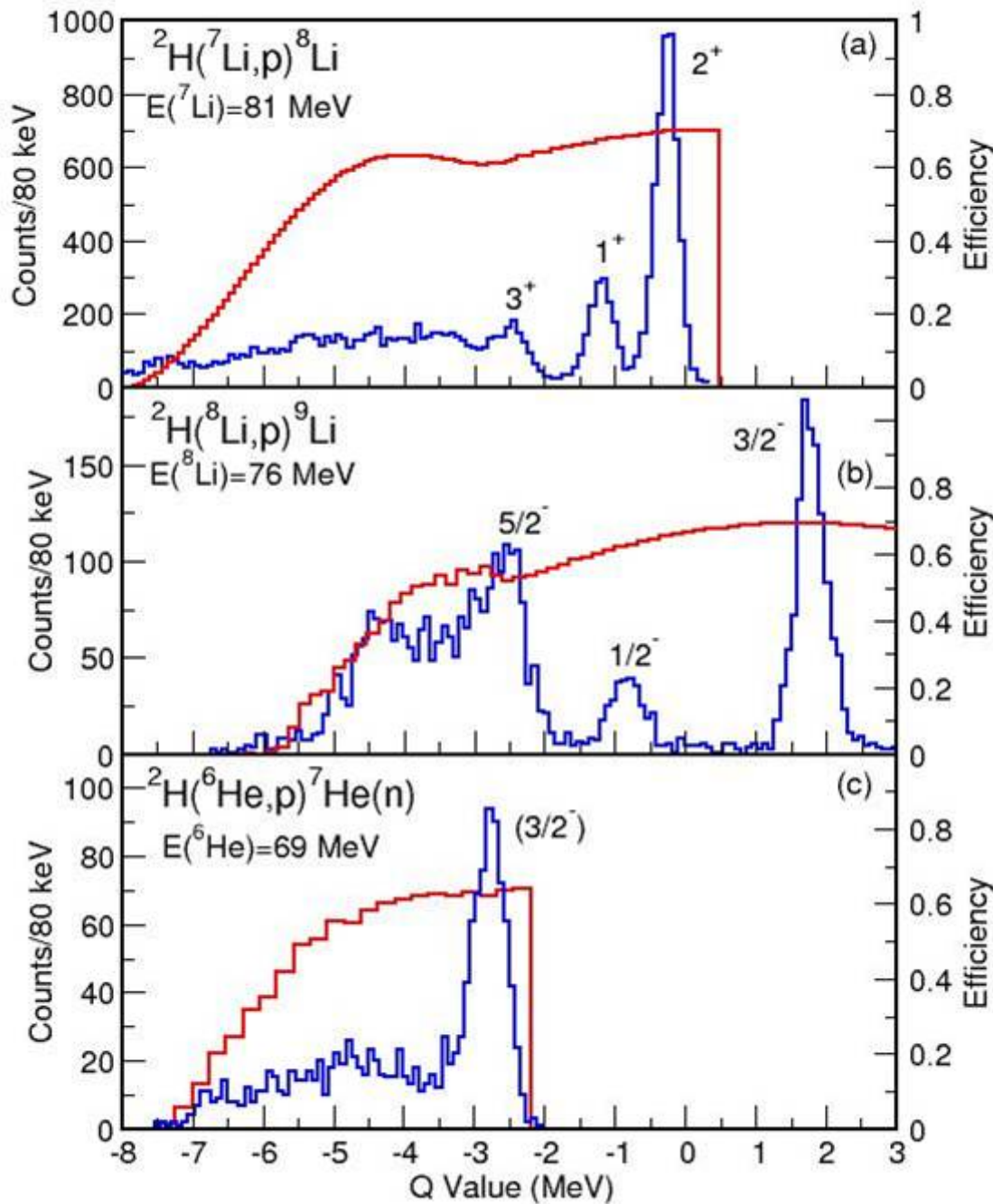
Detectors



Segmented proton detectors

500 μm /1000 μm silicon E Δ E telescope



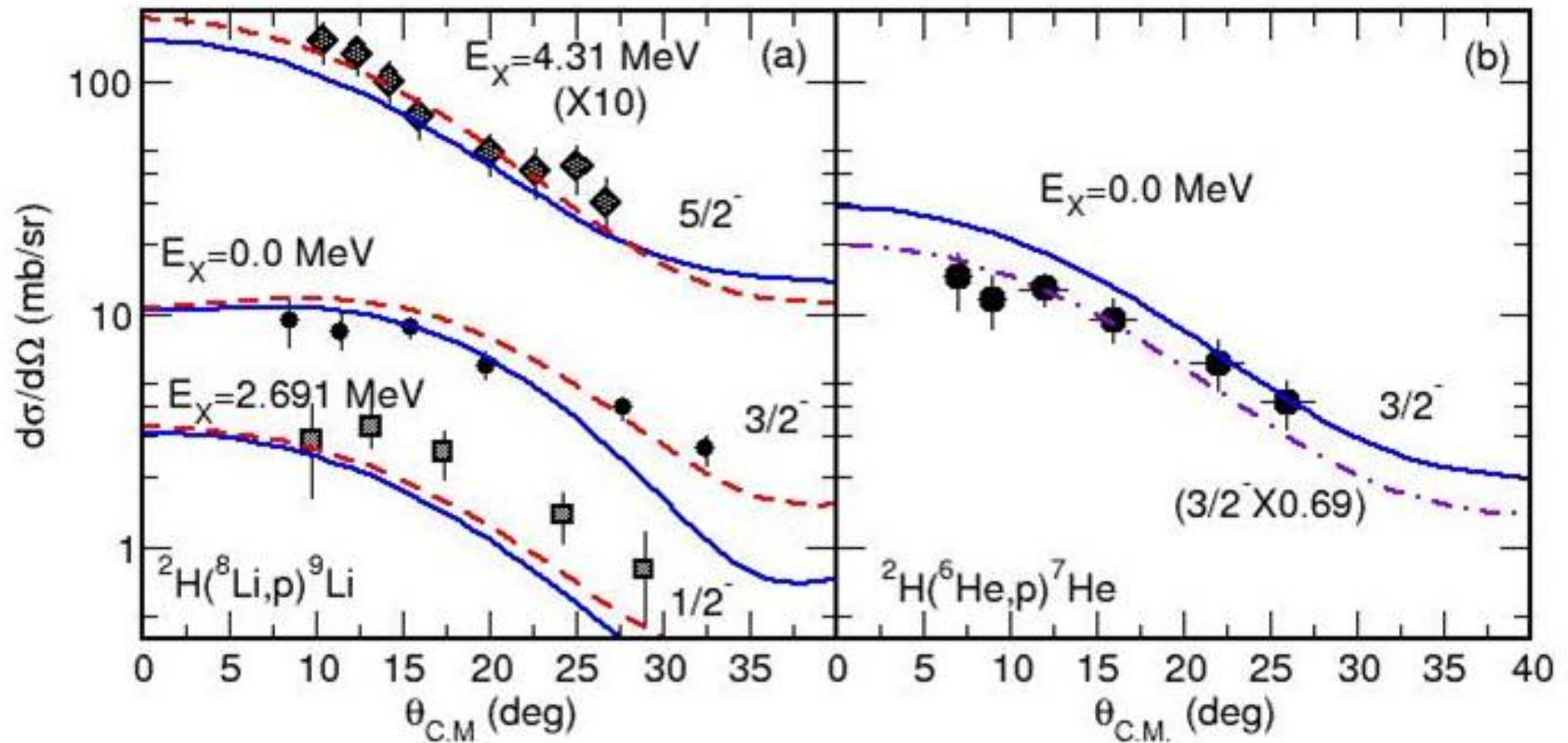


Q-value spectra from (d,p)

p -Li or p - ${}^6\text{He}$ coincidences
 Efficiency from Monte Carlo simulations

PRL **94**, 082502 (2005),
 PRC **72**, 061301(R) (2005)

(d,p) Angular Distributions - narrow states



${}^2\text{H}({}^8\text{Li},p){}^9\text{Li}$ DWBA calculations: **Red**,
blue curves: QMC predictions with
different OMP, no extra normalization

${}^2\text{H}({}^6\text{He},p){}^7\text{He}_{g.s.}$ DWBA calculations
QMC calculations. **Blue**: no normalization
violet- QMC X 0.69.

Optical-model parameters from Schiffer et al, PRC **164**

Shell structure near ^{132}Sn :

Neutron levels in ^{133}Sn via the $^{132}\text{Sn}(d,p)$

Successful “proof of principle” experiment with ^{124}Sn

PHYSICAL REVIEW C 70, 067602 (2004)

Study of the $^{124}\text{Sn}(d,p)$ reaction in inverse kinematics close to the Coulomb barrier

K. L. Jones,¹ R. L. Kozub,² C. Baktash,³ D. W. Bardayan,³ J. C. Blackmon,³ W. N. Catford,⁴ J. A. Cizewski,¹ R. P. Fitzgerald,⁵ M. S. Johnson,⁶ R. J. Livesay,⁷ Z. Ma,⁸ C. D. Nesaraja,² D. Shapira,³ M. S. Smith,³ J. S. Thomas,¹ and D. W. Visser⁵

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³Physics Division, Oak Ridge National Laboratory, Oak Ridge, Tennessee 37831, USA

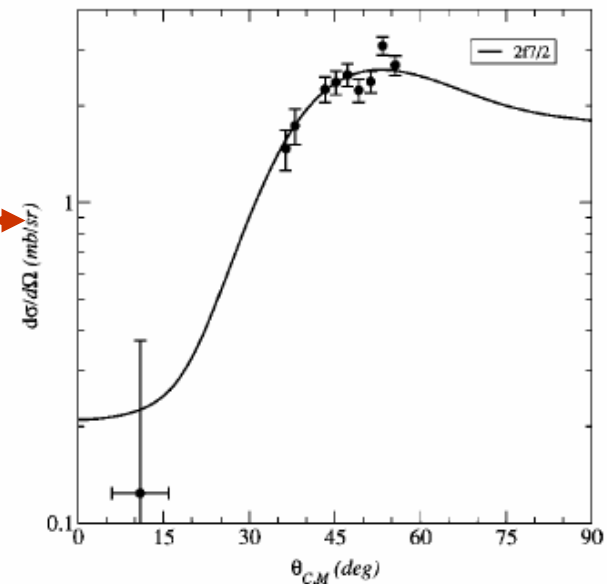
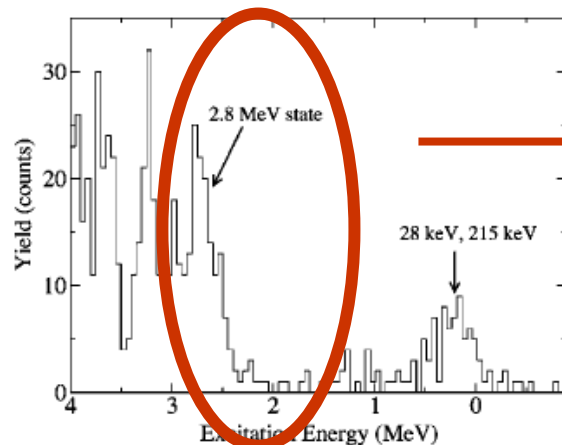
⁴Physics Department, University of Surrey, Guildford, Surrey, GU2 7XH, United Kingdom

⁵Department of Physics and Astronomy, University of North Carolina, Chapel Hill, North Carolina 27599, USA

⁶Oak Ridge Associated Universities, Oak Ridge, Tennessee 37831, USA

⁷Colorado School of Mines, Golden, Colorado 80401, USA

⁸Department of Physics and Astronomy, University of Tennessee, Knoxville, Tennessee 37996, USA



THE experiment with an improved detector array should be running, as we speak, at HRIBF

PAIRING IN NUCLEI

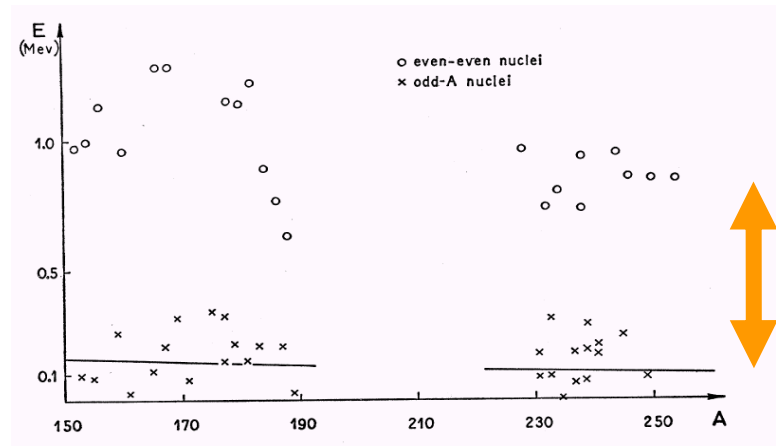
Possible Analogy between the Excitation Spectra of Nuclei and Those of the Superconducting Metallic State

A. BOHR, B. R. MOTTELSON, AND D. PINES*

Institute for Theoretical Physics, University of Copenhagen, Copenhagen, Denmark, and Nordisk Institut for Teoretisk Atomfysik, Copenhagen, Denmark

(Received January 7, 1958)

The evidence for an energy gap in the intrinsic excitation spectrum of nuclei is reviewed. A possible analogy between this effect and the energy gap observed in the electronic excitation of a superconducting metal is suggested.

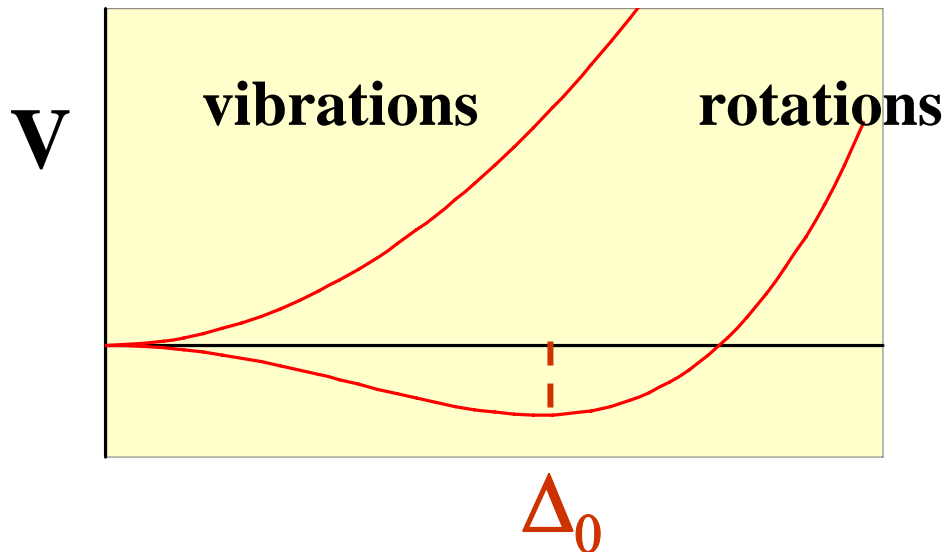


- Gap in the excitation spectra of even- A nuclei
- Odd-even mass differences
- Rotational moments of inertia

Relevance to other finite Fermion systems such as ^3He clusters, Fermi-gas condensates, quantum dots, metal clusters,

Two-nucleon Transfer Reactions

Generalized densities $a+a^+$, aa represent the pair field and in close analogy to the collective excitations corresponding to the ordinary density, they can give rise to collective modes.



$$\Delta = G \langle \sum a_v^+ a_{\bar{v}}^+ \rangle$$

Two particle transfer reactions like (t,p) or (p,t), where 2 neutrons are deposited or picked up at the same point in space provide an specific tool to probe the amplitude of this collective motion. The transition operator $\langle f | a+a^+ | i \rangle$ will be proportional to the pair density of the nucleus.

NOTE ON THE TWO-NUCLEON STRIPPING REACTION

SHIRO YOSHIDA †

Radiation Laboratory, University of Pittsburgh, Pittsburgh, Pennsylvania ††

Received 9 February 1962

Abstract: The magnitude of the two-nucleon s interaction model. The calculation also is a types of reaction a collective enhancement o

* *Proceedings Int. Symp. on nuclear structure*²⁰
(IAEA, Vienna) 1968
p. 179

**PAIR CORRELATIONS AND
DOUBLE TRANSFER REACTIONS**

A. BOHR
THE NIELS BOHR INSTITUTE,
UNIVERSITY OF COPENHAGEN,
COPENHAGEN, DENMARK

PHYSICS REPORT

ISOVECTOR PAIRING VIBRATIONS

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State University of New York at Stony Brook, Physics Department, Stony Brook, New York 11794, USA and
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and

R.A. BROGLIA

*The Niels Bohr Institute, University of Copenhagen, DK-2100 Copenhagen Ø, Denmark and
State University of New York at Stony Brook, Physics Department, Stony Brook, New York 11794, USA*

and

Ole HANSEN and O. NATHAN

The Niels Bohr Institute, University of Copenhagen, DK-2100 Copenhagen Ø, Denmark

Neutron Proton Pairing

NP Pairing

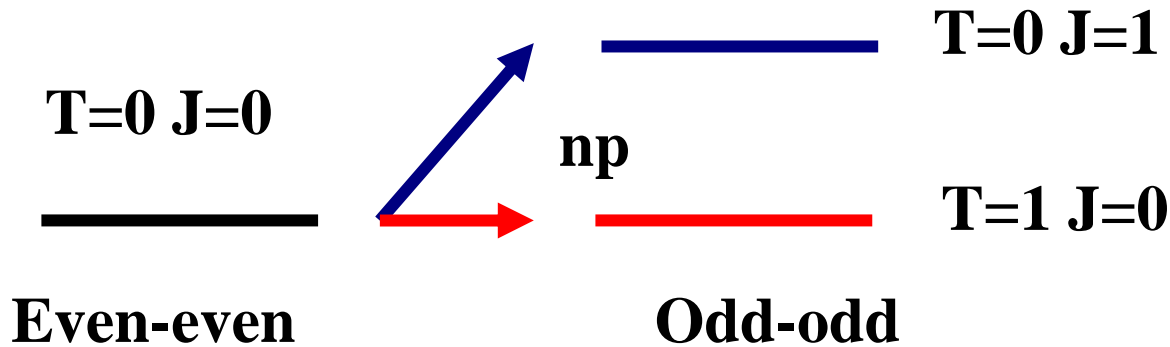
- ★ N=Z nuclei, unique systems to study np correlations
 - As you move out of N=Z nn and pp pairs are favored
- ★ Role of isoscalar ($T=0$) and isovector ($T=1$) pairing
 - Large spatial overlap of n and p
 - Pairing vibrations (normal system)
 - Pairing rotations (superfluid system)
- ★ Does isoscalar pairing give rise to collective modes?
- ★ What is (are) the “smoking-gun(s)”?
 - Binding energy differences
 - Ground states of odd-odd self-conjugate nuclei
 - Rotational properties: moments of inertia, alignments
 - Two-particle transfer cross-sections



(3He,p) Transfer Reactions

$$\sigma \propto < 1 | T | 0 >^2 \sim \Omega \quad \text{or} \quad \Omega^2$$

$\sigma ?$



$(^3\text{He}, p)$

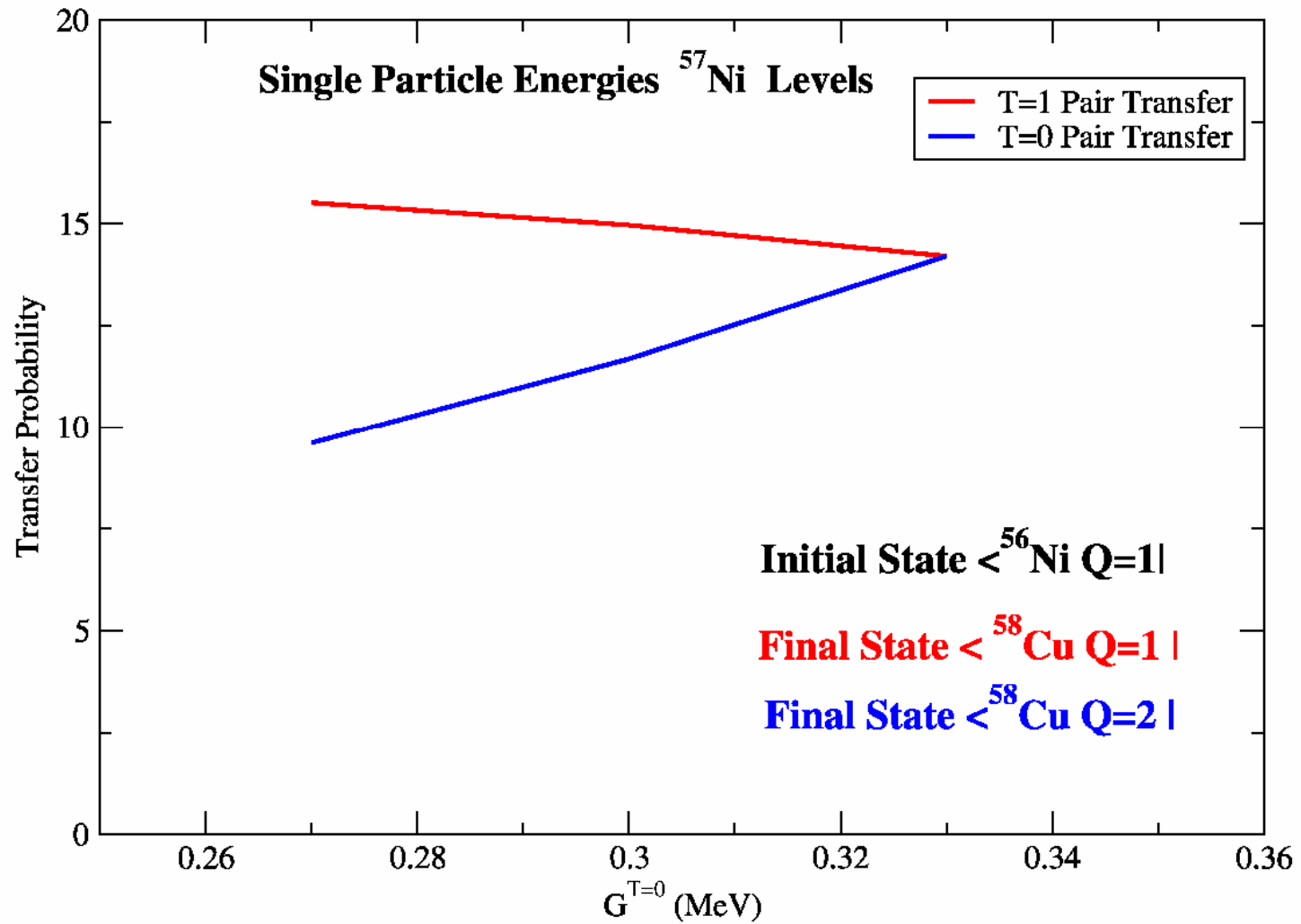
$L=0$ transfer – forward peaked

Measure the np transfer cross section to $T=1$ and $T=0$ states

Both absolute $\sigma(T=0)$ and $\sigma(T=1)$ and relative $\sigma(T=0) / \sigma(T=1)$ tell us about the character and strength of the correlations

n-p Pair Transfer Probability

$$G^{T=1} = 0.33 \text{ MeV}$$



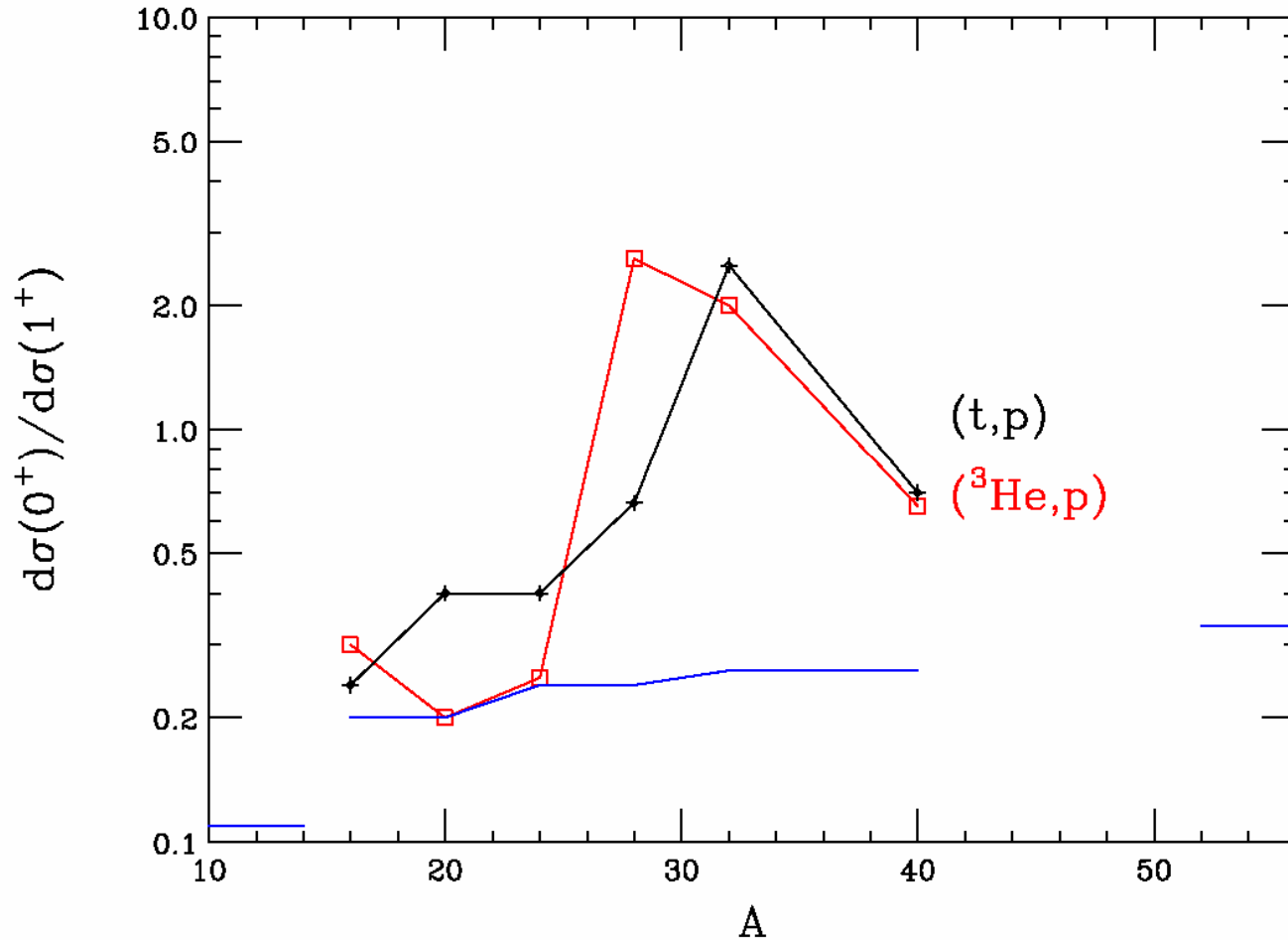
Deuteron Transfer in $N = Z$ NucleiP. Van Isacker,¹ D.D. Warner,² and A. Frank³¹*Grand Accélérateur National d'Ions Lourds, B.P. 55027, F-14076 Caen Cedex 5, France*²*CCLRC Daresbury Laboratory, Daresbury, Warrington WA4 4AD, United Kingdom*³*Instituto de Ciencias Nucleares, UNAM, Apdo. Postal 70-543, 04510 México, D.F. Mexico*

(Received 14 September 2004; published 29 April 2005)

TABLE I. Predicted deuteron-transfer intensities C_T^2 between even-even (EE) and odd-odd (OO) $N = Z$ nuclei in the $SU(4)$ ($b/a = 0$) and $U_T(3) \otimes U_S(3)$ ($|b/a| \gg 1$) limits.

| Limit | Reaction | $C_{T=0}^2$ | $C_{T=1}^2$ |
|--------------|---------------------------|------------------------|------------------------|
| $b/a = 0$ | $EE \rightarrow OO_{T=0}$ | $\frac{1}{2}(N_b + 6)$ | 0 |
| | $EE \rightarrow OO_{T=1}$ | 0 | $\frac{1}{2}(N_b + 6)$ |
| $b/a \ll -1$ | $EE \rightarrow OO_{T=0}$ | $N_b + 3$ | 0 |
| | $EE \rightarrow OO_{T=1}$ | 0 | 3 |
| $b/a \gg +1$ | $EE \rightarrow OO_{T=0}$ | 3 | 0 |
| | $EE \rightarrow OO_{T=1}$ | 0 | $N_b + 3$ |

Systematic of ($^3\text{He},p$) and (t,p) reactions in stable $N=Z$ nuclei



Single-particle estimate $\sim (\text{spin}) \times (^3\text{He}) \times (\text{LS} \rightarrow \text{jj})$

Study of the $^{56}\text{Ni}(d,p)^{57}\text{Ni}$ Reaction and the Astrophysical $^{56}\text{Ni}(p,\gamma)^{57}\text{Cu}$ Reaction Rate

K. E. Rehm,¹ F. Borasi,¹ C. L. Jiang,¹ D. Ackermann,¹ I. Ahmad,¹ B. A. Brown,² F. Brumwell,¹ C. N. Davids,¹
 P. Decroock,¹ S. M. Fischer,¹ J. Görres,³ J. Greene,¹ G. Hackmann,¹ B. Harss,¹ D. Henderson,¹ W. Henning,¹
 R. V. F. Janssens,¹ G. McMichael,¹ V. Nanal,¹ D. Nisius,¹ J. Nolen,¹ R. C. Pardo,¹ M. Paul,⁴ P. Reiter,¹ J. P. Schiffer,¹
 D. Seweryniak,¹ R. E. Segel,⁵ M. Wiescher,³ and A. H. Wuosmaa¹

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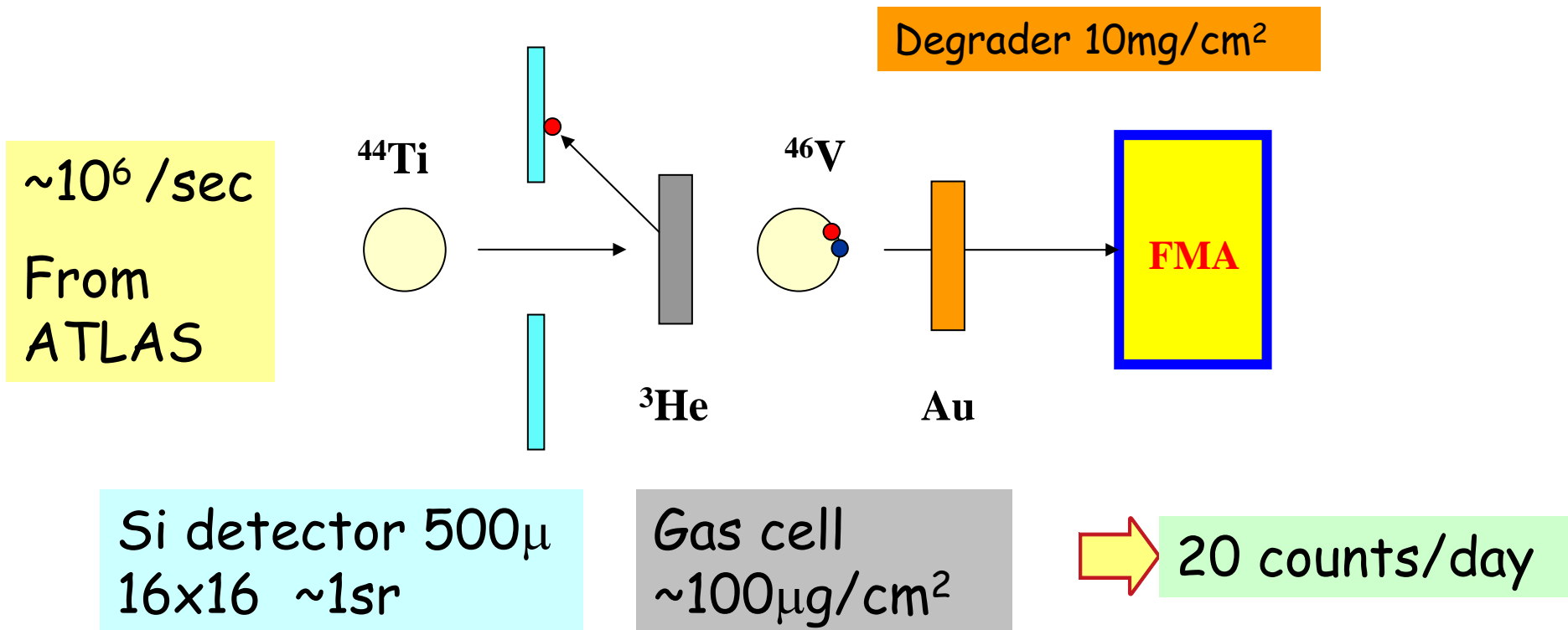
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(Received 29 August 1997)



Proof of principle

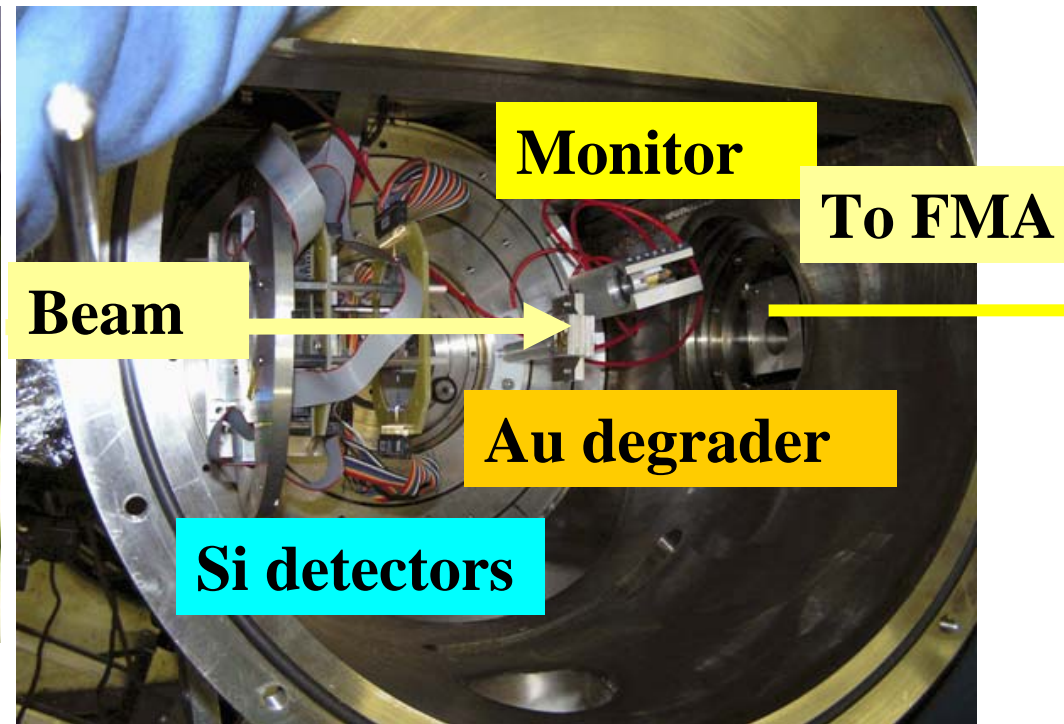
A.O.Macchiavelli¹, E.Rehm², P.Fallon¹, M.Cromaz¹, I.Ahmad¹, C.N.Davis², J.Greene², R.V.F.Janssens², C.L.Jiang²,
E.F.Moore², R.Pardo², D.Seweryniak², J.P.Schiffer², A.Wuosmaa³, J.Cizewski³,

¹ *Lawrence Berkeley National Laboratory*

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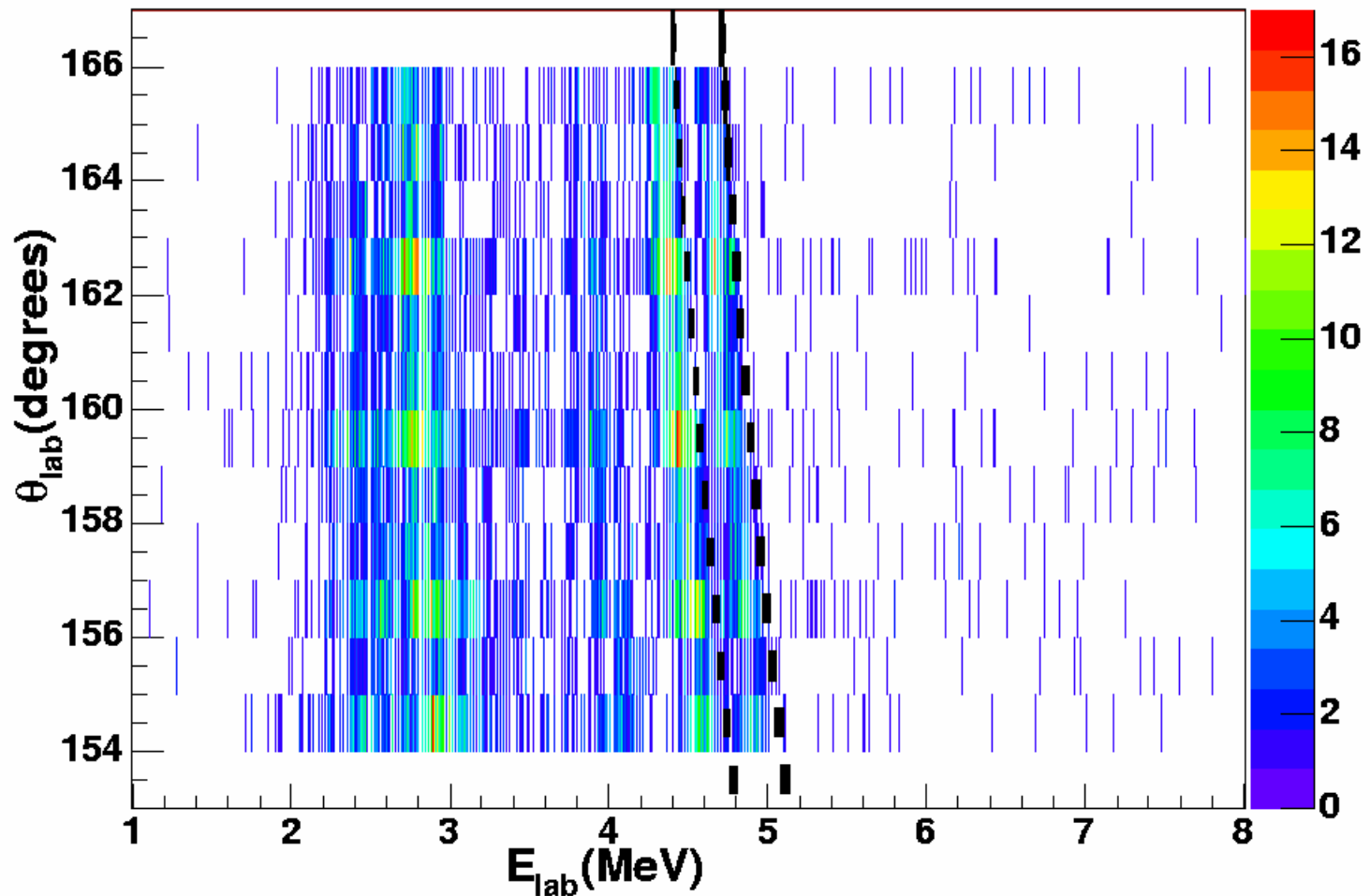
³ *Western Michigan University*

³ *Rutgers University*

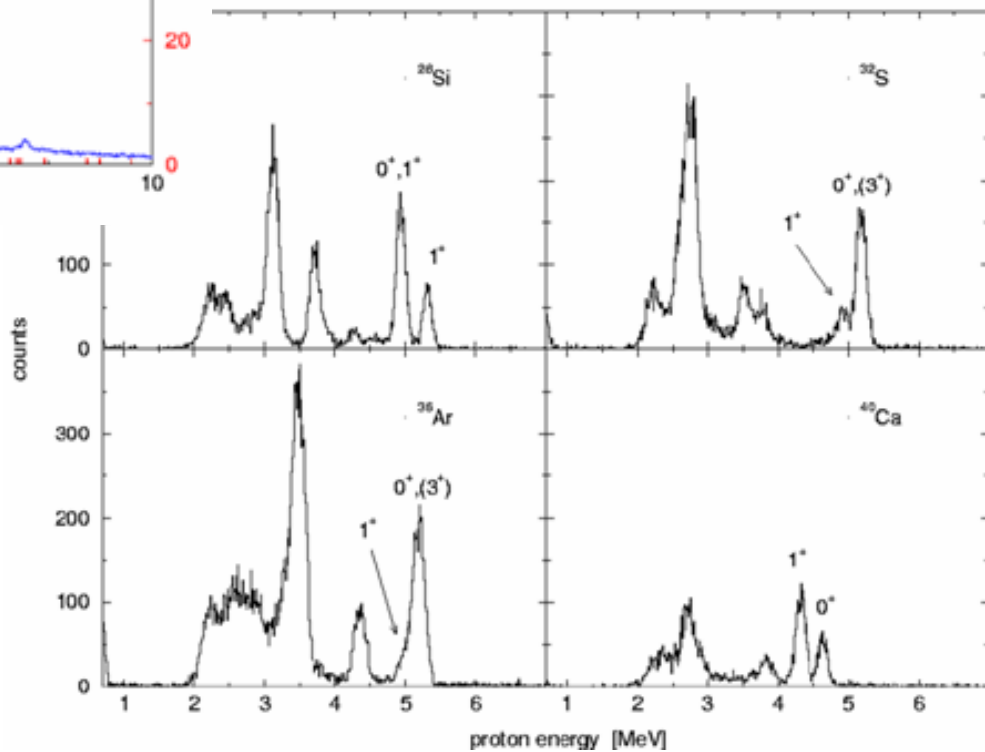
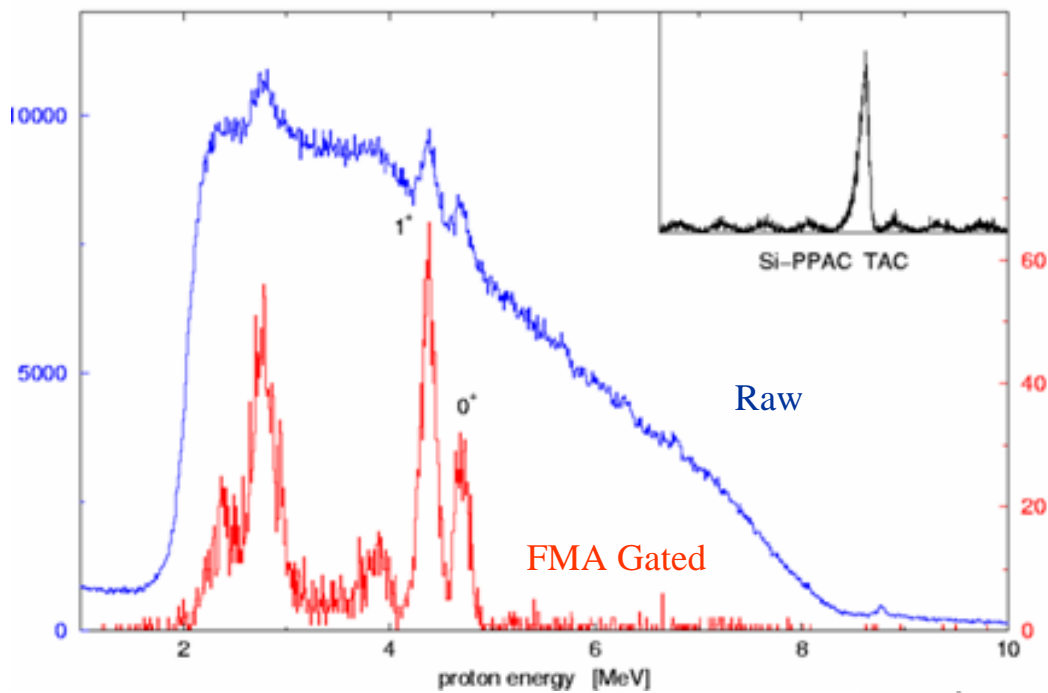


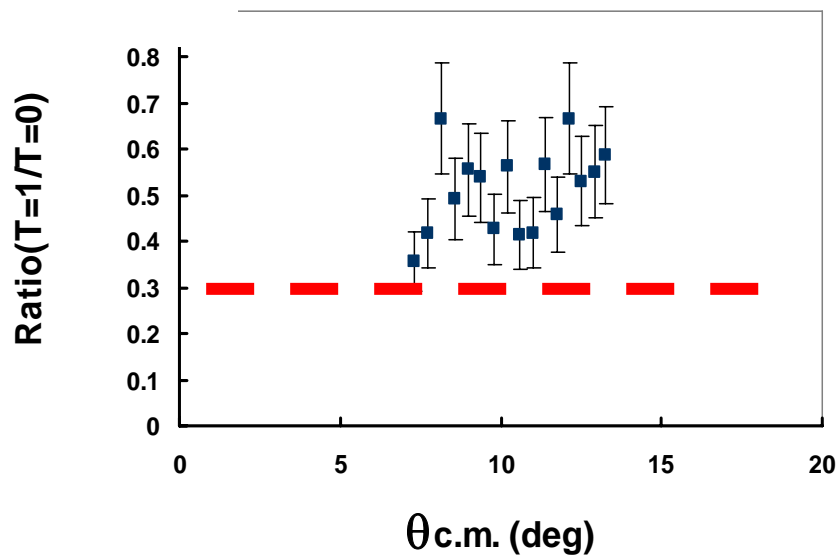
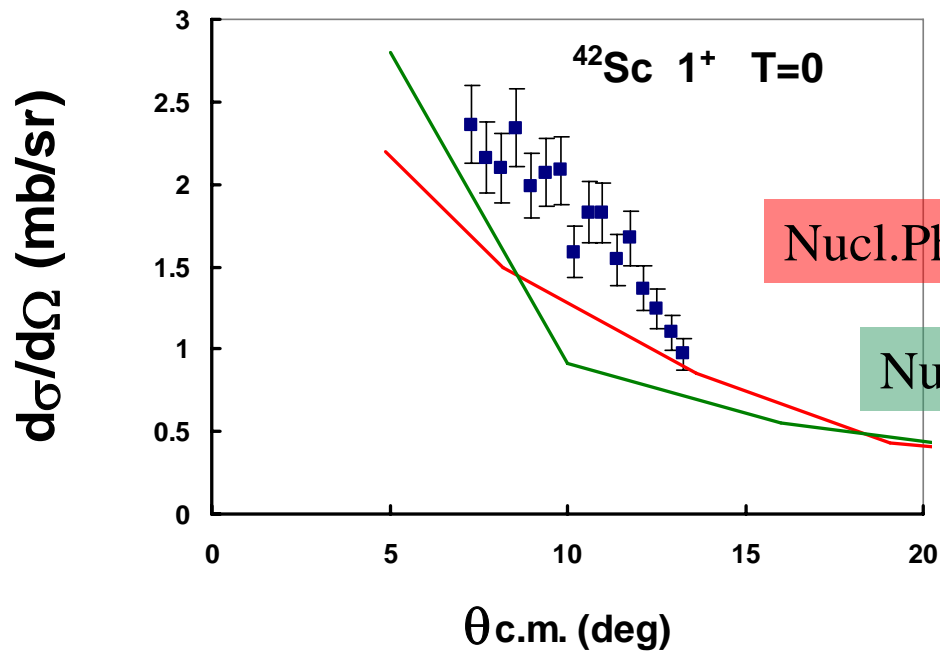
Proof of principle

Inverse kinematics -
Successful test experiment with stable beams



$^{40}\text{Ca}(^3\text{He},p)$ @ 220 MeV

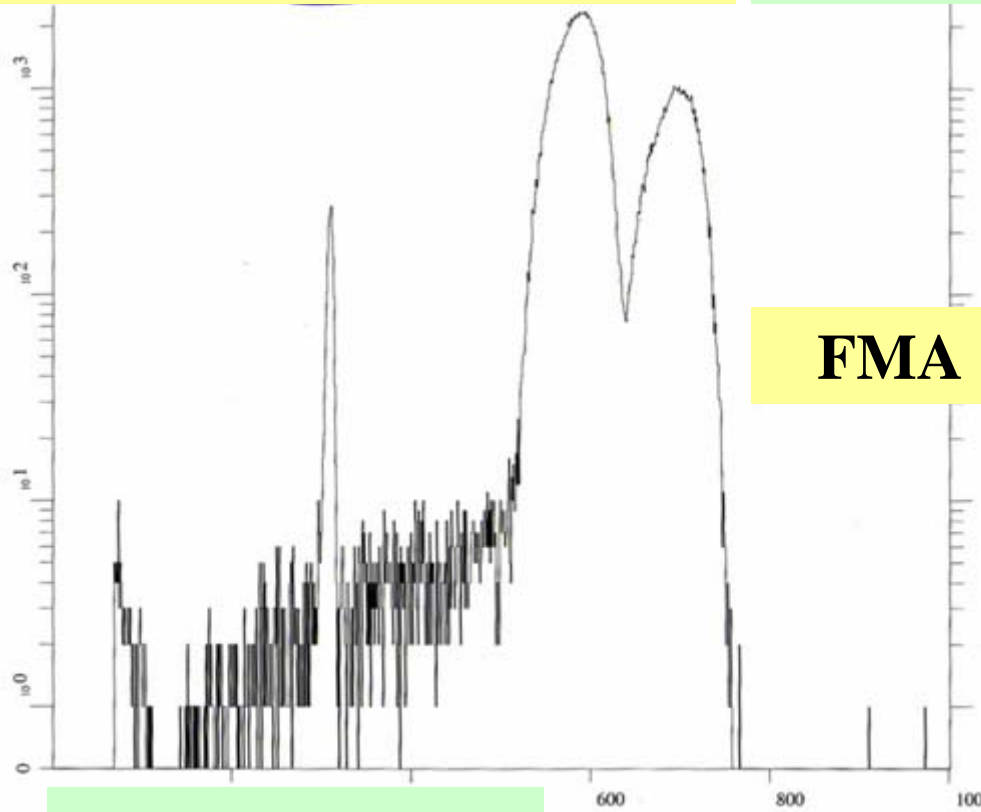




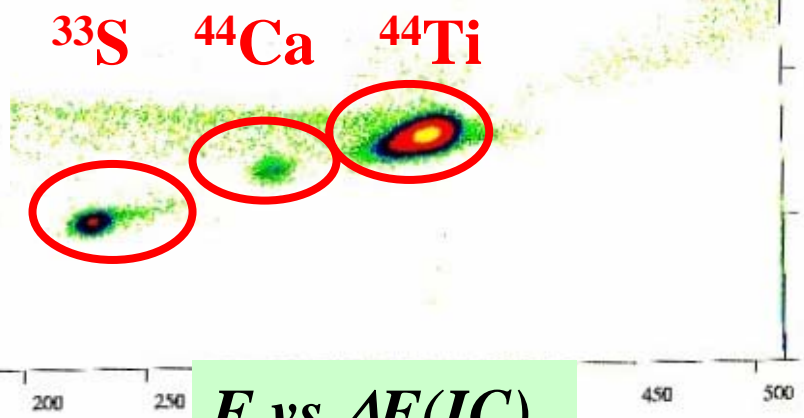
242 MeV ^{44}Ti Beam

$^{44}\text{Ti} / ^{44}\text{Ca} = 2.5 \sim 890/\text{sec}$

$\sim 1.1 \times 10^6 \text{ }^{44}\text{Ti}/\text{sec}$



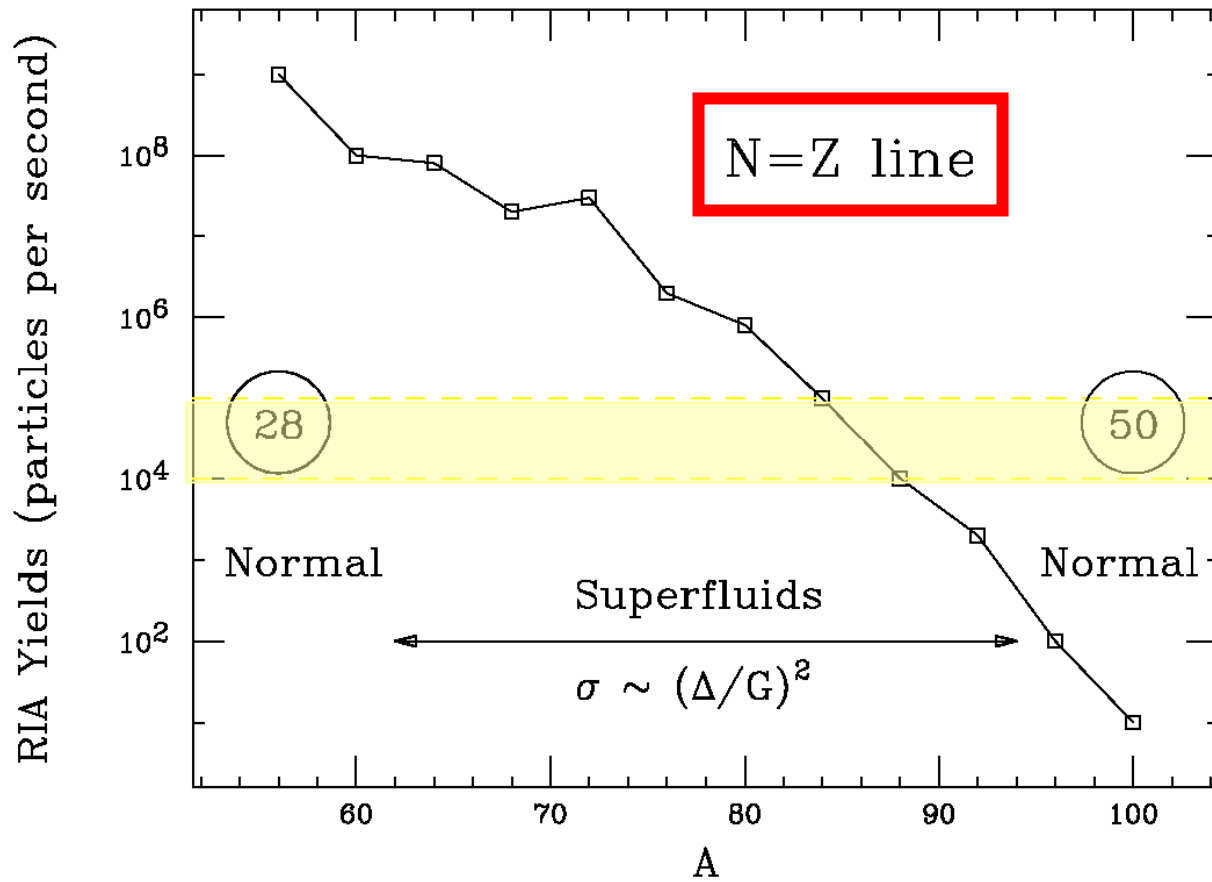
FMA Settings 107MeV, A=44, q=20



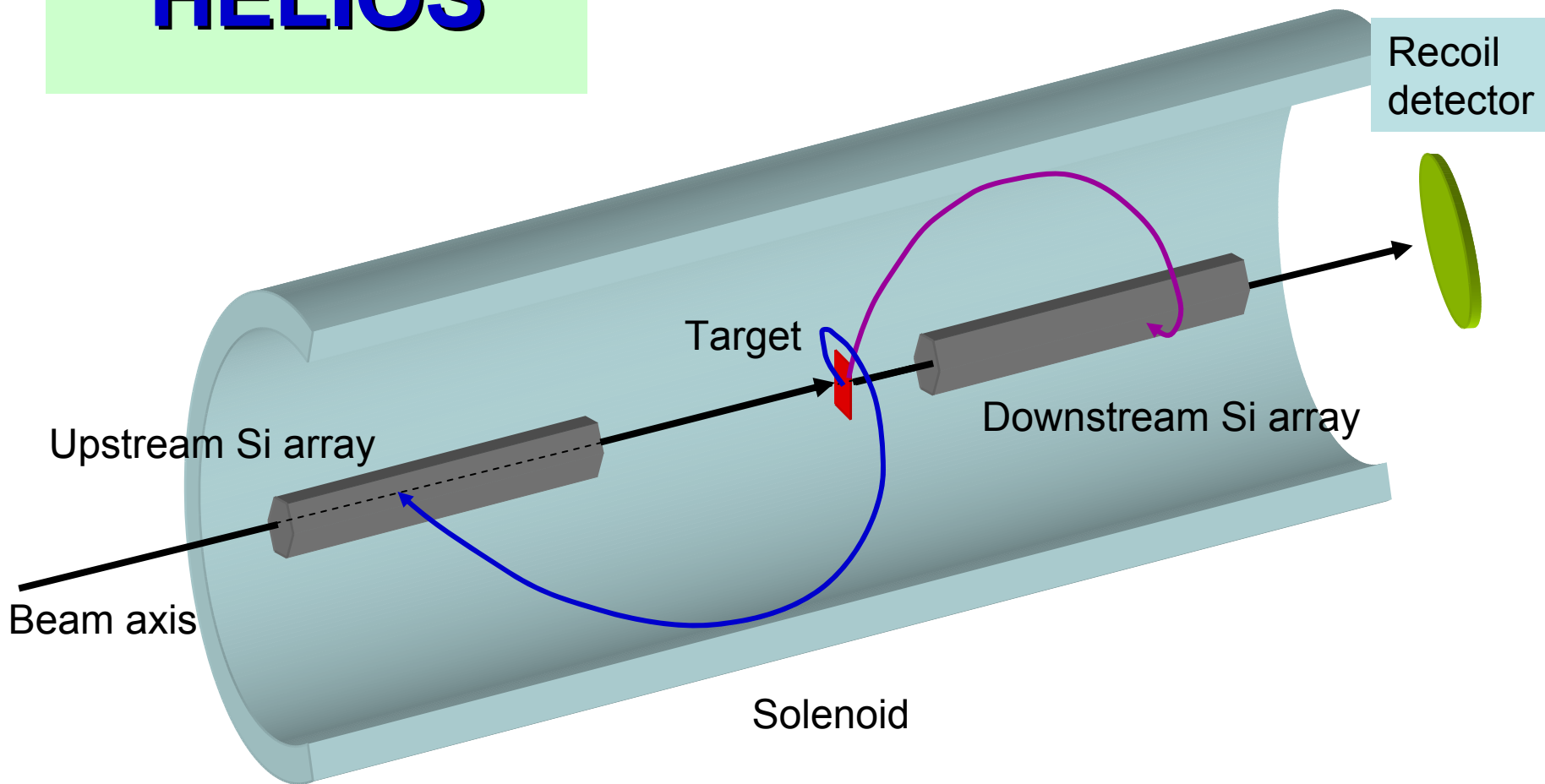
E Monitor @ 0°
X1800 Attenuator

E vs $\Delta E(IC)$

Looking ahead ...



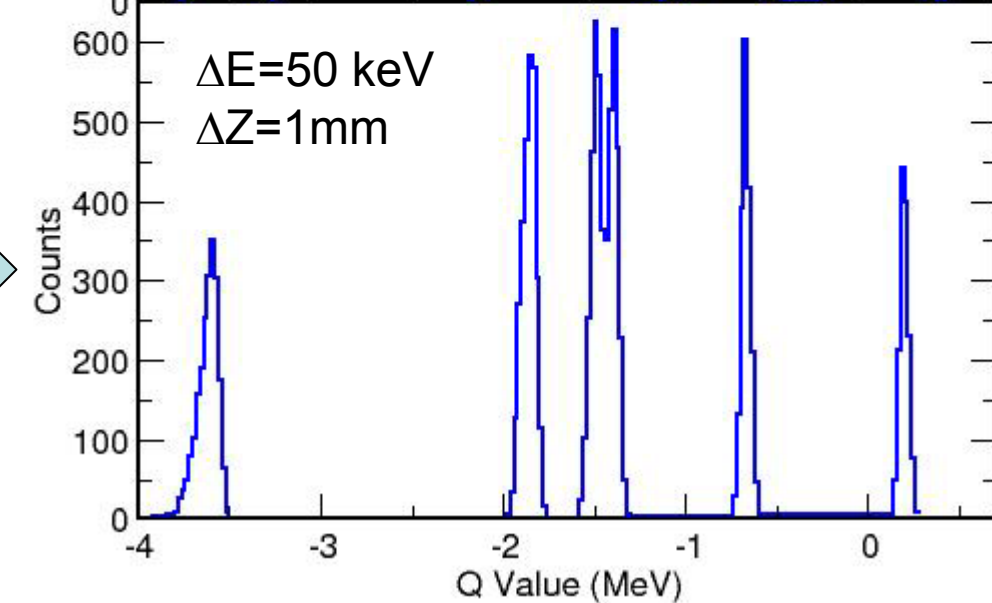
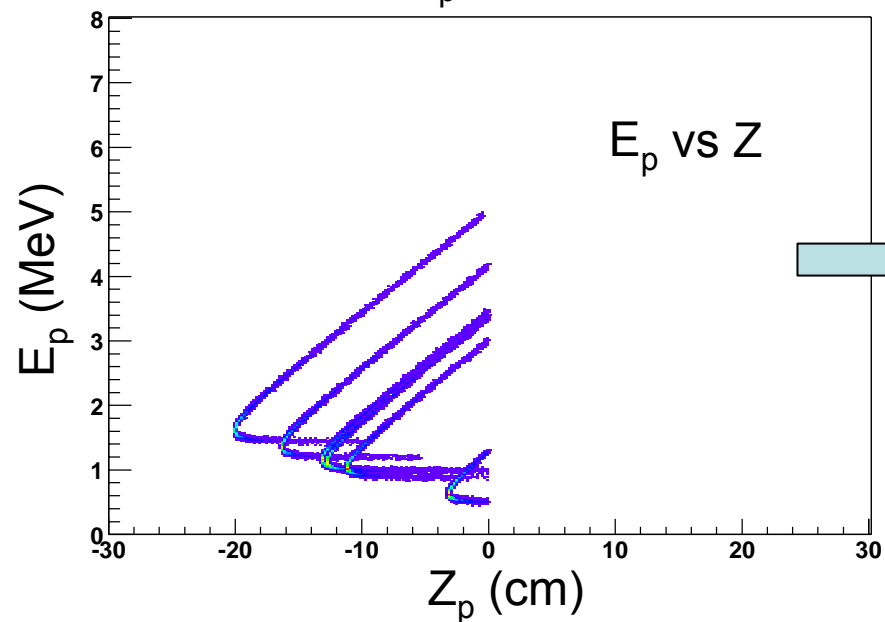
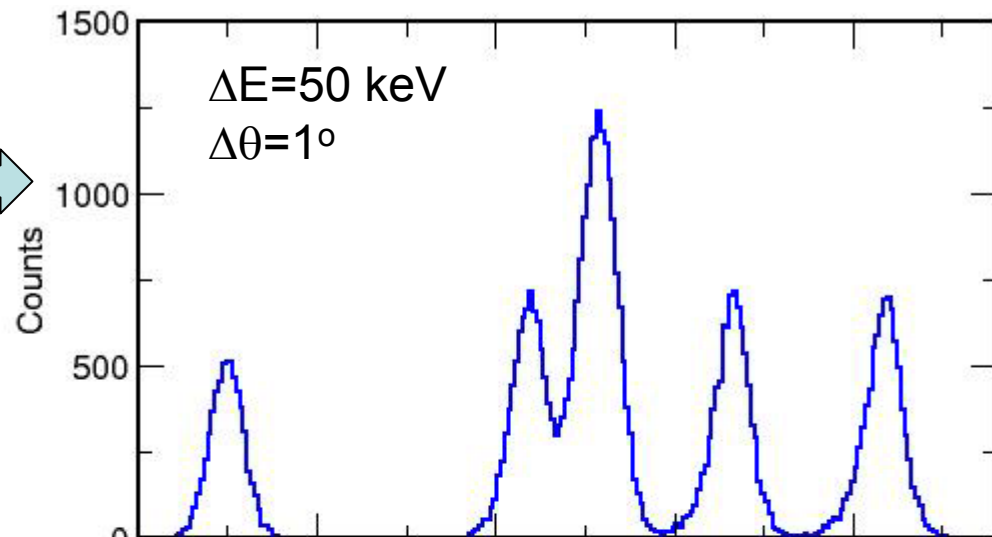
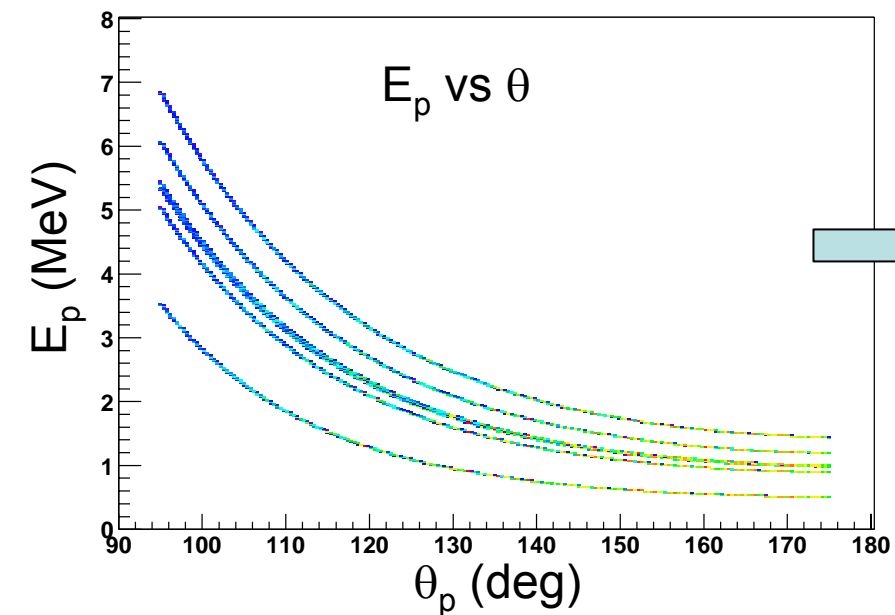
HELIOS



Schematic
design

Improved resolution for (d,p)

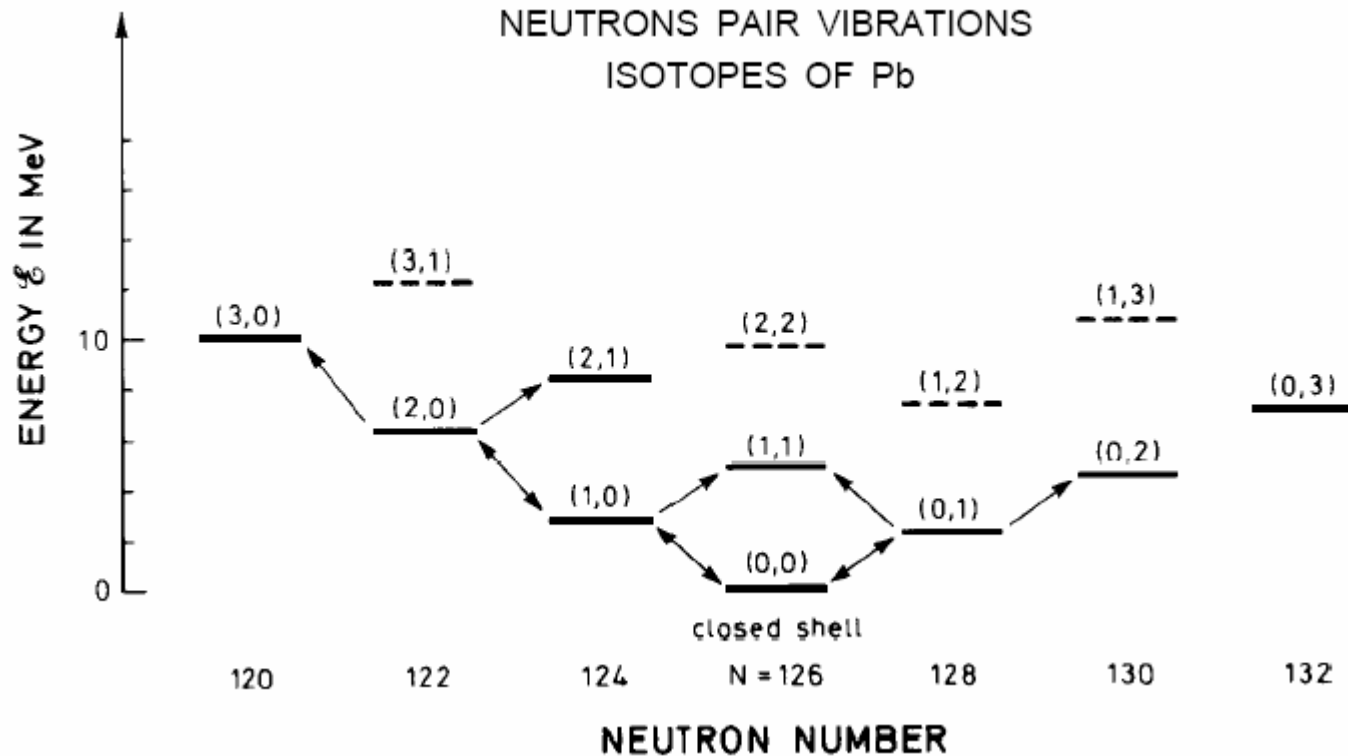
Q Value



Pairing Phase Transition

Pair-Vibrational Structures

(Nobel Lecture, Ben R. Mottelson, 1975 “Elementary Modes of Excitation in the Nucleus”)



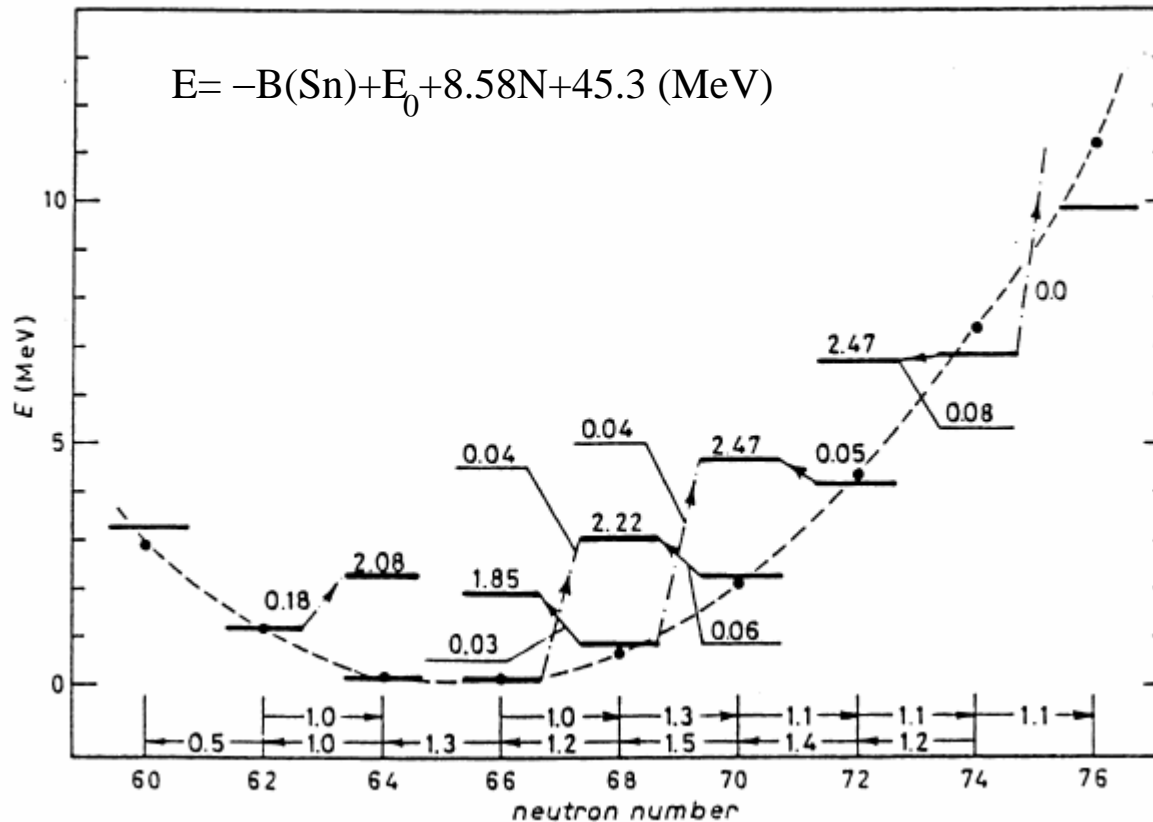
- Near closed shell nuclei (like ^{208}Pb) no static deformation of pair field.
- Corresponds to the “normal” nuclear limit.
- Fluctuations give rise to a vibrational-like excitation spectrum.
- Enhanced pair-addition and pair-removal cross-sections seen in (t,p) and (p,t) reactions (indicated by arrows).



- Large anharmonicities in spectrum must be accounted for.

Pair-Rotational Structures

(R. A. Broglia, J. Terasaki, and N. Giovanardi, Phys. Rep. 335 (2000) 1)



- Many like-nucleon pairs outside a closed-shell configuration (e.g. ^{116}Sn) gives rise to a static deformation of the pair field.
- Corresponds to the “superconducting” limit.
- Rotational-like (parabolic – dashed line) spectrum formed by sequence of ground states of even- N neighbors.
- Angular variable in rotational motion is gauge angle, ϕ .

Critical-Point Descriptions of Shape Transitions

F. Iachello, Phys. Rev. Lett. 85 (2000) 3580, Phys. Rev. Lett. 87 (2001) 052502.

PRL 96, 032501 (2006)

PHYSICAL REVIEW LETTERS

week ending
27 JANUARY 2006

Critical-Point Description of the Transition from Vibrational to Rotational Regimes in the Pairing Phase

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²Dipartimento di Fisica Galileo Galilei, INFN, Sezione di Padova, Padova, Italy

³Physik Department E12, Technische Universität München, Garching, Germany

The Collective Pairing Hamiltonian

D.R. Bès, R.A. Broglia, R.P.J. Perazzo, K. Kumar, Nucl. Phys. A 143 (1970) 1.

$$-\frac{\hbar^2}{2B} \frac{\partial^2 \psi}{\partial \alpha^2} - \frac{\hbar^2}{2B\alpha} \frac{\partial \psi}{\partial \alpha} + \left(\frac{\hbar^2 M^2}{8B\alpha} + V(\alpha) - E \right) \psi = 0$$

α is the deformation of the pair field (can be related to the gap parameter, Δ).

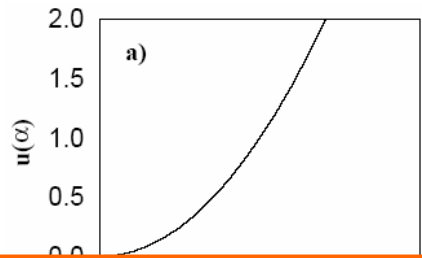
B is a mass parameter. $M=(A-A_0)$ (number of particles, A , relative to reference, A_0).

$V(\alpha)$ is the potential.

Analogy between Shapes and Pairing

$$\frac{\partial^2 \psi}{\partial \alpha^2} + \frac{1}{\alpha} \frac{\partial \psi}{\partial \alpha} + \left(\varepsilon - u(\alpha) - \frac{M^2}{4\alpha^2} \right) \psi = 0$$

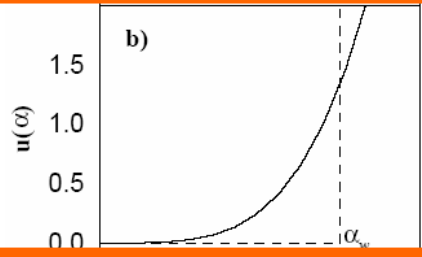
Vibrational Limit
("Normal")



$$u(\alpha) = \alpha^2$$

E(2)

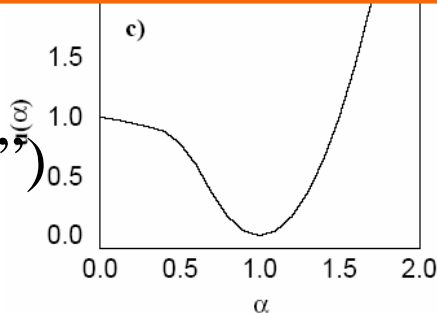
Transitional
("Critical-Point")



$$u(\alpha) = 0, \alpha \leq \alpha_w$$

$$u(\alpha) = \infty, \alpha > \alpha_w$$

Rotational Limit
("Superconducting")



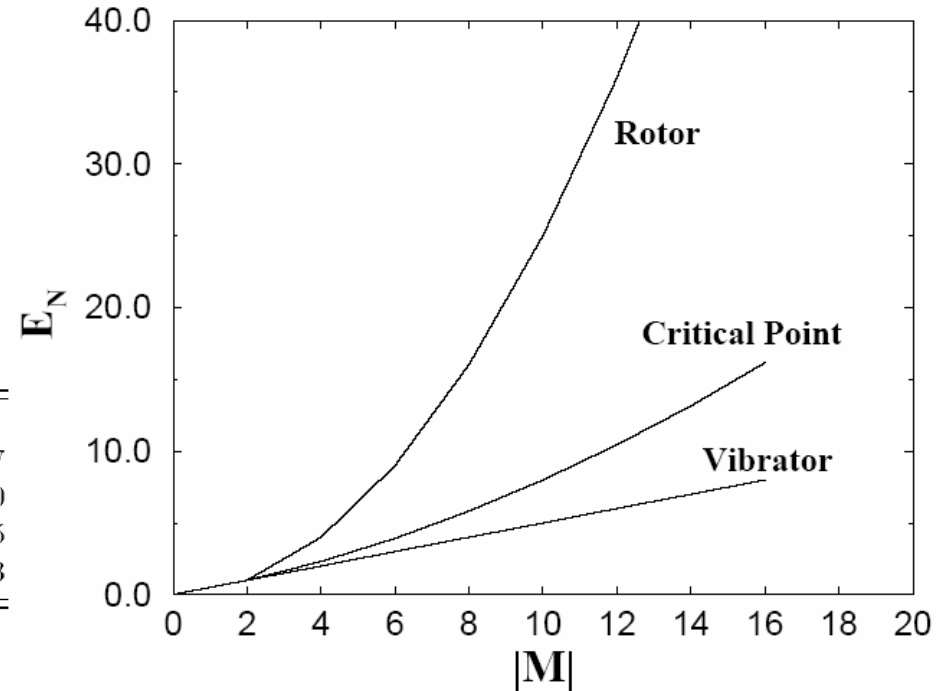
$$u(\alpha) = \delta_{\alpha_0}$$

The Reduced Energy Spectrum

Normalizing energies of excited states to that of the first excited state:

$$E_N = \frac{x_{\xi, M}^2 - x_{1,0}^2}{x_{1,2}^2 - x_{1,0}^2}$$

| | $\xi=1$ | $\xi=2$ | $\xi=3$ | $\xi=4$ |
|---------|---------|---------|---------|---------|
| $ M =0$ | 0.00 | 2.77 | 7.77 | 14.97 |
| $ M =2$ | 1.00 | 4.88 | 10.98 | 19.30 |
| $ M =4$ | 2.31 | 7.31 | 14.52 | 23.95 |
| $ M =6$ | 3.92 | 10.06 | 18.39 | 28.93 |

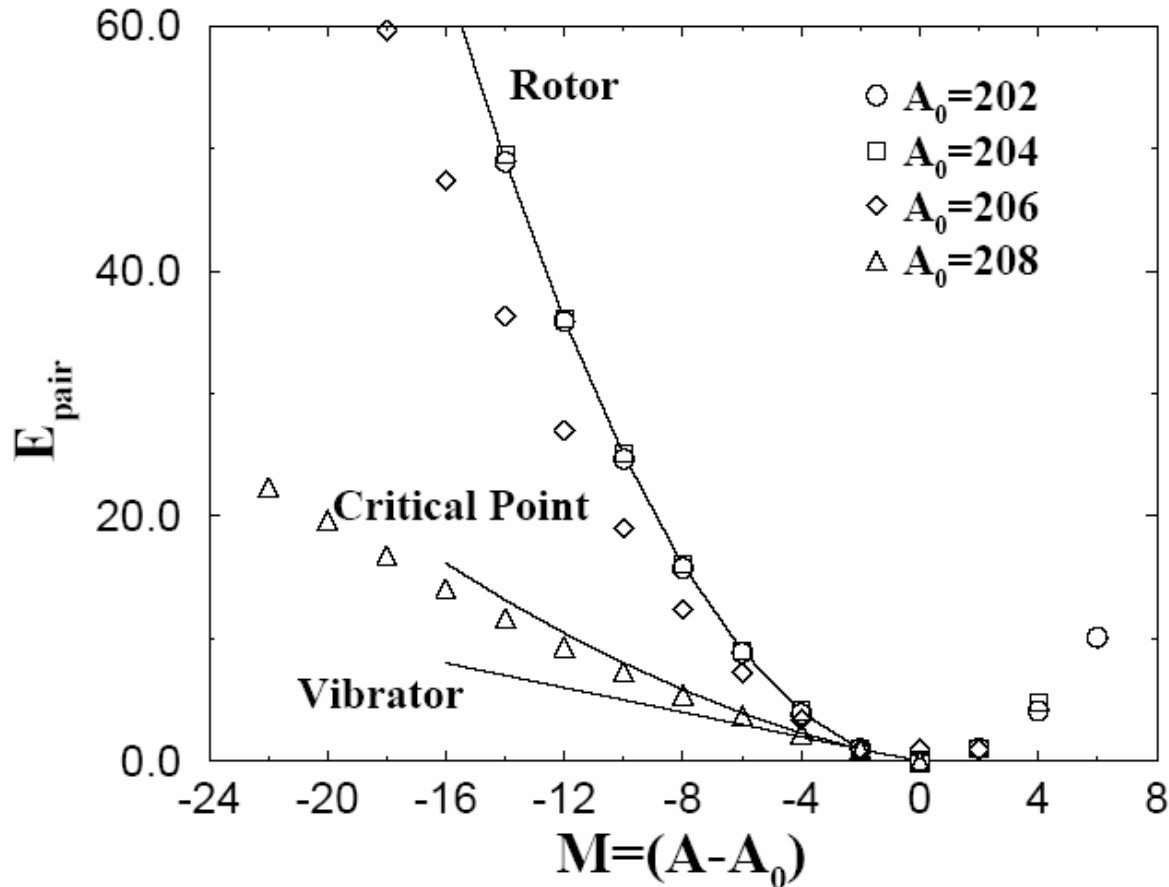


- $\xi=1$ sequence of states correspond empirically to the sequence formed by the 0^+ ground-states of neighboring even-even nuclei along isotopic or isotonic chain.
- $\xi=1$ sequence of states follows behavior between the linear dependence of a harmonic vibrator and the parabolic dependence of a deformed rotor, as expected.
- $\xi>1$ correspond to excited 0^+ states formed from pair excitations.

Comparison with Data

Empirical neutron pairing energy defined as: $E_{pair} = [\varepsilon(A) - \varepsilon(A_0)] - C \cdot (A - A_0)$

$\varepsilon(A) - \varepsilon(A_0)$ is difference in mass excess between isotope of mass A and reference nucleus of mass A_0 (G. Audi et al., Nucl. Phys. A 729 (2003) 337)



Only a few nucleons outside of the closed shell are required for a static pair deformation (“superconductivity”) and pair rotational sequences develop.

A comment on anharmonicities

From Bohr & Mottelson Vol. II pag. 646

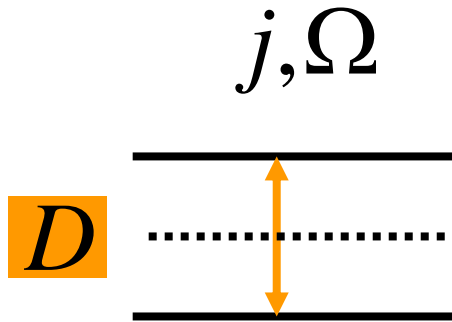
$$E \sim n\hbar\omega_0 + \frac{1}{2}V_{--}n(n-1)$$

$$V_{--} \sim 700\text{keV}$$

Our simple estimate is

$$V_{--} \sim (2.31 - 2)\hbar\omega_0 \sim 800\text{keV}$$

A simple microscopic model: Two j-shells



$$H = \frac{1}{2} D(N_2 - N_1) - G\Omega(A_1^+ + A_2^+)(A_1 + A_2)$$

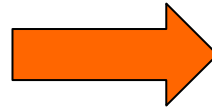
$$x = 2\Omega G / D$$

“Control parameter”

$$\Omega \sim A^{2/3}$$

$$G \sim 20/A$$

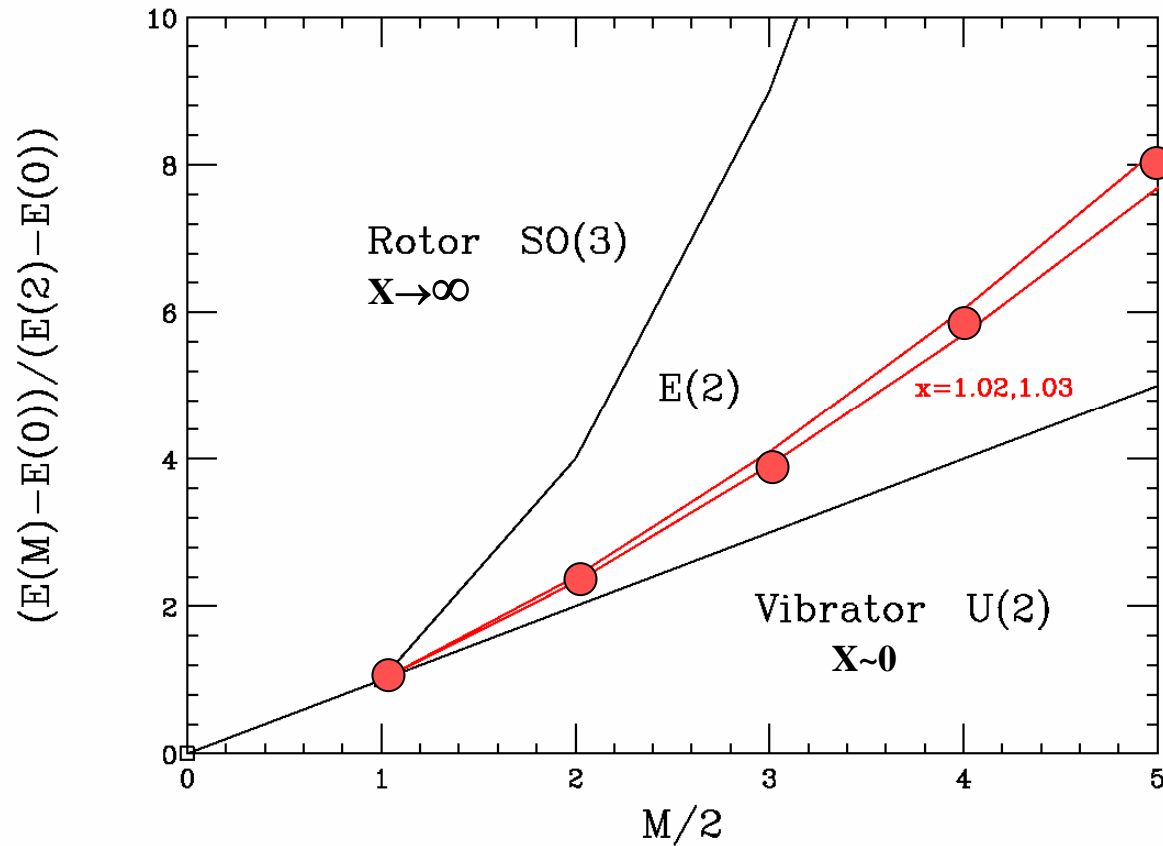
$$D \sim 41/A^{1/3}$$



$$x \sim 1$$

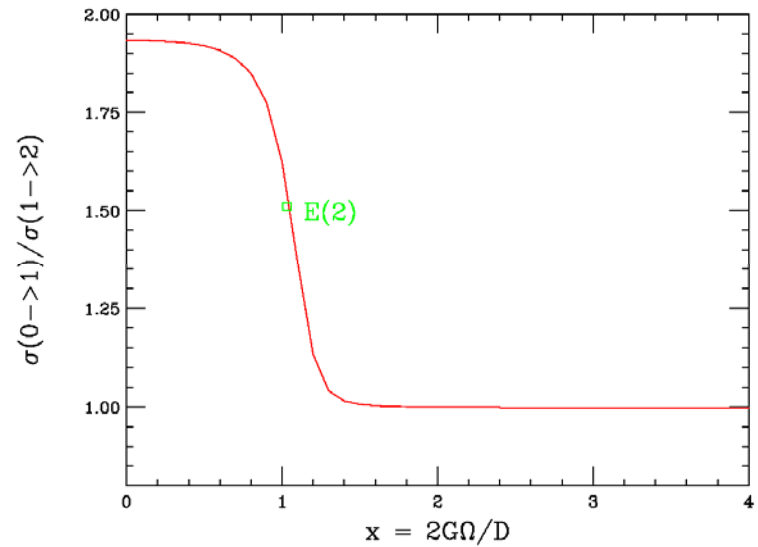
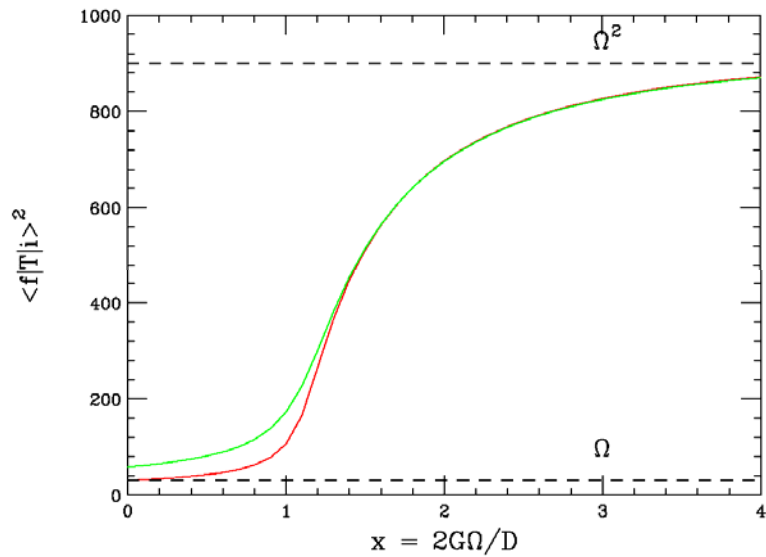
A simple microscopic model: Two j-shells

Energies



A simple microscopic model: Two j-shells

Transition probabilities



Transition Probabilities

The transition matrix elements are related to two-nucleon transfer probabilities:

$$\langle \psi_{\xi', M'} | \hat{O} | \psi_{\xi, M} \rangle = \int \psi_{\xi', M'}^* \hat{O} \psi_{\xi, M} d\tau$$

where, the volume element is: $d\tau = 2B\alpha \cdot d\alpha \cdot d\phi$

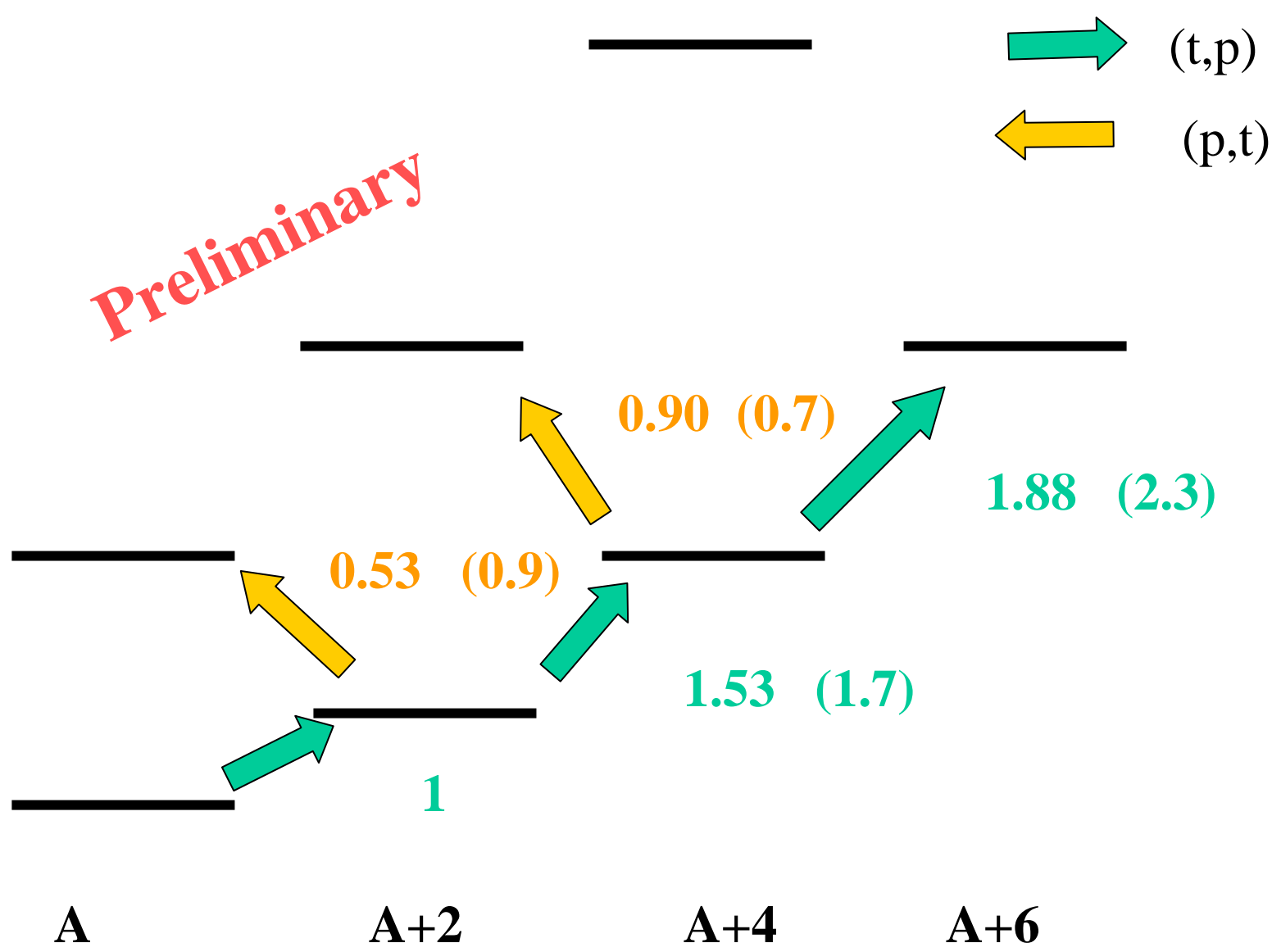
and the pair-transfer operator is: $\hat{O} = \alpha_{\pm 2} = \alpha e^{\pm 2i\phi}$

One then finds:

$$\langle \psi_{\xi', M'} | \hat{O} | \psi_{\xi, M} \rangle \propto \int_0^{\alpha_w} \psi_{\xi', M'}^* \alpha^2 \psi_{\xi, M} d\alpha$$

The integrals can be solved numerically.

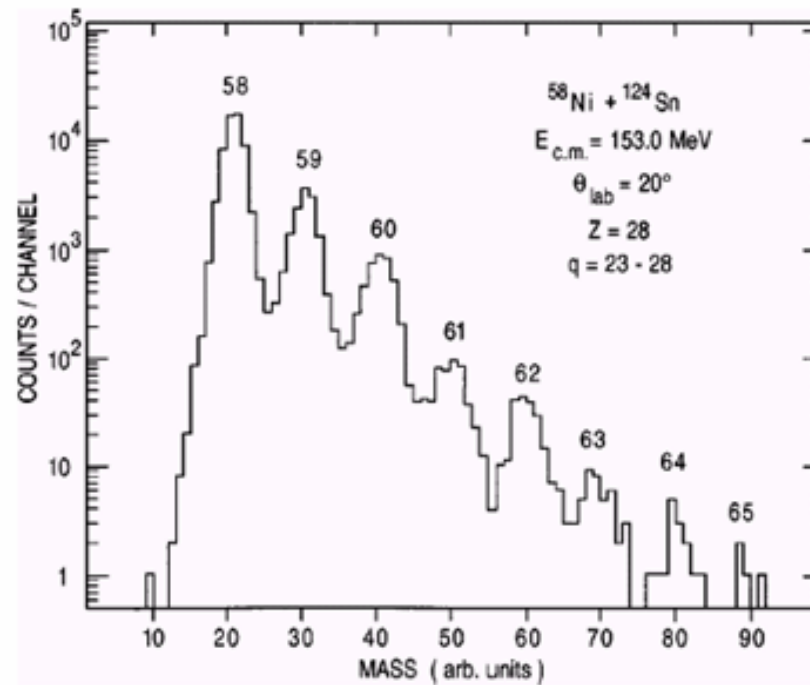
Transition Probabilities



Heavy Ion Transfer

A Few Words on Heavy Ion Transfer Reactions

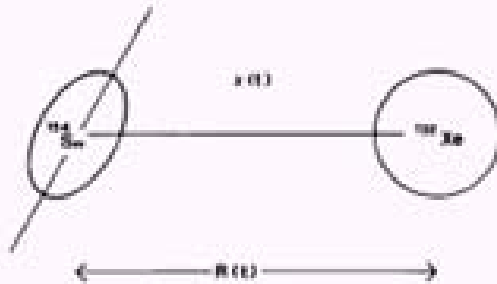
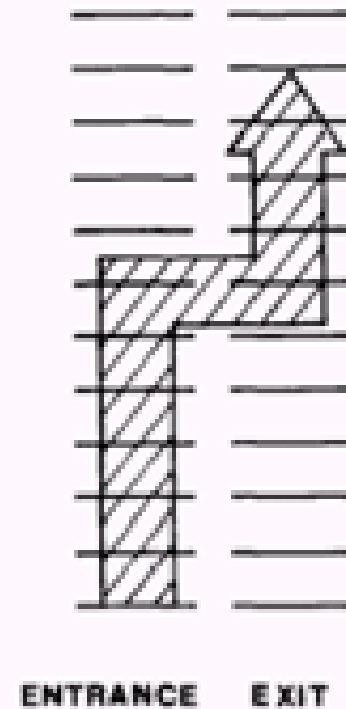
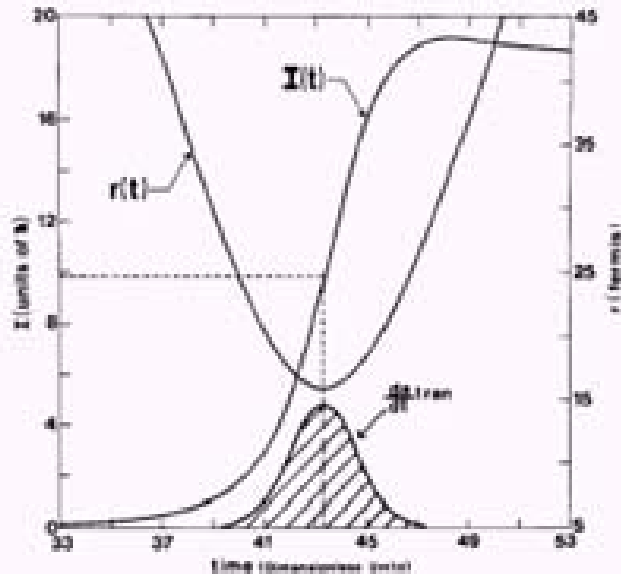
Multiple nucleon transfer



C.L.Jiang et al. PRC57(1998)

$$\Delta M/M, \Delta L/L \text{ and } \Delta E/E \ll 1$$

Large Sommerfeld parameter $\eta \gg 1$ allows for a semi-classical description. \Rightarrow Classical trajectory and tunneling



Population of high angular momentum and high excitation energy states

⇒ population mechanism, followed by γ -ray spectroscopy

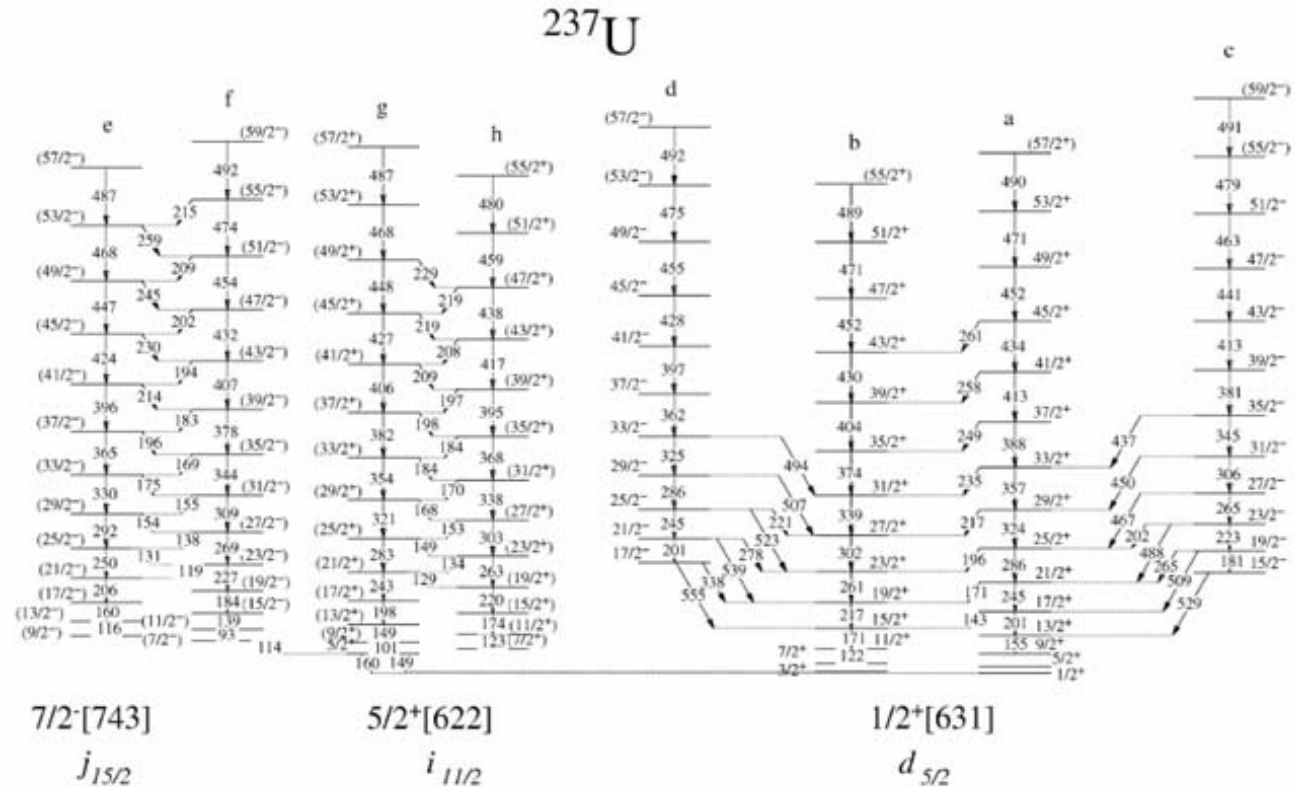


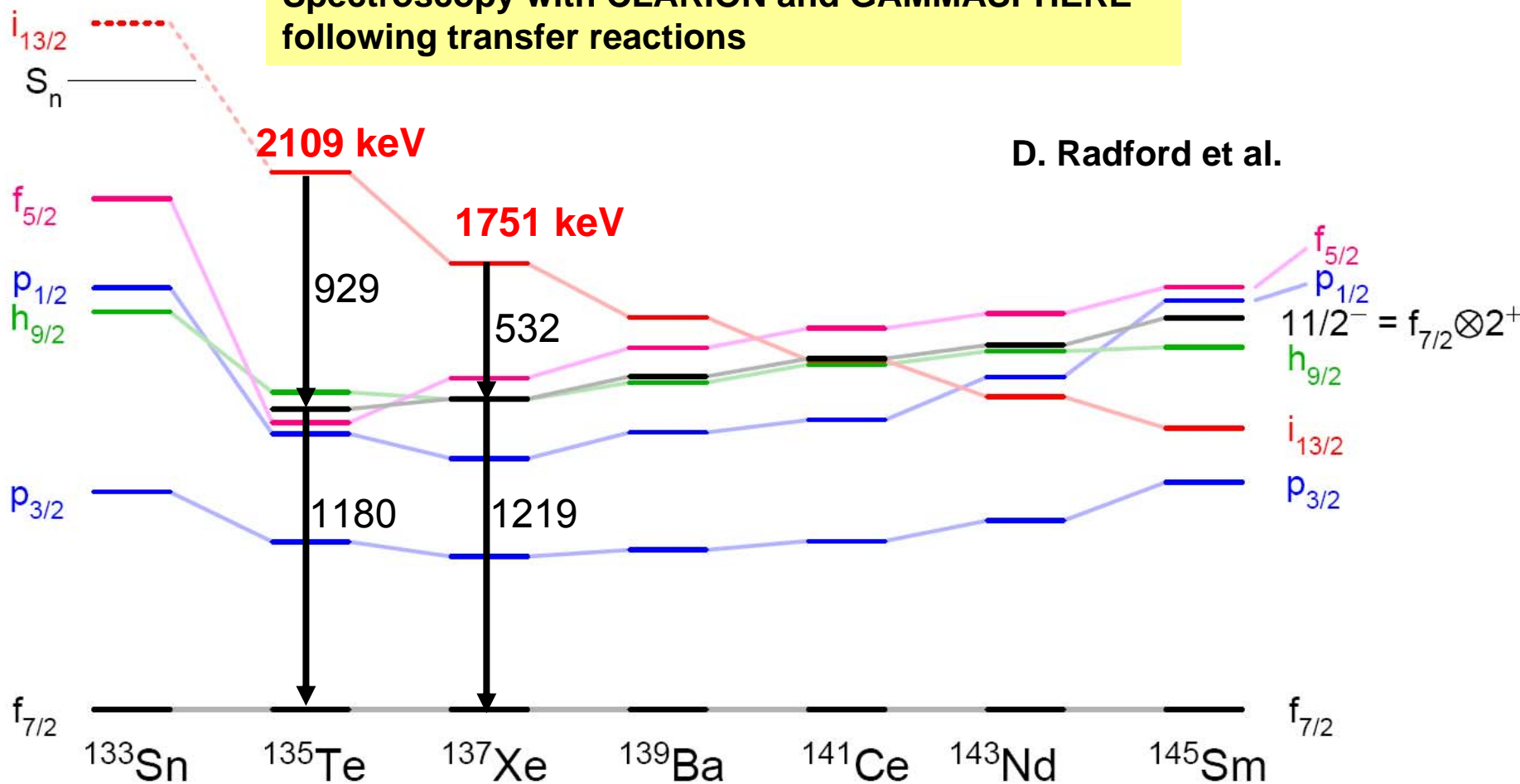
Fig. 3. Proposed level scheme for ^{237}U . The energies of the transitions are given in keV. Bands c and d are the octupole bands under discussion in the present Letter.

One neutron pickup reaction, $^{207}\text{Pb} + ^{238}\text{U}$ in GAMMASPHERE
to study octupole correlations in actinides

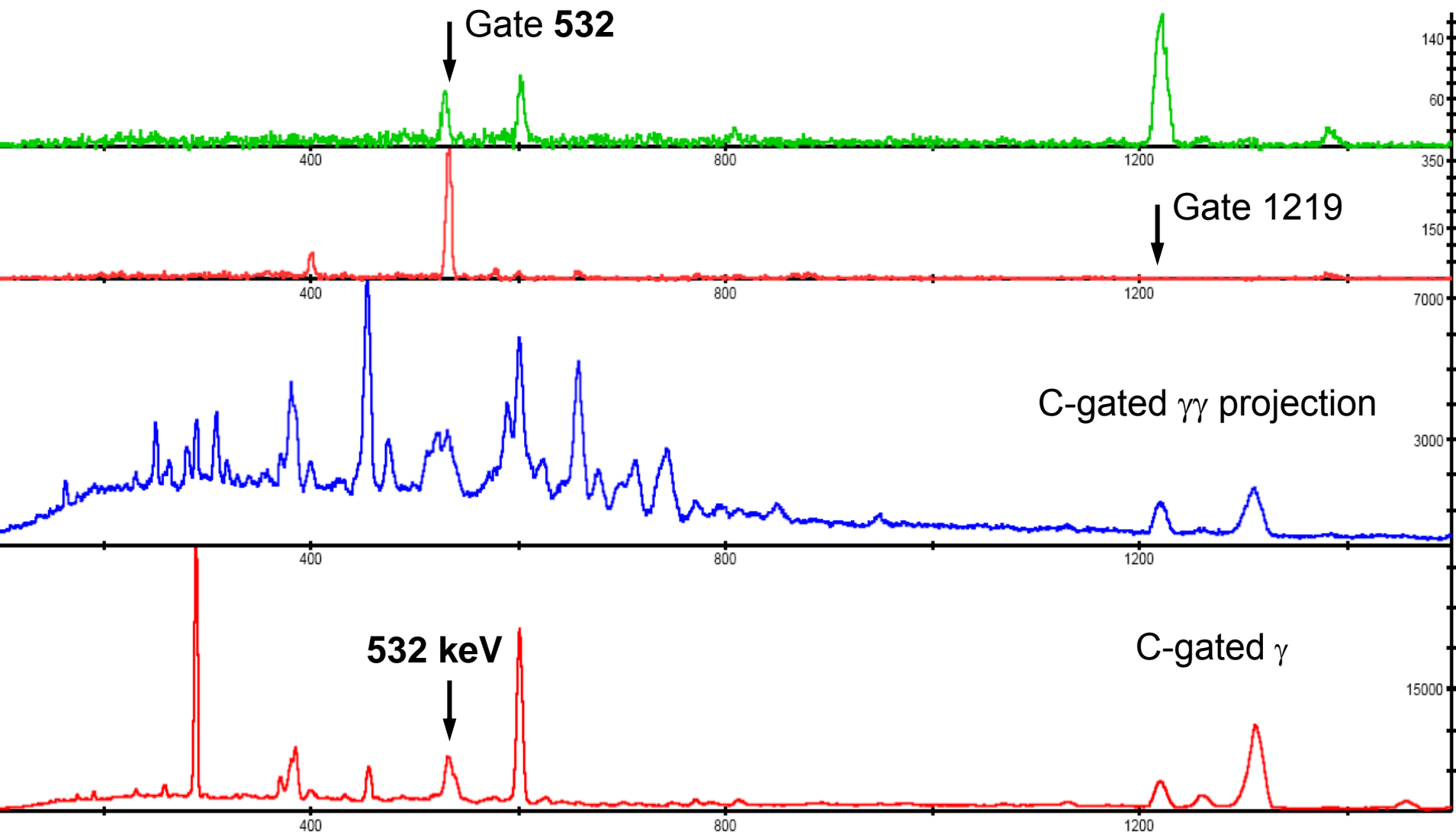
Spin-orbit splitting again

N = 83 level energy systematics

Spectroscopy with CLARION and GAMMASPHERE following transfer reactions



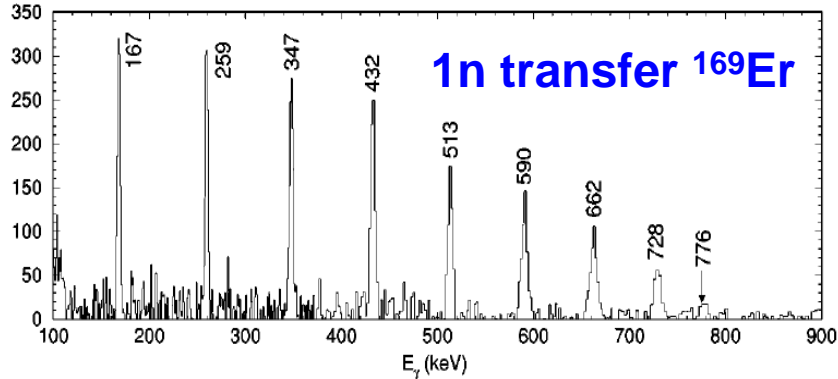
$^{13}\text{C}(^{136}\text{Xe}, ^{12}\text{C})^{137}\text{Xe}$ 560 MeV



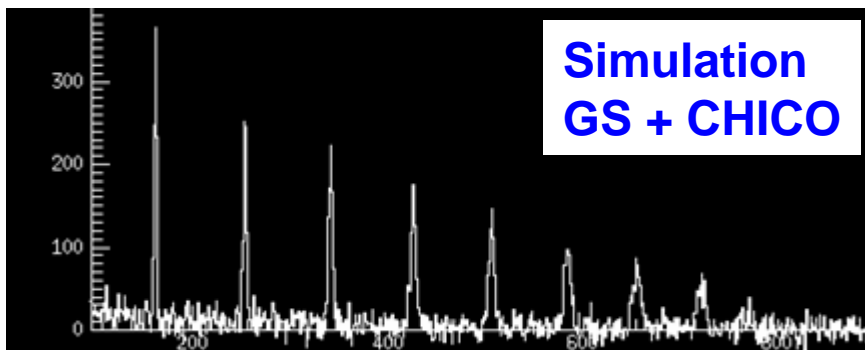
GRETINA: 1π array

Transfer reactions with re-accelerated RIBS and stable beams

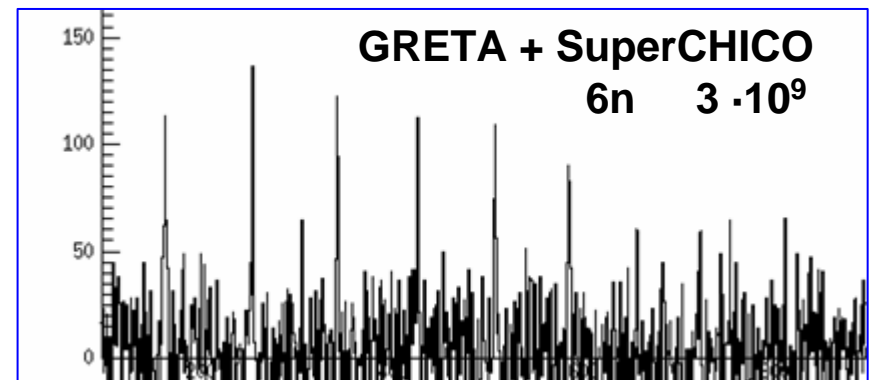
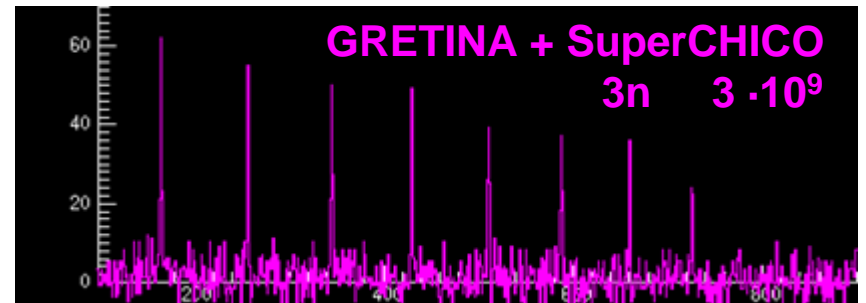
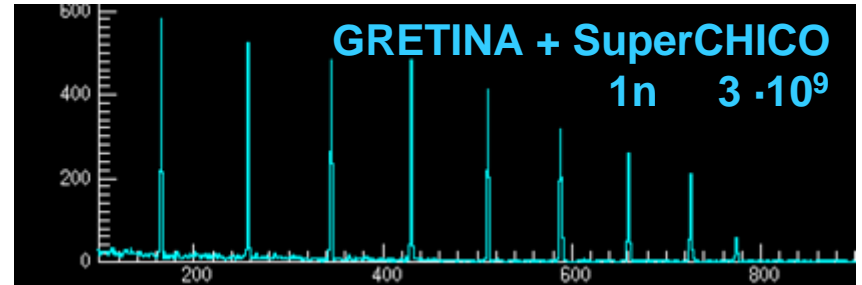
$^{238}\text{U} + ^{170}\text{Er}$ 5.7 MeV/u
GS + CHICO
 $3 \cdot 10^9$ p/s (0.5 pna), 0.5 mg/cm²
3 days, γ - γ - γ



C.Y. Wu et al., PRC 70, 014313 (2004)



Simulation $^{170}\text{Er} + ^{238}\text{U}$ 5.7 MeV/u



Conclusions

Transfer Reactions have provided a wealth of information that has shaped our current understanding of the structure of atomic nuclei. They will continue to provide a unique tool as we embark in our experimental study of very-neutron (proton) rich nuclei.

Existing and planned exotic beam facilities worldwide and new detector systems with increased sensitivity and resolving power not only will allow us to answer some burning questions we have today, but most likely will open up a window to new and unexpected phenomena.

As so eloquently expressed in the title of the conference, the nucleus and its structure are of paramount importance to many aspects of physics.

I believe the road ahead looks inviting!

Thank You!