

2014 ATLAS User's Meeting

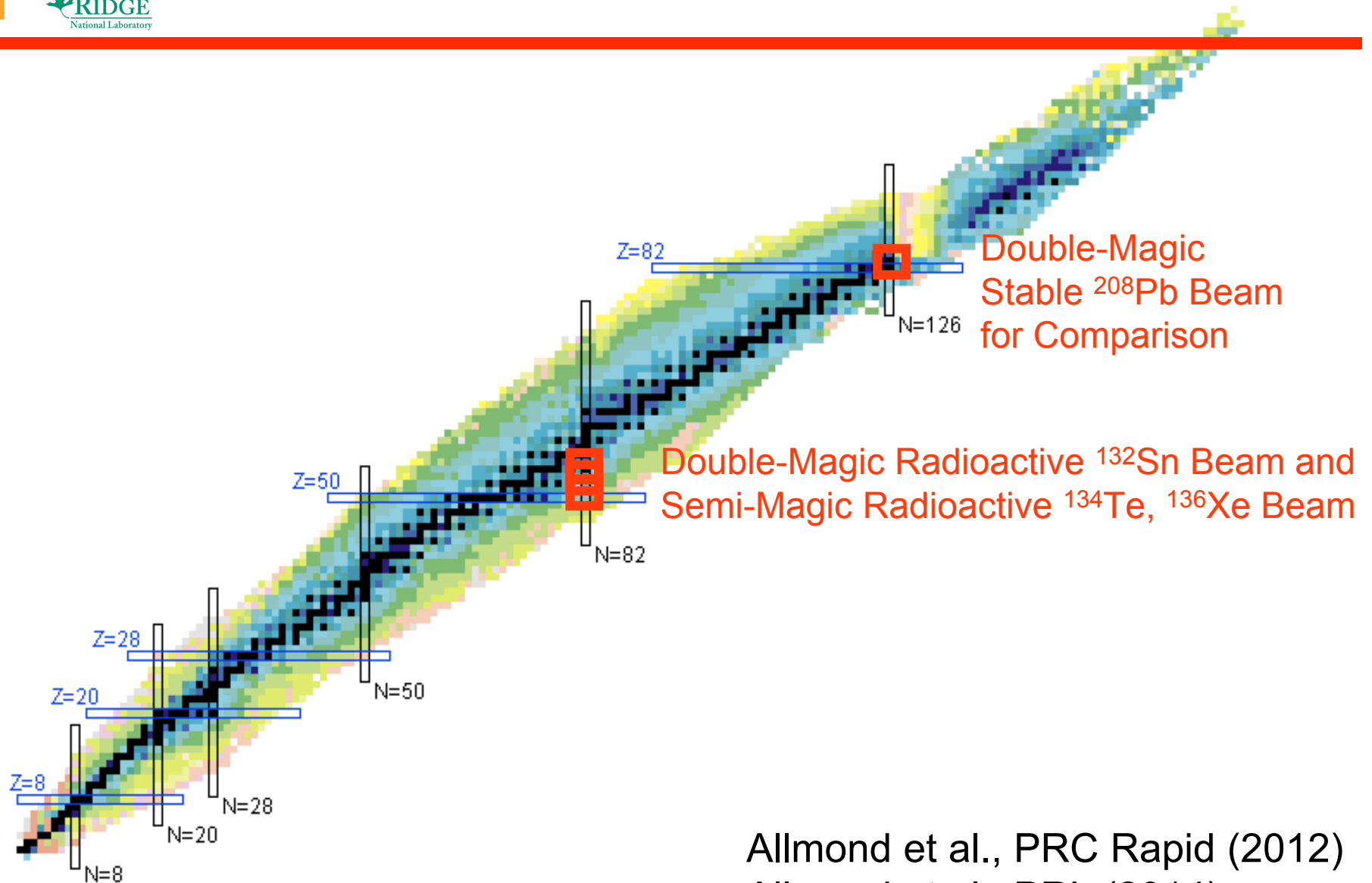
Heavy-Ion Induced Transfer Reactions using Particle- γ Coincidence Spectroscopy (Sub Coulomb)



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Example One-Neutron Transfer Studies



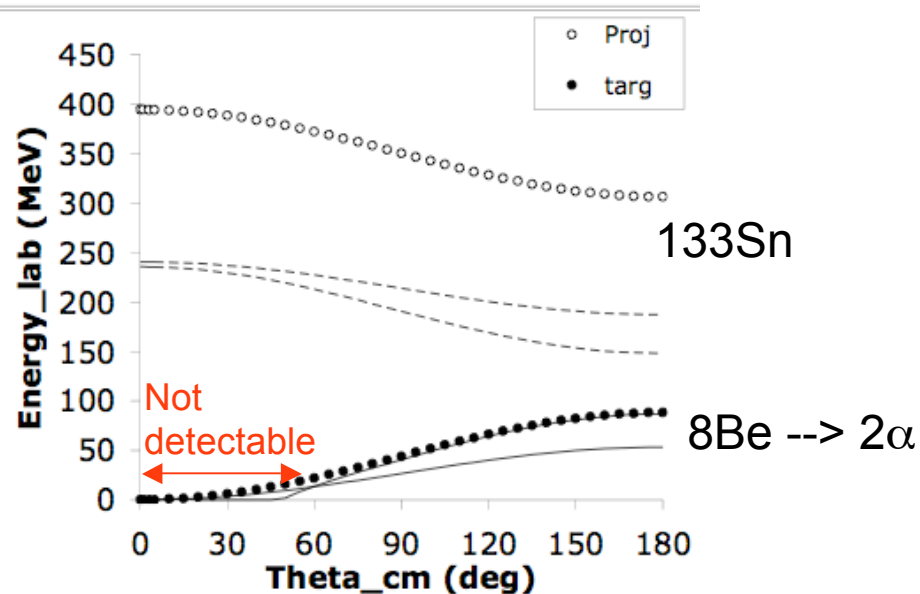
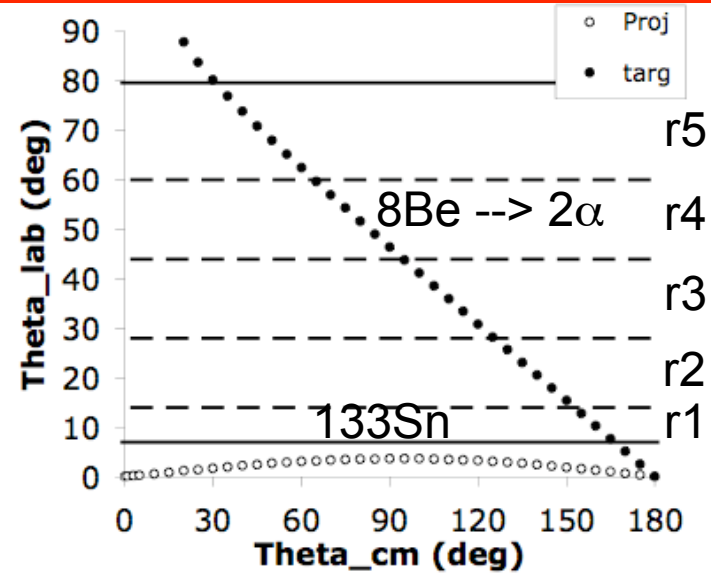
Allmond et al., PRC Rapid (2012)
Allmond et al., PRL (2014)

Inverse Kinematics: (${}^9\text{Be}$, ${}^8\text{Be} \rightarrow 2\alpha$)

Energetic/detectable target-like recoils predominately at backward θ_{cm}

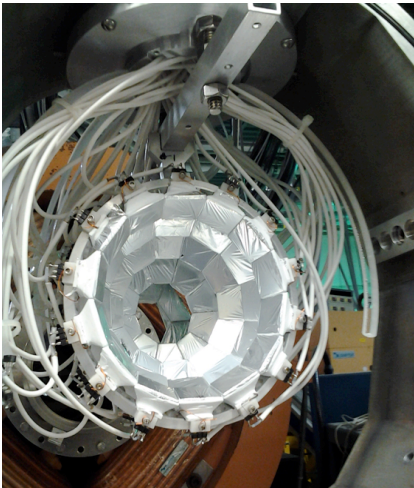
*Only need 2π particle detector

*Use sub-Coulomb to obtain reliable absolute cross-section normalization via Rutherford



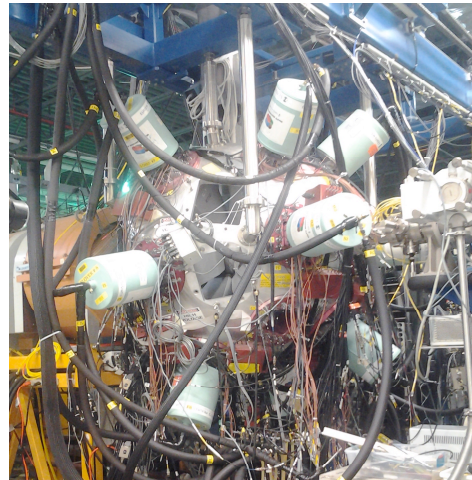
Hardware

CsI or Si*



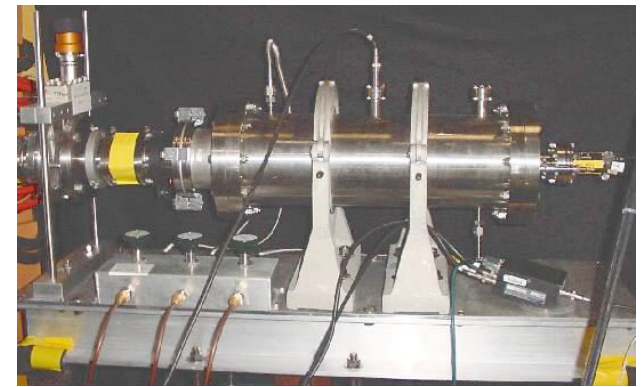
Particles

HPGe



γ -Rays

Zero-Degree Bragg



Target Thickness and
Beam Composition

*Recoiling target-like Heavy Ion may not make it through ΔE of Si telescope for PID: use (9Be, 8Be \rightarrow 2a) for clean tag in this case.

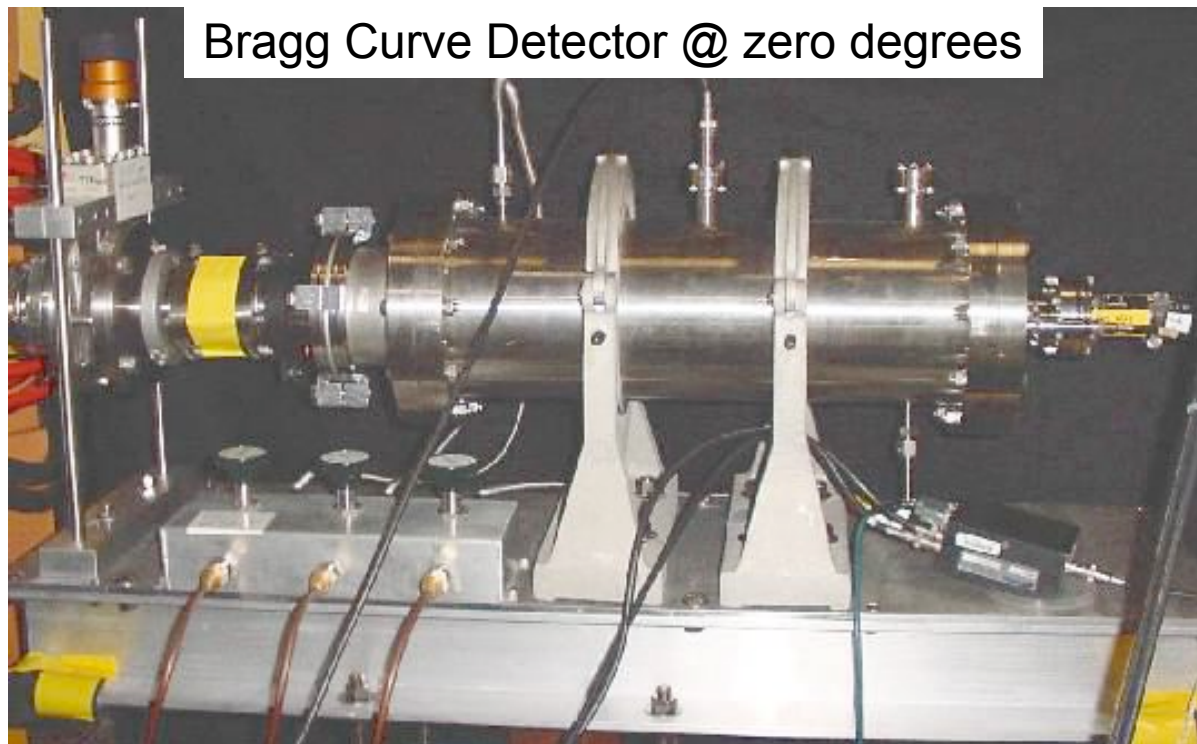
Measure Target Thickness / Eloss

Stopping powers are not known to high precision

$\frac{\sigma_{\text{exp}}}{\sigma_{\text{thy}}}$ and DSAM are sensitive to the target thickness and E_{loss}

WARNING!!!

Do not trust the energy loss calculated from the “nominal” target thickness.



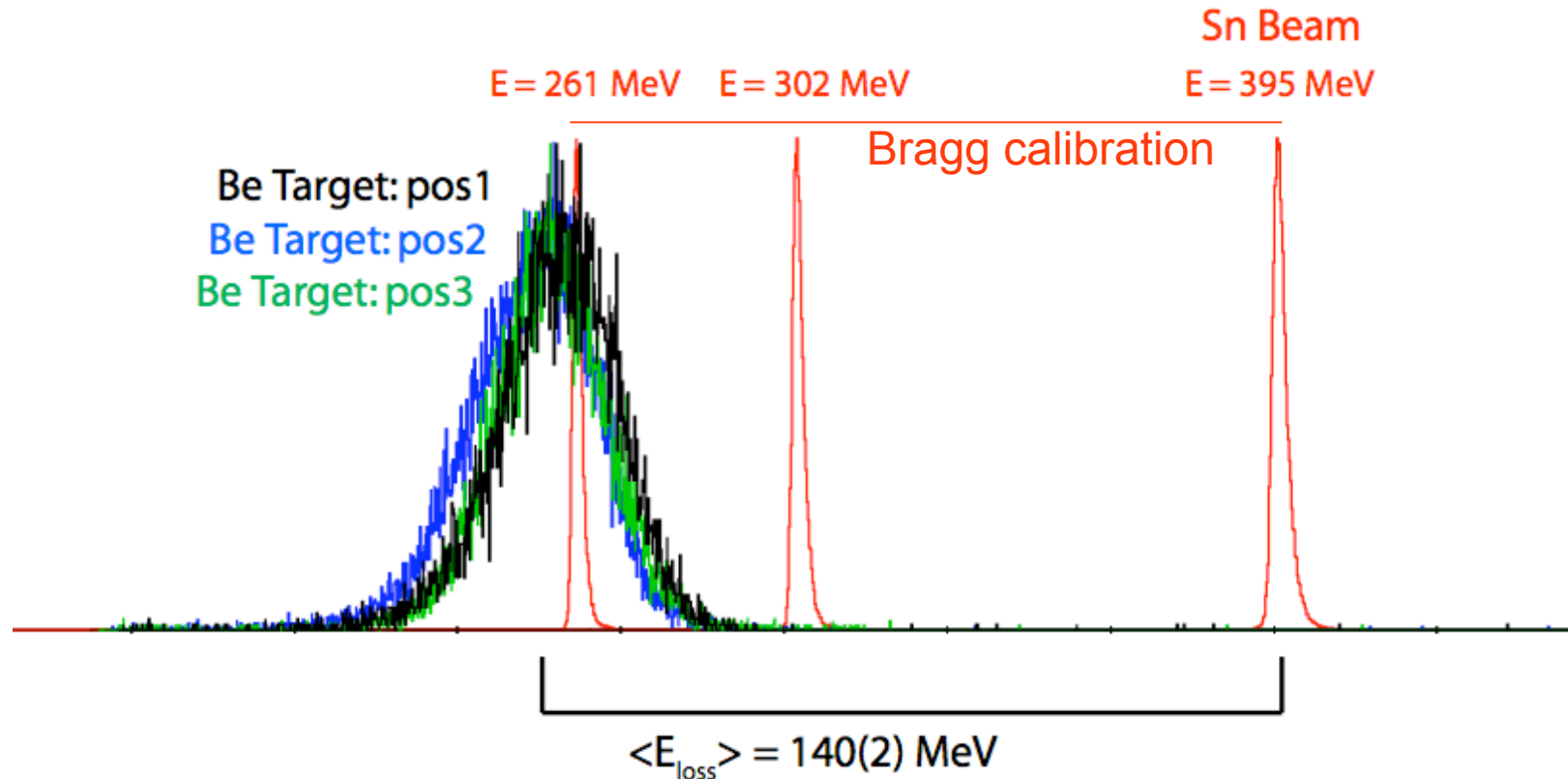
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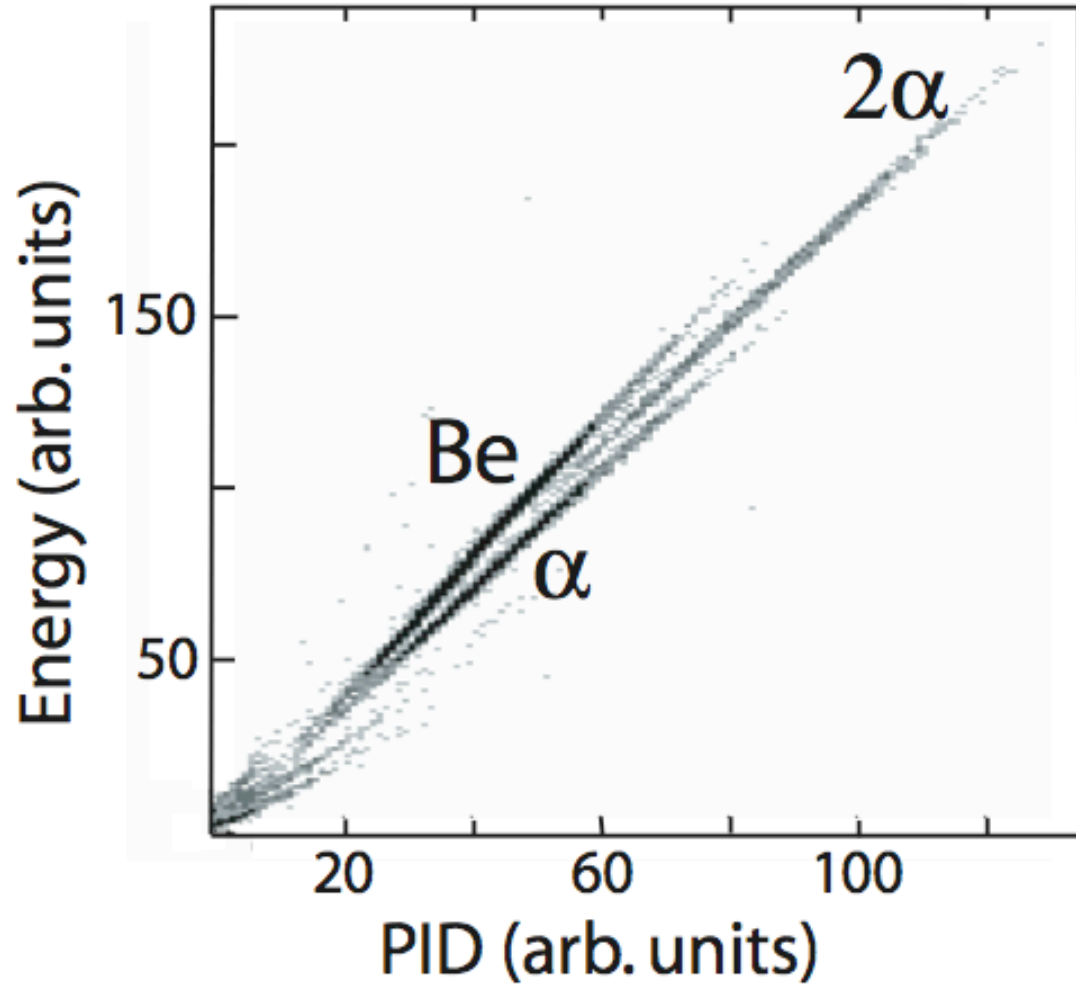
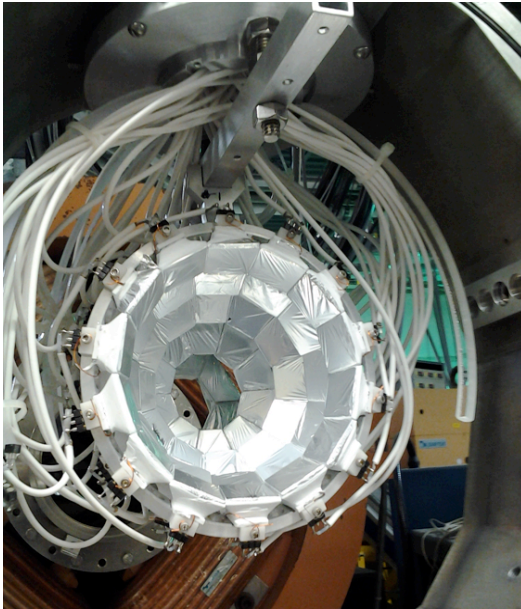


h r i b f



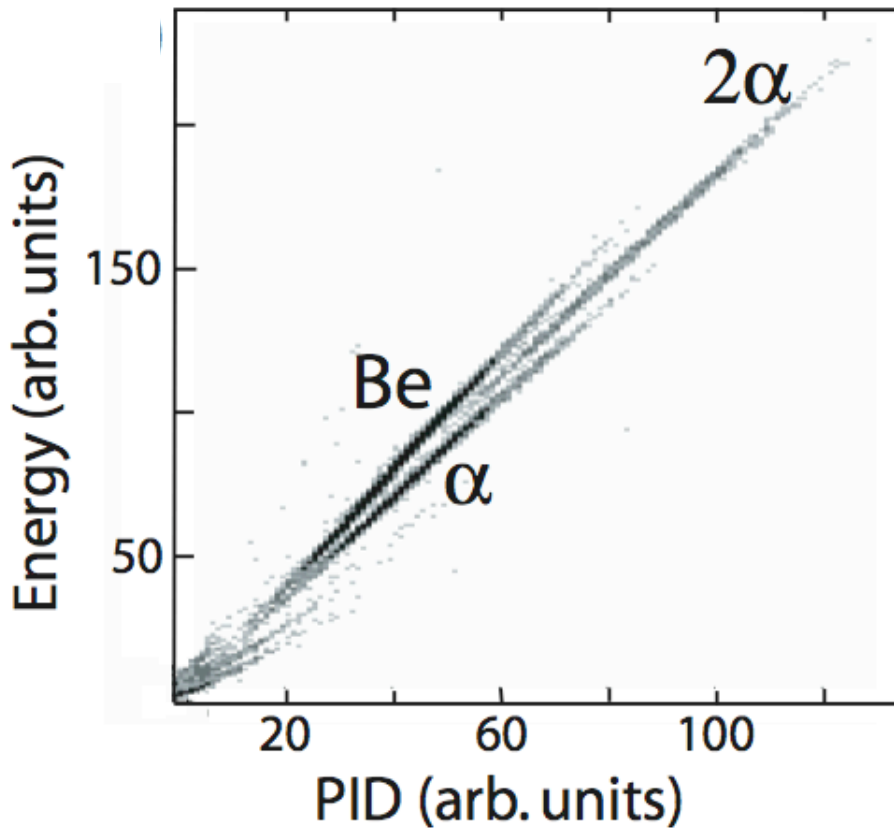
PID with CsI: (${}^9\text{Be}$, ${}^8\text{Be} \rightarrow 2\alpha$)

${}^8\text{Be}$ ($T_{1/2} = 8.2 \times 10^{-17}\text{s}$) decay to two correlated alphas provides clean trigger

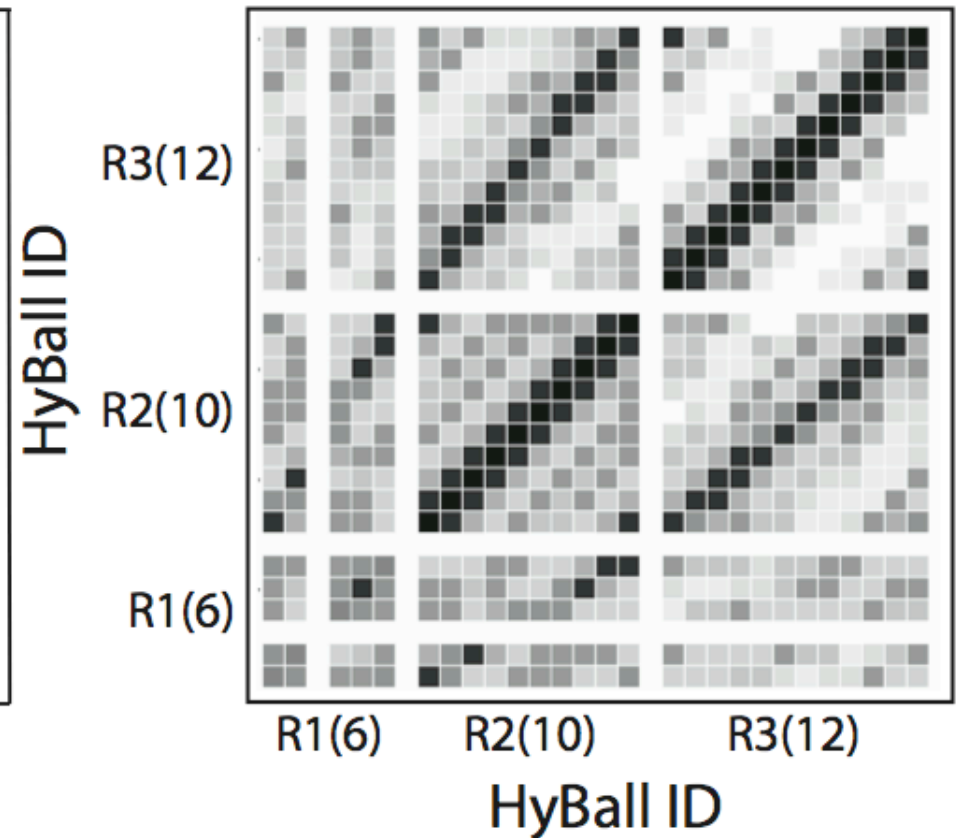


PID with CsI: (^9Be , $^8\text{Be} \rightarrow 2\alpha$)

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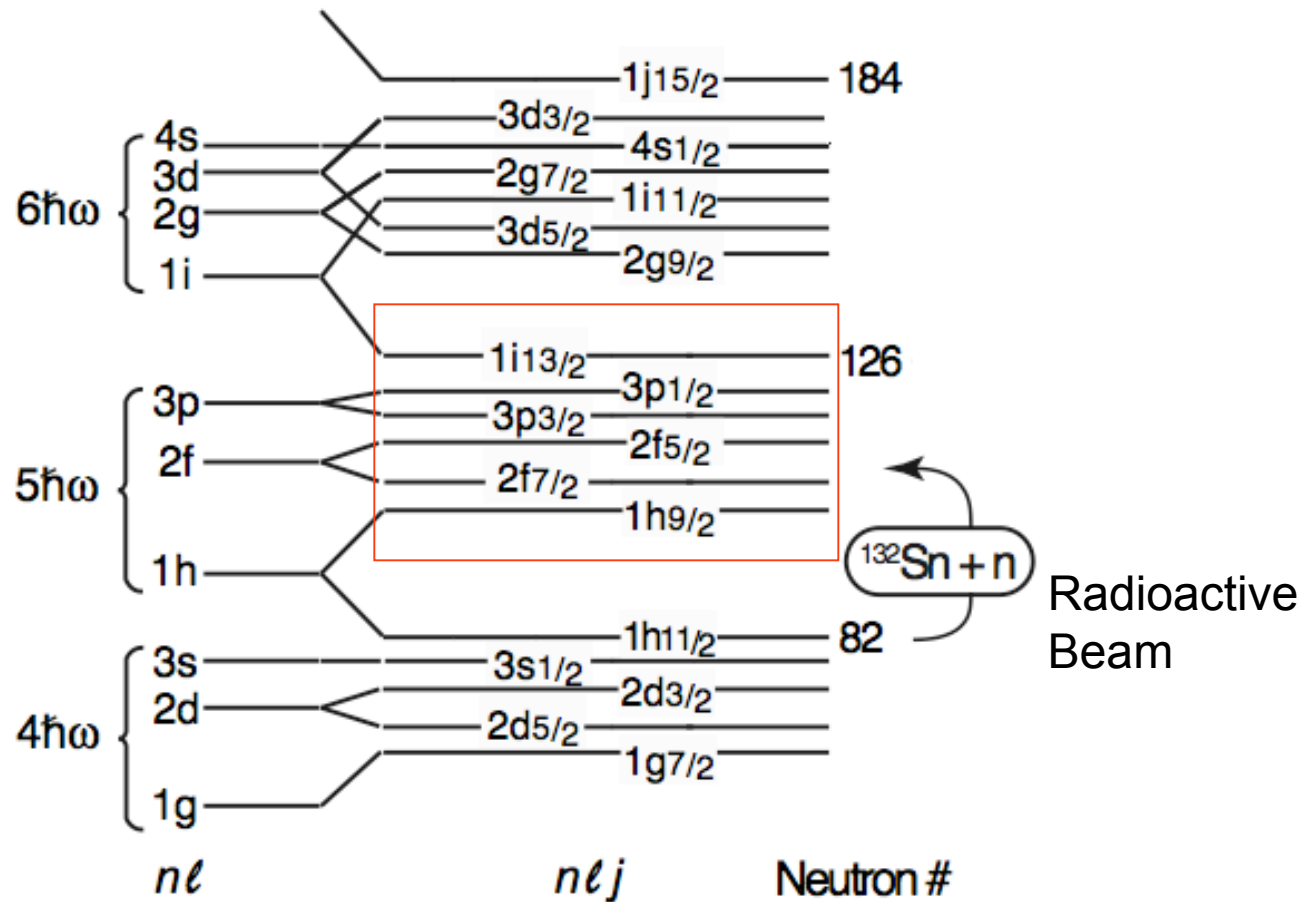
$^8\text{Be} \rightarrow 2 \times 1\alpha$ correlation
between CsI detector segments



*For Si detectors, use 2α hit in ΔE with equivalent energies for clean tag.

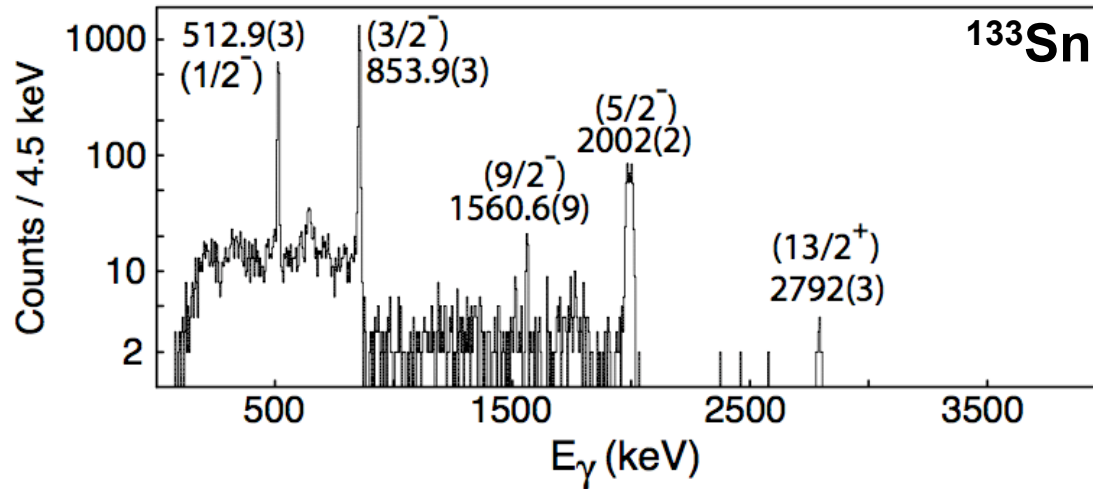
N=83 Single-Particle States

One-neutron transfer should select s.p. states above N=82 shell closure



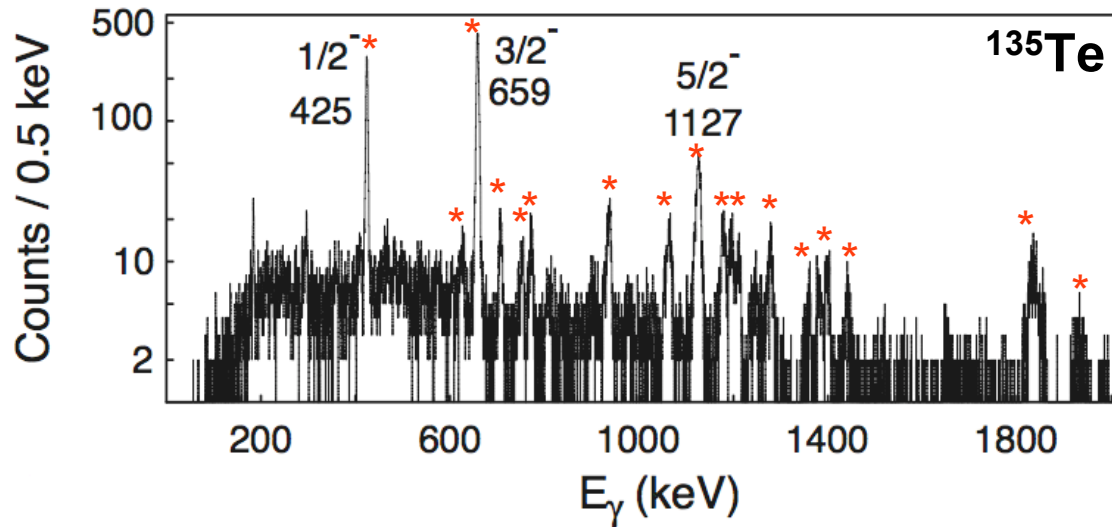
^{133}Sn versus $^{133}\text{Sn} + 2 \text{ Protons}$

Going from system of 1n to 1n+2p adds a lot of complexity



1 Active Neutron
"Simple"

Pure s.p. states

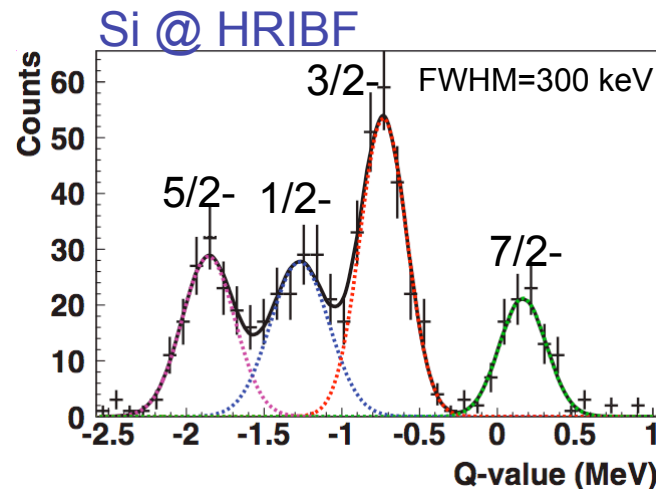


1 Active Neutron
2 Active Protons
"Complex"

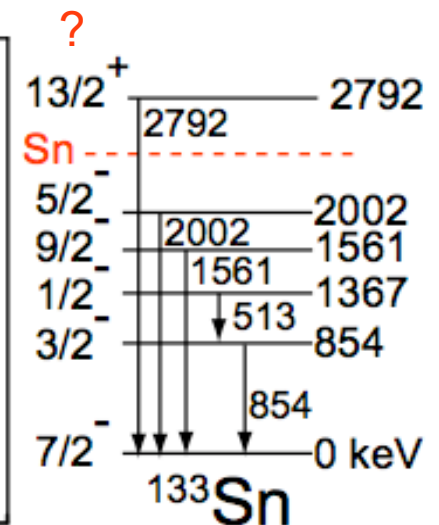
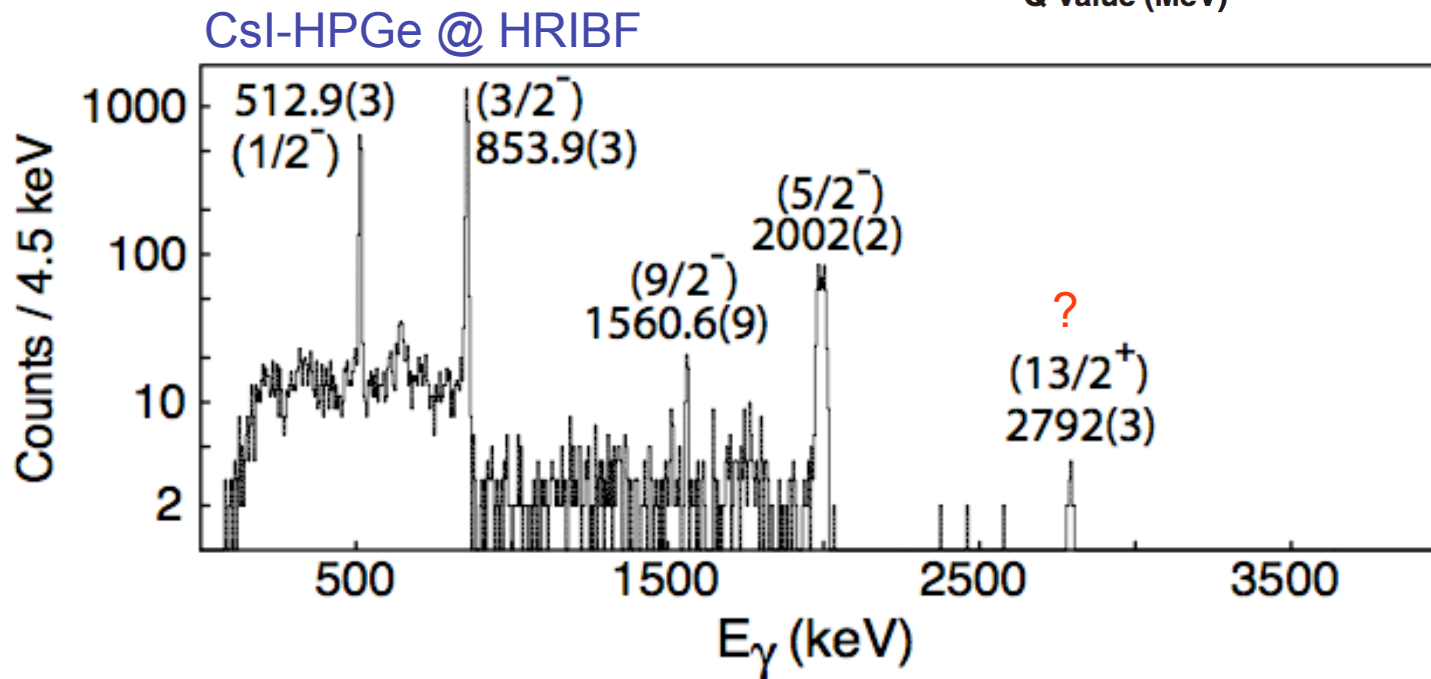
Fragmented s.p.
states

Particle vs Particle- γ Spectra: ^{133}Sn

Transitions / states are well resolved in inverse kinematics by particle- γ

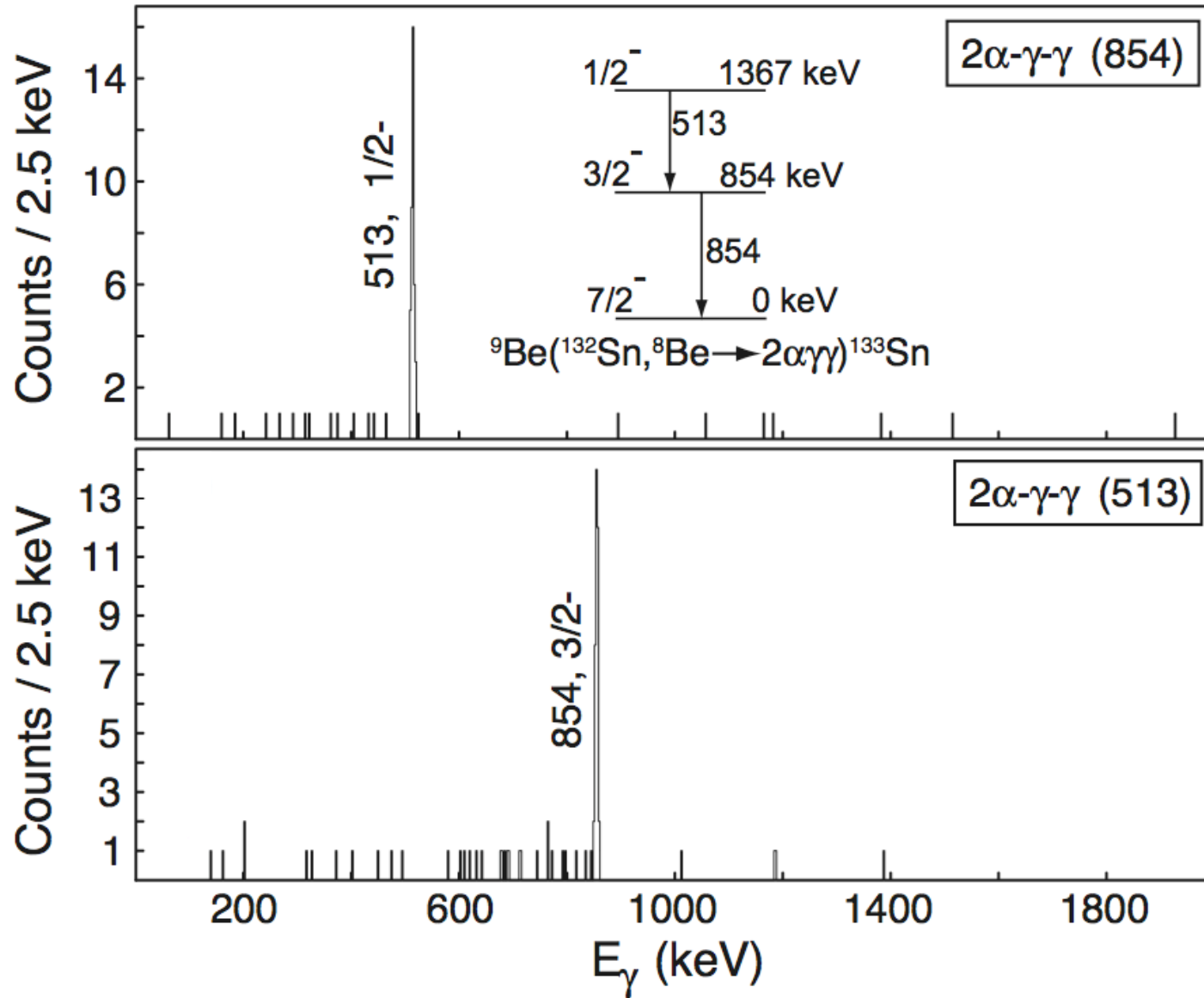


$d(^{132}\text{Sn},p)^{133}\text{Sn}$
 K. L. Jones et al.,
 Nature (London)
 465, 454 (2010).



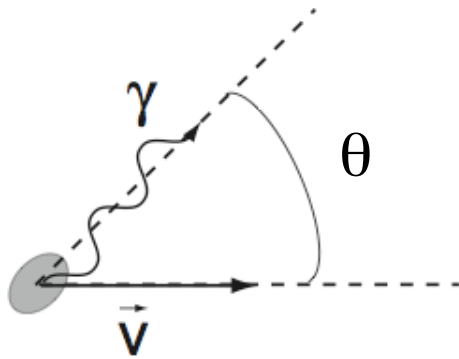
$^{133}\text{Sn}(N=83)$ Decay Paths by $\gamma\text{-}\gamma$

$\gamma\text{-}\gamma$ coincidences can be used to determine decay paths



Lifetimes by Doppler Shift

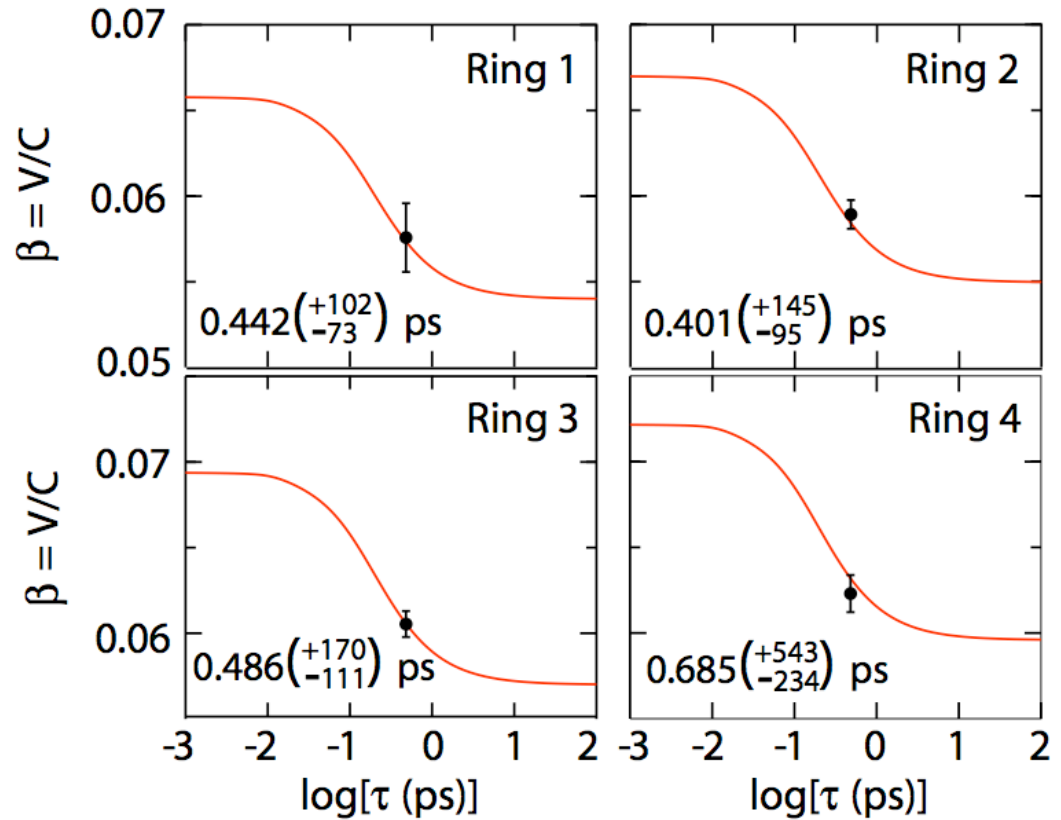
Lifetimes can be measured if comparable to flight time through target



$$E_{\gamma} = E_0 \frac{\sqrt{1 - \beta^2}}{1 - \beta \cos \theta}$$

Experimentally Measure

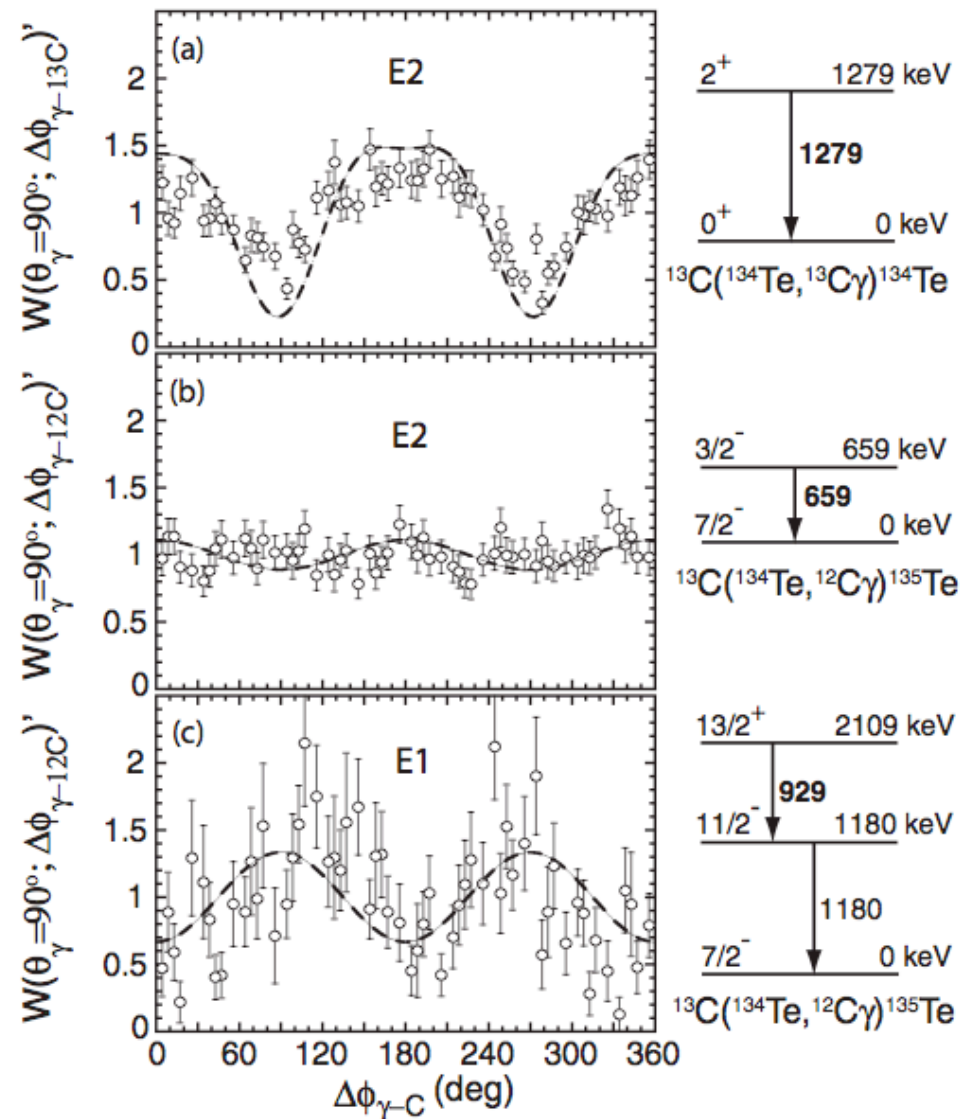
1/2- @ 1363 keV



$\tau = 0.48(13)$ ps

Particle- γ Angular Correlations

Can use particle- γ correlations to determine multipolarity of transitions

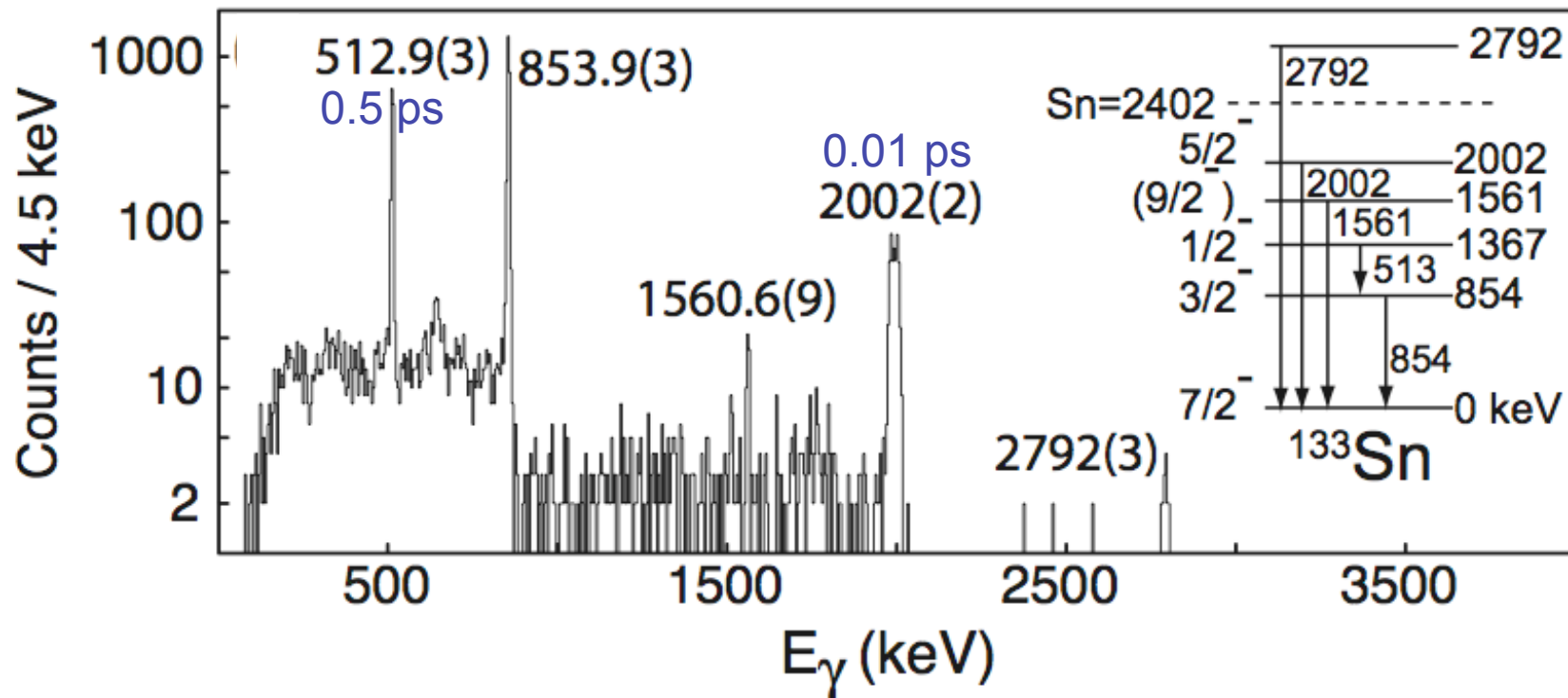


$(^9\text{Be}, ^8\text{Be}\gamma)^{133}\text{Sn}$ Summary

Extensive spectroscopic information determined

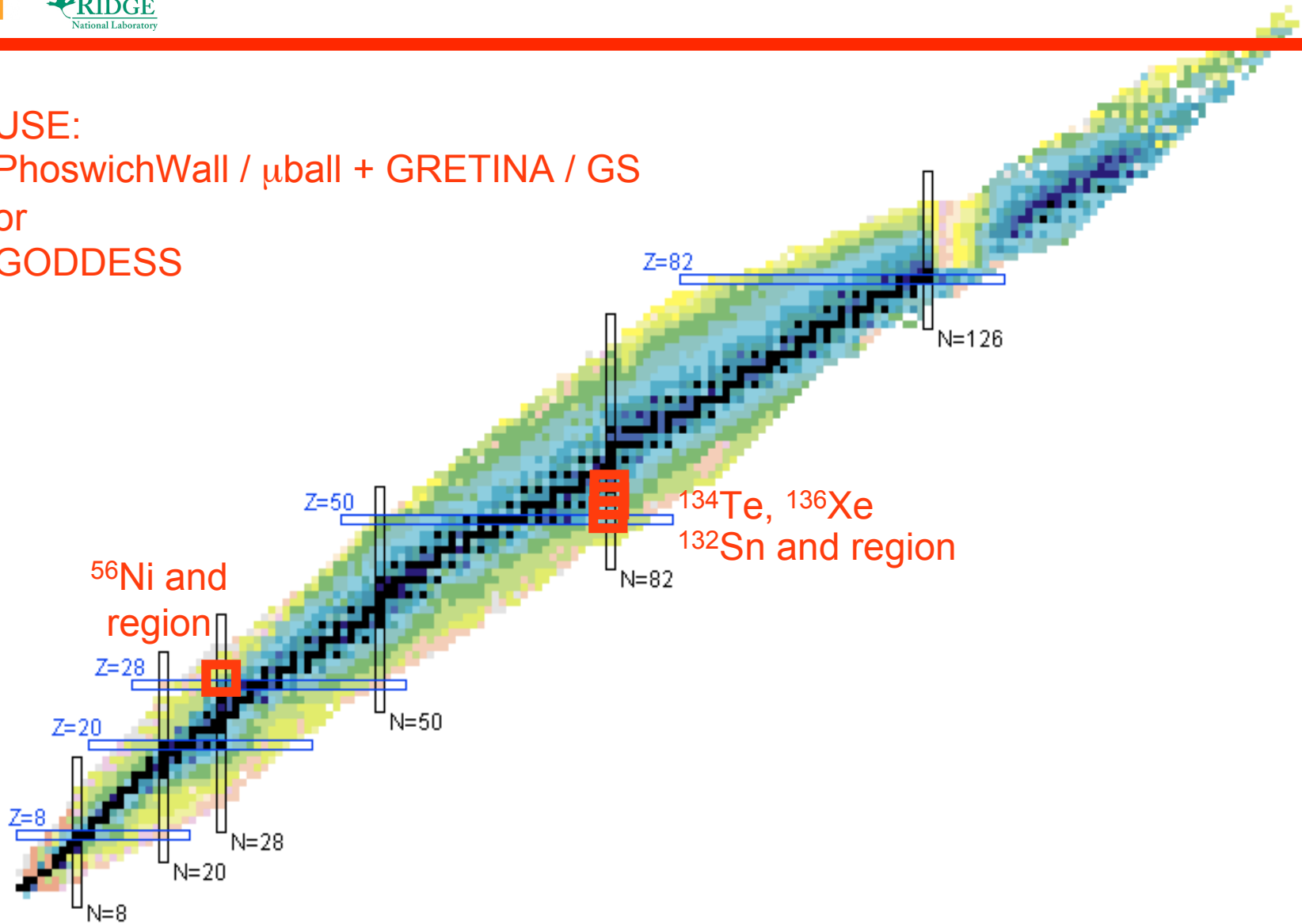
Complete Spectroscopy by Particle- γ

- J^π -- p- γ , γ paths, and target $l \pm 1/2$ selectivity
- $E(J^\pi)$ -- γ -ray Energies
- $\tau(J^\pi)$ -- Doppler Shift
- $\sigma(J^\pi)$ -- γ -ray Intensities
- Spectroscopic Factors (questionable)
- Asymptotic Normalization Coefs. (reliable)



Initial Cases for CARIBU / AIRIS

USE:
PhoswichWall / μ ball + GRETINA / GS
or
GODDESS





Thanks to all of the Collaborators

**D.C. Radford², A. Galindo-Uribarri^{2,3}, A.E. Stuchbery⁴, J.R. Beene², R.L. Varner², E. Padilla-Rodal⁵,
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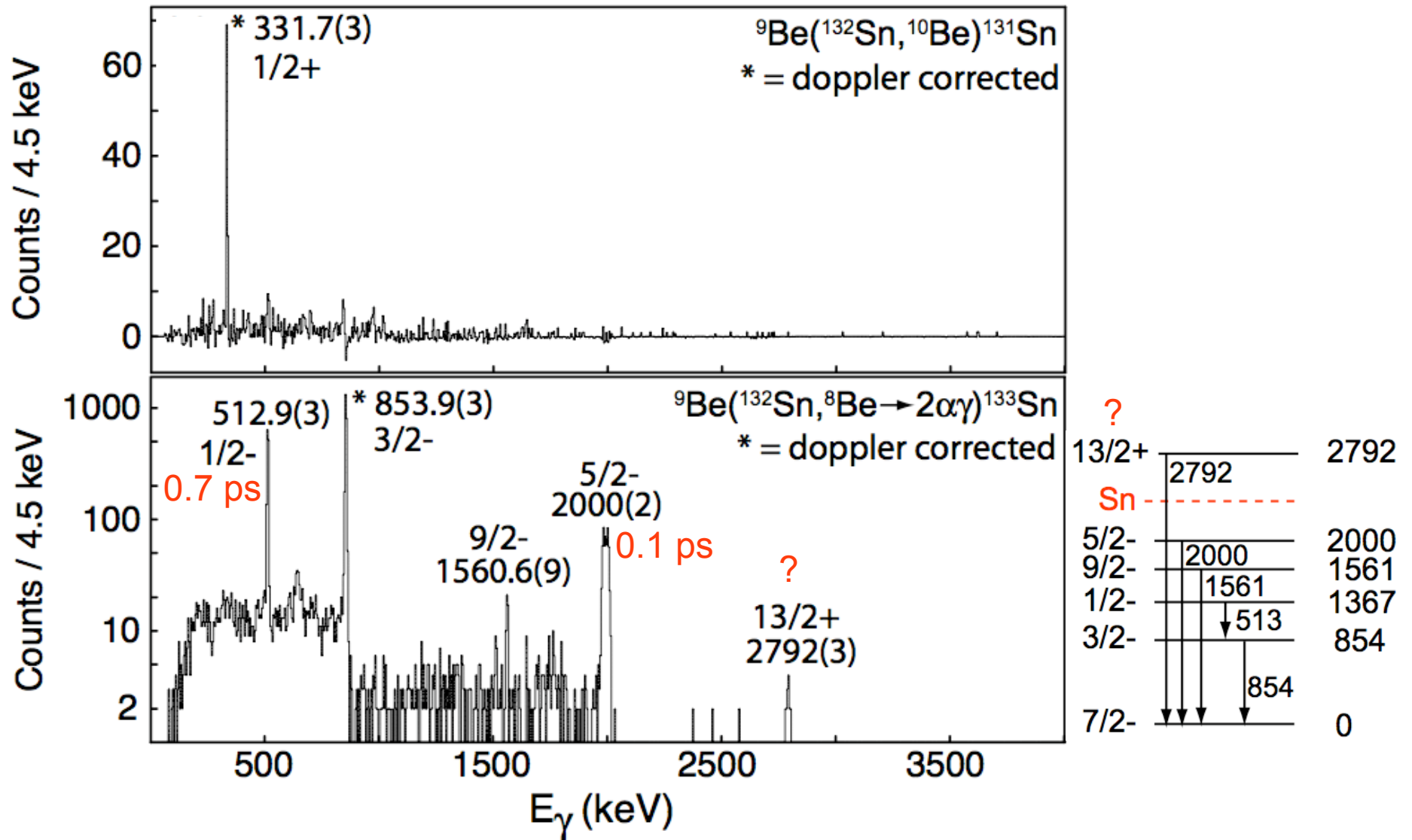


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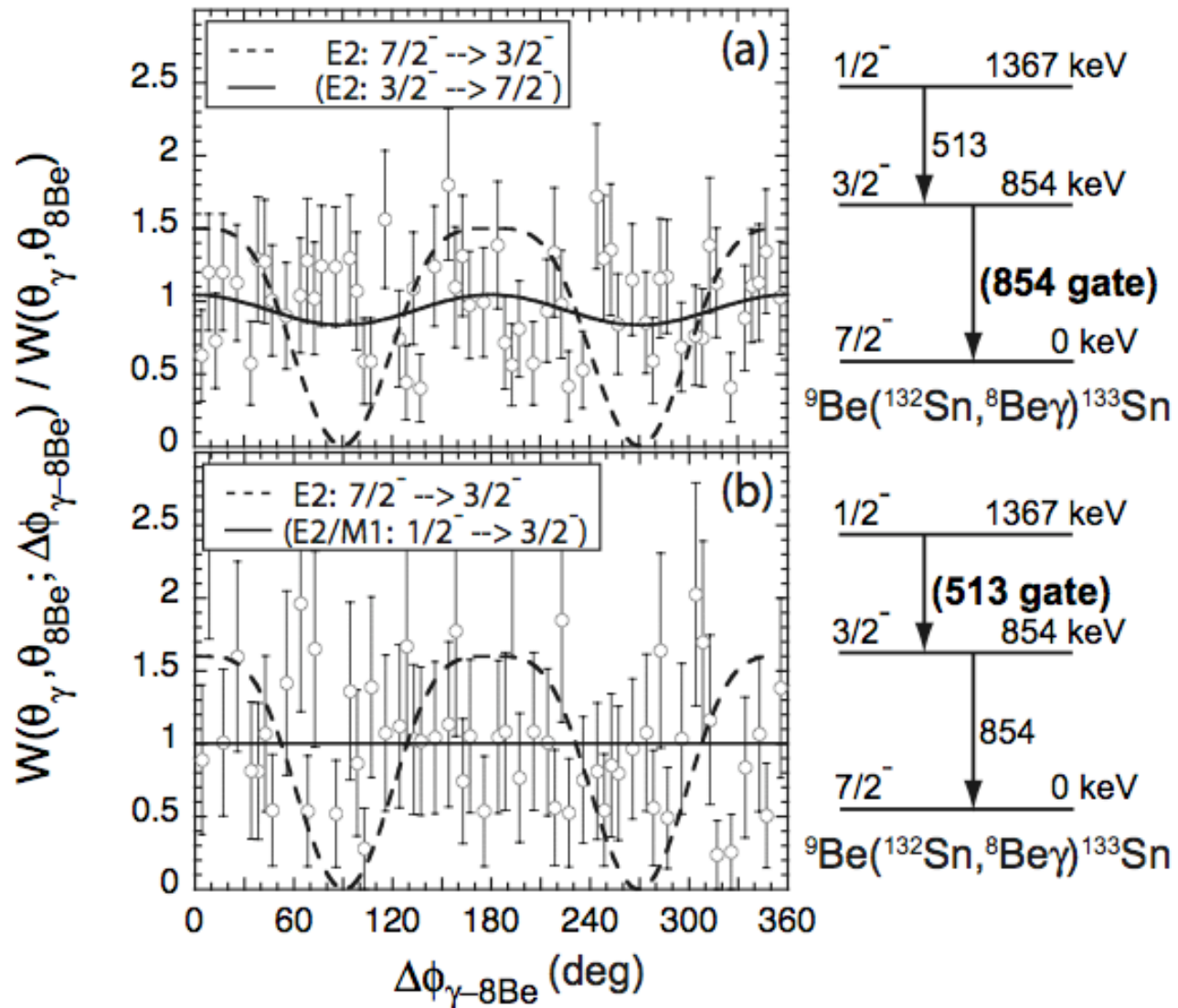
$(^9\text{Be}, ^{10}\text{Be}\gamma)^{131}\text{Sn}$ and $(^9\text{Be}, ^8\text{Be}\gamma)^{133}\text{Sn}$

All expected N=83 single-particle states accounted for in ^{133}Sn



Particle- γ Angular Correlations in ^{133}Sn

The two transitions are consistent with low-spin initial states



Results

E_x (keV)	J^π	E_γ (keV)	τ (fs)	σ (mb)	σ_{thy} (mb)	Present (^9Be , $^{8,10}\text{Be}$) S	[25,51] (d,p) S	Present (^9Be , $^{8,10}\text{Be}$) C^2 (fm $^{-1}$)	[25,51] (d, p) C^2 (fm $^{-1}$)	[53] (^{13}C , ^{12}C) C^2 (fm $^{-1}$)
^9Be (^{208}Pb , ^8Be) ^{209}Pb										
0	$9/2^+$				0.0013(4)		1.21(36)		2.20(17)	2.25(29)
778.9(3) ^b	$11/2^+$				0.0005(2)		1.57(47)		0.00187(13)	0.0037(5)
1423(1) ^b	$15/2^-$				0.0001(1)		1.19(36)		$2.5(2) \times 10^{-5}$	
1566.0(9)	$5/2^+$	1566.0(9)		0.13(4)	0.084(21)	1.5(6)	1.08(32)	14(5)	13.0(7)	
2031(1)	$1/2^+$	464.5(4)		0.28(2)	0.22(5)	1.3(3)	1.04(31)	45(8)	48.7(30)	41.7(54)
2489(2)	$7/2^+$	2489(2)		0.10(2)	0.062(19)	1.6(6)	1.27(38)	0.026(6)	0.025(2)	
2535(1)	$3/2^+$	969.4(6)	87(24)	0.43(3)	0.38(9)	1.1(3)	1.11(33)	2.3(4)	2.93(20)	
^9Be (^{132}Sn , ^{10}Be) ^{131}Sn										
0	$(3/2^+)$				0.15(11)					
331.7(3)	$(1/2^+)$	331.7(3)		0.68(8)	0.17(12)	4(3)				
1654.53(8) ^b	$(5/2^+)$				0.03(2)					
^9Be (^{132}Sn , ^8Be) ^{133}Sn										
0	$7/2^-$				3(1)		0.86(7)		0.64(5)	
853.9(3)	$3/2^-$	853.9(3)		12(1)	13(3)	0.9(2)	0.92(7)	6.0(14)	5.6(4)	
1366.8(4)	$1/2^-$	512.9(3)	$480^{(+160)}_{(-100)}$	11(1)	12(3)	0.9(2)	1.1(2)	2.5(5)	2.6(6)	
1560.6(9)	$(9/2^-)$	1560.6(9)		0.58(10)	1.1(4)	0.5(2)		$5.1(15) \times 10^{-6}$		
2002(2)	$5/2^-$	2002(2)	$13^{(+10)}_{(-13)}$	8.6(6)	9.6(24)	0.9(2)	1.1(2)	0.0020(4)	0.0009(2)	
2792(3)		2792(3)		0.38(9)	0.18(7) ^c					

Results

Nuclide	Transition	$B(M1)^{exp}$	$B(M1)^{thy}$
^{209}Pb	$3d_{3/2} \rightarrow 3d_{5/2}$	0.72(20)	0.71 ^b
^{207}Pb	$3p_{3/2}^{-1} \rightarrow 3p_{1/2}^{-1}$	0.47(6) ^a	0.40 ^b
^{133}Sn	$2f_{5/2} \rightarrow 2f_{7/2}$	0.55($^{+\infty}_{-14}$)	0.52 ^c
^{133}Sn	$3p_{1/2} \rightarrow 3p_{3/2}$	0.88($^{+23}_{-22}$)	0.67 ^c

Results

