The NJL-jet model
for quark fragmentation functions

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Introduction

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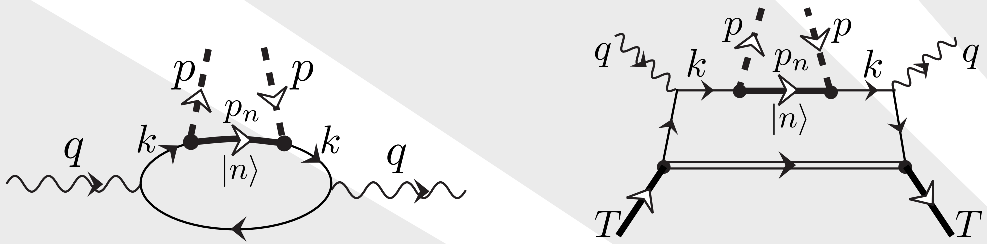
Fragmentation function $D_q^h(z)$ $(z = \frac{2p\cdot q}{q^2})$ describes **semi-inclusive hadron production** in e^+e^- annihilation and (e, e') DIS processes.
Parton model diagrams for cross sections:



- This simple picture of "**independent fragmentation**" was formulated by Field and Feynman (Phys. Rev. **D 15** (1977) 250). (Same as "factorization".)
- **Empirical fragmentation functions** were extracted from data. For example: M. Hirai et al: PRD **75** (2007) 094009.
- Model calculations using **effective quark theories**: Almost all calculations introduced artificial "normalization factors" (or other ad-hoc parameters) to enlarge the calculated fragmentation functions.

Definitions and interpretation

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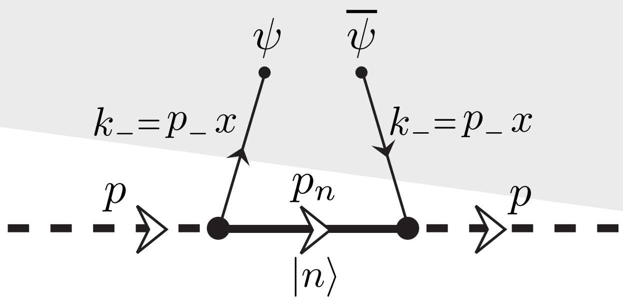
Compare definitions of **distributions** and **fragmentations**:

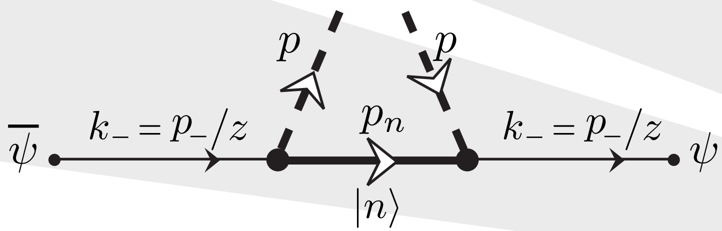
$$f_q^h(x) = \frac{1}{2}\sum_n \delta(p_- x - p_- + p_{n-})\langle p|\overline{\psi}|p_n\rangle\gamma^+\langle p_n|\psi|p\rangle \quad (\mathbf{p}_T = 0)$$

$$= p_- \int \mathrm{d}^2 k_T \sum_\alpha \frac{\langle p|b_\alpha^\dagger(k) b_\alpha(k)|p\rangle}{\langle p|p\rangle}, \quad (\mathbf{p}_T = 0): \ \textbf{quarks in hadron.}$$

$$D_q^h(z) = \frac{z}{12}\sum_n \delta\left(\frac{p_-}{z} - p_- - p_{n-}\right)\langle p, p_n|\overline{\psi}|0\rangle\gamma^+\langle 0|\psi|p, p_n\rangle \quad (\mathbf{p}_T = 0)$$

$$= \frac{k_-}{6}\int \mathrm{d}^2 p_\perp \sum_\alpha \frac{\langle k(\alpha)|a_h^\dagger(p) a_h(p)|k(\alpha)\rangle}{\langle k(\alpha)|k(\alpha)\rangle}, \quad (\mathbf{k}_\perp = 0): \ \textbf{hadrons in quark!}$$





Momentum sum rules

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- Interpretation of $D_q^h(z)$: **Probability that a hadron (p) in the cloud of a virtual quark (k) has fraction z of the quark's light cone momentum: $\mathbf{p_- = z k_-}$.**
- Formal relation between distribution and fragmentation (from crossing and charge conjugation): $D_q^h(z) = (-1)^{2(s_q+s_h)+1}\frac{z}{6}\, f_q^h(x = \frac{1}{z})$. However, in practice this "**Drell-Levy-Yan**" relation is (almost) useless.
- Momentum sum rules (sums include antiparticles):

$$\sum_q \int_0^1 x\, \mathrm{d}x\, f_q^h(x) = 1 : \textbf{hadron consists of quarks.}$$

$$\sum_h \int_0^1 z\, \mathrm{d}z\, D_q^h(z) = 1 : \textbf{quark hadronizes completely!}$$

Derivation of sum rule for D_q^h assumes that **quark is an eigenstate of the momentum operator** $\hat{P}_- = \sum_h \int_0^\infty \mathrm{d}p_- \int \mathrm{d}^2 p_\perp \left(p_- a_h^\dagger(p) a_h(p)\right)$ **expressed in terms of hadrons!**

Elementary NJL $q \to \pi$ fragmentation: $d_q^\pi(z)$

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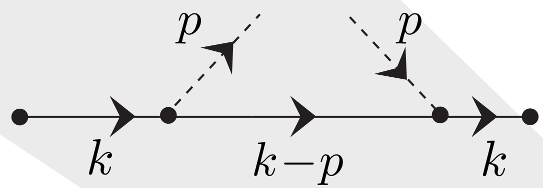
Simplest approximation: Truncate $|n\rangle$ to **one-quark state**:



- This approximation leads to **disastrous results**: Extermely small compared to the empirical functions. To avoid this, previous calculations using effective quark models introduced "normalization constants" or other ad-hoc parameters.
- The lowest order fragmentation process $q \to q\pi$ is completely inadequate to describe fragmentation functions, although the "crossed" process $\pi \to q\bar{q}$ describes distribution functions well.

Reason for failure

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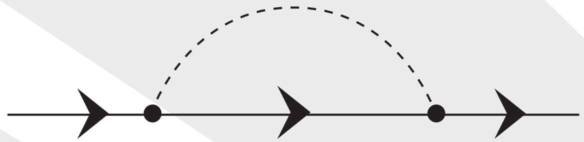
In constituent - like quark models: **Large probability ($Z_Q \simeq 0.85$) to have a quark "without its pion cloud"**. Here Z_Q is the residue of quark propagator including pion-loop self energy:



- Elementary NJL fragmentation function corresponds to the following "**number of pions per quark**":

$$\int_0^1 \mathrm{d}z \sum_{\pi} d_q^{\pi}(z) = 1 - Z_Q \simeq 0.15$$

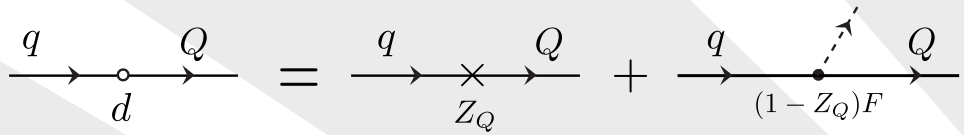
  Therefore the **pion momentum sum is small**:

$$\int_0^1 z \, \mathrm{d}z \sum_{\pi} d_q^{\pi}(z) \simeq 0.1 < 1 - Z_Q.$$

- On the other hand, empirical functions show that $\simeq 74\%$ of the initial quark momentum is converted to pions!
- $\Rightarrow$ Expect: **High-energy quark may radiate a large number of pions, and we must sum up the momenta of *all* pions!**

Product ansatz for multifragmentations

Auxiliary quantity: $d_q^Q(\eta)$ = **fragmentation function for: quark (q) $\rightarrow$ quark (Q).** [Same as: **distribution of Q inside q.**]



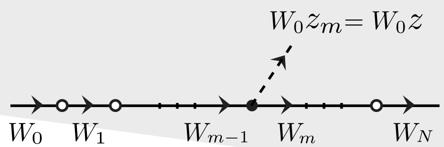
$$6\, d_q^Q(\eta) = Z_Q \delta(\eta - 1) + d_q^\pi(1-\eta) \equiv Z_Q\delta(\eta-1) + (1-Z_Q)F(\eta)$$

(Isospin indices omitted.) **Here $F(\eta)$ is normalized to 1.**

d_q^Q describes the elementary $q \rightarrow Q$ splitting $\Rightarrow$ **Product ansatz** for $D_q^\pi(z)$: **If a quark can produce a maximum of N pions, then**

$$D_q^\pi(z) = \int_0^1 \mathrm{d}\eta_1 \ldots \int_0^1 \mathrm{d}\eta_N \; 6d(\eta_1)\cdot 6d(\eta_2)\cdot \cdots \cdot 6d(\eta_N) \left(\sum_{m=1}^{N} \delta(z - z_m) \right)$$

$$D_q^\pi(z) = \sum_{m=1}^{N}$$



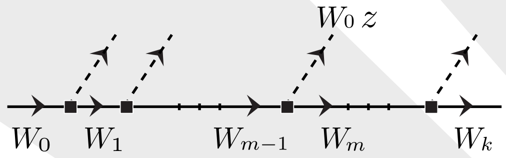
$$\left(z_m = \frac{W_{m-1} - W_m}{W_0} = \eta_1 \eta_2 \ldots \eta_{m-1}(1-\eta_m) \right)$$

Quark cascades (NJL-jet model)

What is the **physical meaning of this ansatz**?

- Rewrite the product *identically* as follows:

$$D_q^\pi(z) = \sum_{k=1}^{N} P(k) \int_0^1 d\eta_1 \ldots \int_0^1 d\eta_k \, F(\eta_1) \ldots F(\eta_k) \left(\sum_{m=1}^{k} \delta(z - z_m) \right)$$

$$D_q^\pi(z) = \sum_{k=1}^{N} P(k) \Big(\sum_{m=1}^{k}$$ 

$$\Big)$$

  $P(k)$ is the **probability that k pions are produced**:

$$P(k) = \binom{N}{k} (1 - Z_Q)^k \; Z_Q^{N-k} \quad \Rightarrow \quad \sum_{k=0}^{N} P(k) = 1.$$

  In the limit $N \to \infty$, $P(k)$ becomes a **normal distribution** with mean number (multiplicity) $\langle k \rangle = N(1 - Z_Q)$.

- In each elementary process, a fraction $\alpha \equiv \langle z F(z) \rangle < 1$ is left to the quark $\Rightarrow$ Fraction left to the final quark remainder is:

$$\sum_{k=0}^{N} P(k) \alpha^k \xrightarrow{N \to \infty} 0.$$

  $\Rightarrow$ **For N $\to \infty$, 100% of quark momentum is converted to pions!** **(Price: Divergent multiplicity.)**

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- **Pion multiplicity, momentum sum, isospin sum** for finite N:

$$\int_0^1 dz \sum_\pi D_q^\pi(z) = \sum_{k=1}^{N} kP(k) = N(1 - Z_Q)$$

$$\int_0^1 dz \sum_\pi z\, D_q^\pi(z) = 1 - \sum_{k=0}^{N} P(k)\langle zF\rangle^k = 1 - \left(Z_Q + (1 - Z_Q)\langle zF\rangle\right)^N$$

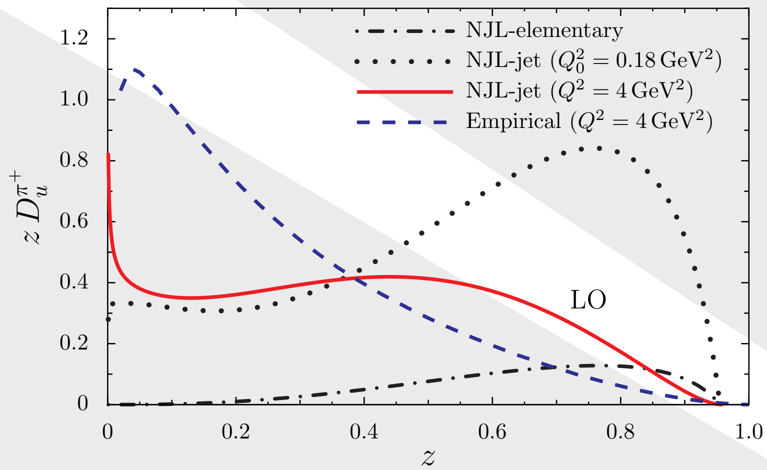
$$\int_0^1 dz \sum_\pi \tau_\pi\, D_q^\pi(z) = \frac{\tau_q}{2}\left[1 - \sum_{k=0}^{N} P(k)\left(-\frac{1}{3}\right)^k\right]$$

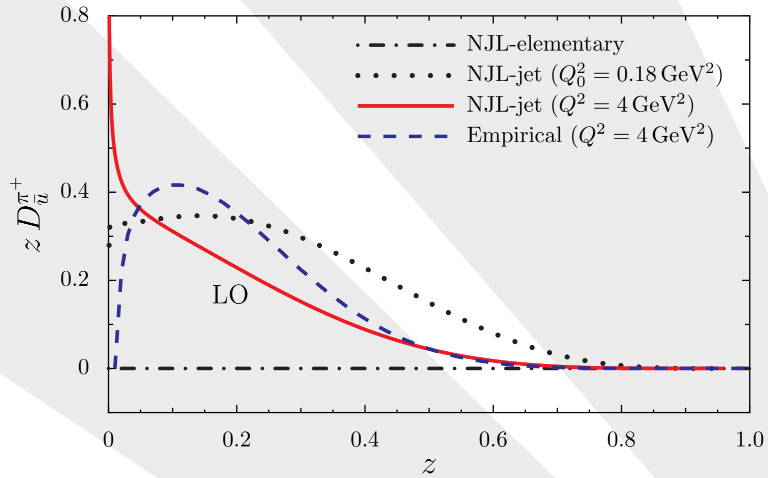
$$= \frac{\tau_q}{2}\left[1 - \left(Z_Q - \frac{1}{3}(1 - Z_Q)\right)^N\right]$$

- For $N \to \infty$, $D_q^\pi(z)$ satisfies the **same integral equation (chain equation) as in the original Field-Feynman model**:

$$D_q^\pi(z) = d_q^\pi(z)/(1 - Z_Q) + \sum_Q \left[F_q^Q \otimes D_Q^\pi\right](z)$$

NJL-jet model: Pions only ($M = 300$ MeV)





- **Cascade-like processes enhance the fragmentation functions tremendously!**
- Calculated functions are still **too stiff** because:
  - Q^2 evolution should be performed in **NLO**. (At present, codes are not available for the public . . .)
  - Some of observed pions are **secondary pions** (from decay of vector mesons).

NJL-jet model: Pions + Kaons ($M_s = 450$ MeV)

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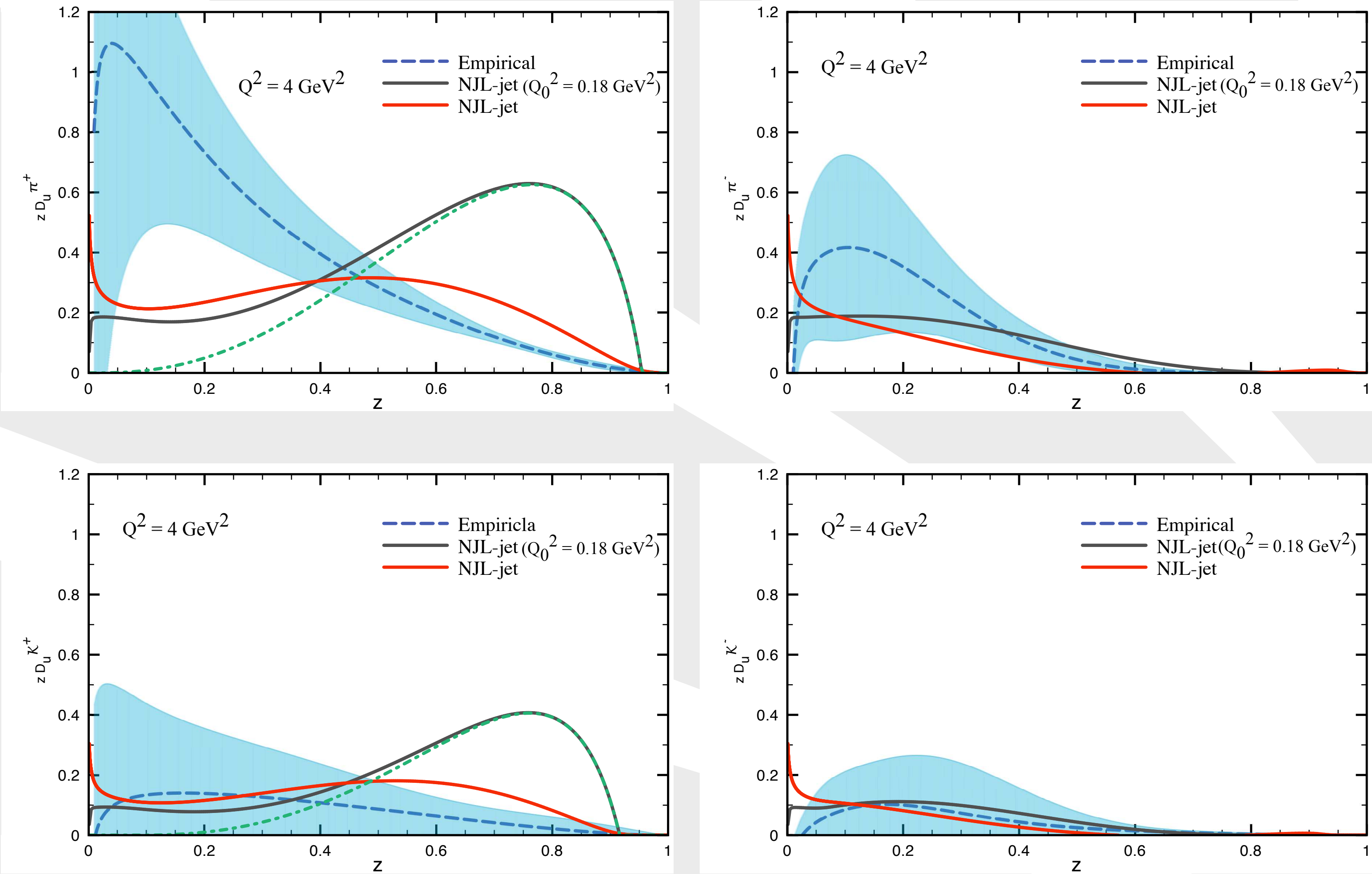

- Momentum of u-quark: 67$\%$ to π, 33$\%$ to K.
  $D_u^\pi(z)$ is approximately scaled down by factor 2/3.
- Isospin of u-quark: 80$\%$ to π, 20$\%$ to K.

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- **Cascade - type multifragmentation processes are extremely important to describe fragmentation functions.**
- **The "NJL-jet model" describes qualitatively the empirical fragmentation functions without any new parameters.**
- **Straight forward extensions will improve the description: NLO effects in Q^2 evolution; inclusion of vector meson and nucleon channel.**
- **Important: The product ansatz should be derived from field theory ($\Leftrightarrow$ rainbow-ladder approximation for quark self energy).**

Outlook: SIDIS on nuclear targets

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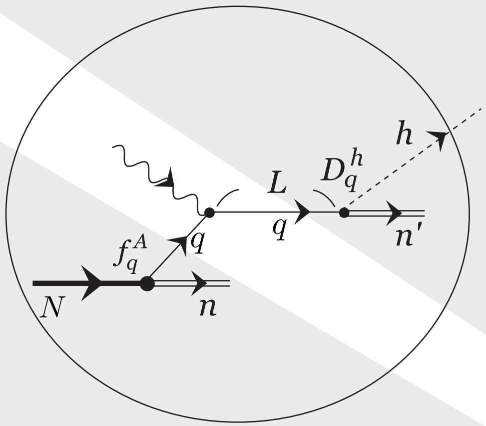
Recent data for *nuclear* targets (HERMES, JLab) indicate **medium modification of SIDIS process**:



Nucleus *A*

- Medium modification of $f_q^A \Leftrightarrow$ EMC effect (talk of Ian Cloet).
- Average hadron formation length L.
- Energy loss of propagating quark in medium.
- Medium modification of D_q^h.
- Interaction of hadron h with medium.