

Ab Initio No Core Full Configuration approach to the structure of light nuclei

Pieter Maris
pmaris@iastate.edu
Iowa State University

IOWA STATE
UNIVERSITY



UNEDF SciDAC Collaboration

Universal Nuclear Energy Density Functional

SciDAC project – UNEDF

spokespersons: Rusty Lusk (ANL), Witek Nazarewicz (ORNL/UT)

<http://www.unedf.org>

INCITE award – Computational Nuclear Structure

PI: David Dean (ORNL)



PetaApps award

PIs: Jerry Draayer (LSU), Umit Catalyurek (OSU)

Masha Sosonkina, James Vary (ISU)

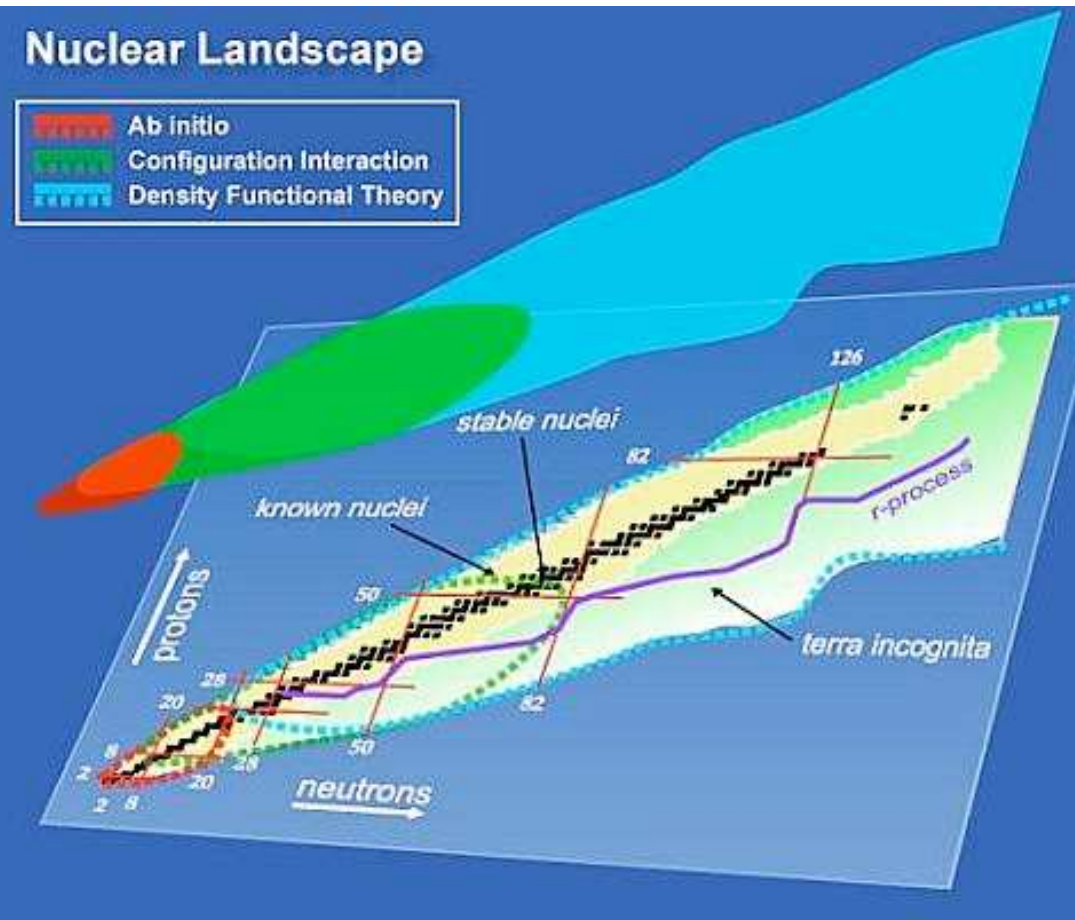


SciDAC/UNEDF – Uniform description of nuclear structure

SciDAC project – UNEDF

spokespersons: Rusty Lusk (ANL), Witek Nazarewicz (ORNL/UT)

<http://www.unedf.org>



● Universal Nuclear Energy Density Functional theory that spans the entire mass table based on **ab initio** calculations for light nuclei

- Coupled Cluster (Papenbrock *et al*, ORNL)
- No-Core Configuration Interaction calculations
numerical code: **MFDn**
- Greens Function Monte Carlo (Pieper *et al*, ANL)

Ab initio nuclear structure calculations

Eigenvalue problem $\hat{H}|\Psi\rangle = E|\Psi\rangle$ with Hamiltonian

$$\hat{H} = \sum_{i < j} \frac{(\vec{p}_i - \vec{p}_j)^2}{2 m A} + \sum_{i < j} V_{ij} + \sum_{i < j < k} V_{ijk} + \dots$$

- Self-bound quantum many-body problem, with $3A$ degrees of freedom in coordinate or momentum space
- Strong interactions, with both short-range and long-range pieces
- Multiple scales, from keV's to MeV's



Jaguar



Franklin



Blue Gene/p



Atlas

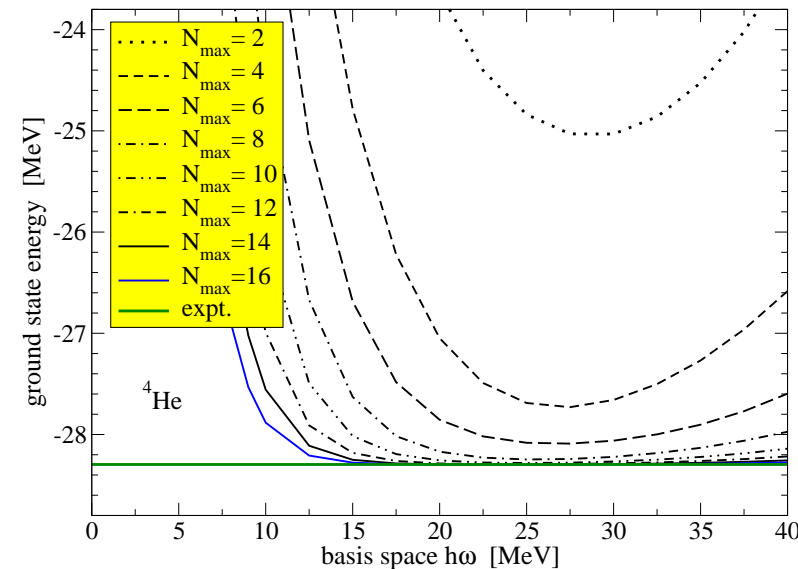


Configuration Interaction Methods

- Expand wave function in basis states $|\Psi\rangle = \sum a_i |\psi_i\rangle$
- Express Hamiltonian in basis $\langle \psi_j | \hat{\mathbf{H}} | \psi_i \rangle = H_{ij}$
- Diagonalize Hamiltonian matrix H_{ij}
- Complete basis \longrightarrow exact result
 - caveat: complete basis is infinite dimensional
- In practice
 - truncate basis
 - study behavior of observables as function of truncation
- Computational challenge
 - construct large ($10^{10} \times 10^{10}$) sparse symmetric real matrix H_{ij}
 - use Lanczos algorithm to obtain lowest eigenvalues & eigenvectors

Configuration Interaction Methods

- Expand wave function in basis states $|\Psi\rangle = \sum a_i |\psi_i\rangle$
- Express Hamiltonian in basis $\langle \psi_j | \hat{H} | \psi_i \rangle = H_{ij}$
- Diagonalize Hamiltonian matrix H_{ij}
- **Variational**: for any finite truncation of the basis space, eigenvalue is an upper bound for the ground state energy
- Smooth approach to asymptotic value with increasing basis space:
No-Core Full Configuration calculation
- Convergence: **independence** of N_{\max} and H.O. basis $\hbar\omega$
 - different methods (NCFC, CC, GFMC, ...) using the same interaction should give same results within numerical errors

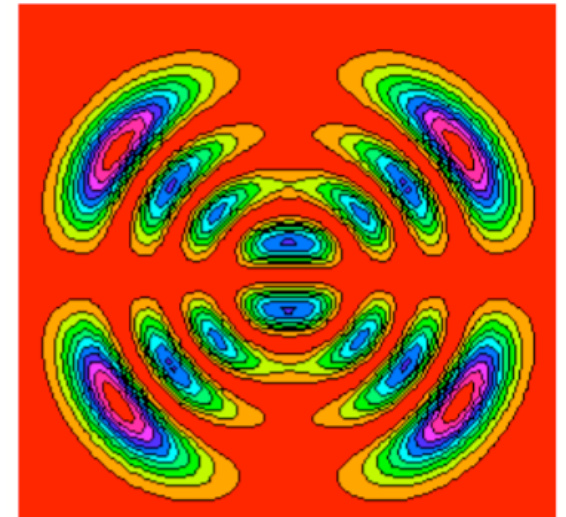


Many-Body Basis Space

- Expand wave function in basis states $|\Psi\rangle = \sum a_i |\psi_i\rangle$
- Many-Body basis states $|\psi_i\rangle$
 - Slater Determinants of single-particle states $|\phi\rangle$

$$|\psi\rangle = |\phi_1\rangle \otimes \dots \otimes |\phi_A\rangle$$

- single-particle basis states
eigenstates of SU(2) operators
 $\hat{\mathbf{L}}^2, \hat{\mathbf{S}}^2, \hat{\mathbf{J}}^2 = (\hat{\mathbf{L}} + \hat{\mathbf{S}})^2$, and $\hat{\mathbf{J}}_z$
w. quantum numbers $|\phi\rangle = |n, l, s, j, m\rangle$
- radial wavefunctions:
Harmonic Oscillator



sample harmonic oscillator basis function

- M -scheme: many-body basis states eigenstates of $\hat{\mathbf{J}}_z$

$$\hat{\mathbf{J}}_z |\psi\rangle = M |\psi\rangle = \sum_{i=1}^A m_i |\psi\rangle$$

- Alternatives: LS -scheme, **Total- J -scheme**, **Symplectic basis**, . . .

Truncation Schemes

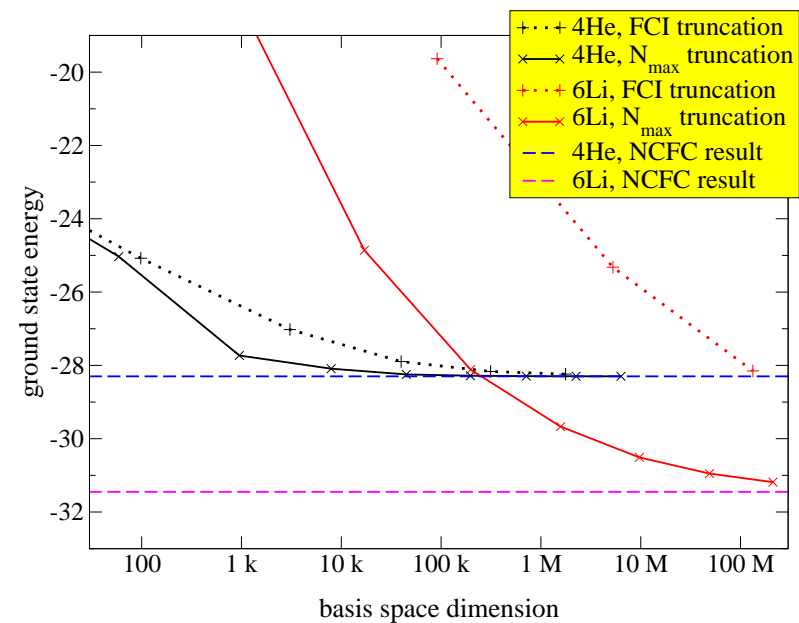
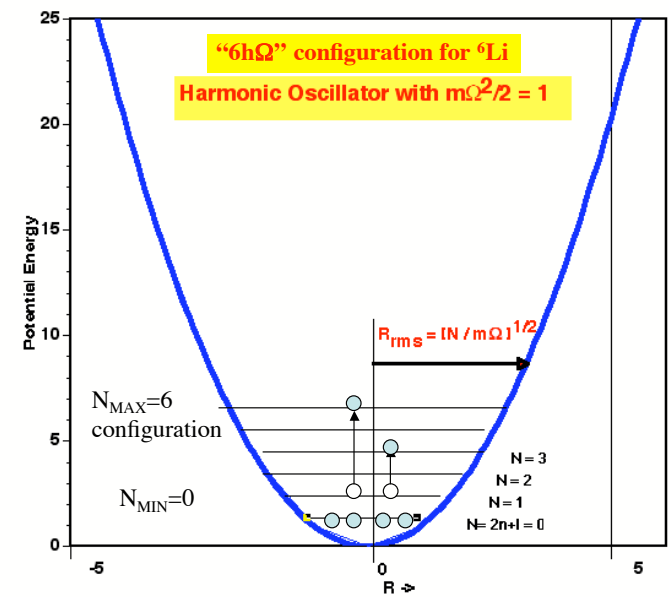
- N_{\max} truncation
 - truncation on the total number of H.O. oscillator quanta above minimal configuration for that nucleus

- allows for exact separation of Center-of-Mass motion and intrinsic motion

- Alternative truncation schemes

- **FCI** – Full Configuration Interaction – truncation on single-particle basis only

- Importance Sampling, **Monte Carlo Sampling**, **Symplectic**, . . .



Intermezzo: Center-of-Mass excitations

- Use single-particle coordinates, not relative (Jacobi) coordinates
 - straightforward to extend to many particles
 - have to separate Center-of-Mass motion from intrinsic motion
- Add Lagrange multiplier to Hamiltonian

$$\hat{\mathbf{H}}_{\text{rel}} \longrightarrow \hat{\mathbf{H}}_{\text{rel}} + \Lambda_{CM} \left(\hat{\mathbf{H}}_{CM}^{H.O.} - \frac{3}{2} \left(\sum_i m_i \right) \omega \right)$$

with $\hat{\mathbf{H}}_{\text{rel}} = T_{\text{rel}} + V_{\text{rel}}$ the relative Hamiltonian

- separates CM excitations from CM ground state $|\Phi_{CM}\rangle$
- Center-of-Mass wave function **factorizes** for **H.O. basis functions** in combination with **N_{max} truncation**

$$\begin{aligned} |\Psi_{\text{total}}\rangle &= |\phi_1\rangle \otimes \dots \otimes |\phi_A\rangle \\ &= |\Phi_{\text{Center-of-Mass}}\rangle \otimes |\Psi_{\text{intrinsic}}\rangle \end{aligned}$$

where

$$\hat{\mathbf{H}}_{\text{rel}} |\Psi_j, \text{intrinsic}\rangle = E_j |\Psi_j, \text{intrinsic}\rangle$$

Configuration Interaction Methods

- Expand wave function in basis states $|\Psi\rangle = \sum a_i |\psi_i\rangle$
 - Slater Determinants of single-particle states
 - N_{\max} truncation on total number of H.O. quanta
- Express Hamiltonian in basis $\langle \psi_j | \hat{\mathbf{H}} | \psi_i \rangle = H_{ij}$

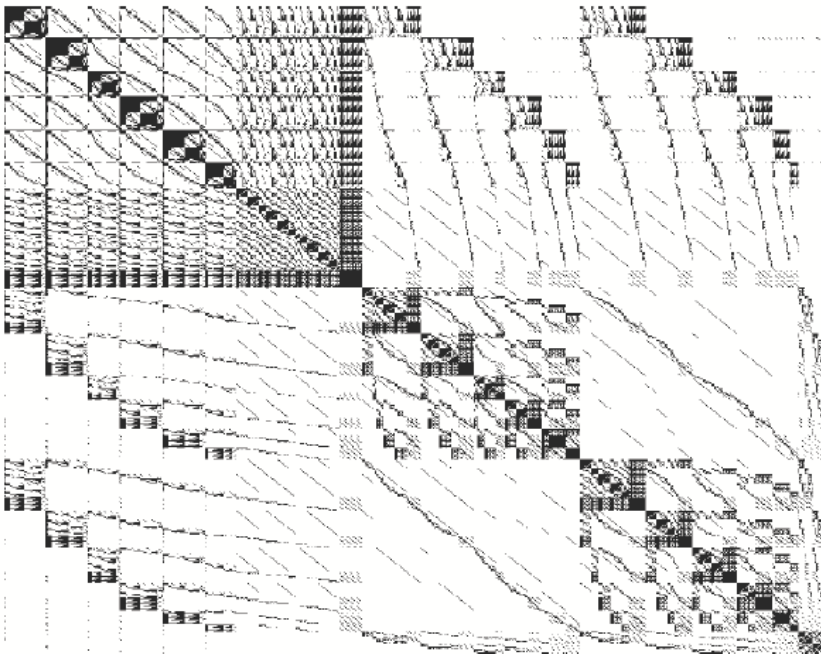
$$\hat{\mathbf{H}} = \hat{\mathbf{T}}_{\text{rel}} + \Lambda_{CM} \left(\hat{\mathbf{H}}_{CM}^{H.O.} - \frac{3}{2} \left(\sum_i m_i \right) \omega \right) + \sum_{i < j} V_{ij} + \sum_{i < j < k} V_{ijk} + \dots$$

- Argonne potentials: AV8, AV18 (plus Illinois NNN interactions)
- Bonn potentials
- Chiral NN interactions (plus chiral NNN interactions)
- ...
- JISP16 (phenomenological NN potential)
- ...

Configuration Interaction Methods

- Expand wave function in basis states $|\Psi\rangle = \sum a_i |\psi_i\rangle$
 - Slater Determinants of single-particle states
 - N_{\max} truncation on total number of H.O. quanta
- Express Hamiltonian in basis $\langle \psi_j | \hat{H} | \psi_i \rangle = H_{ij}$
 - large sparse symmetric matrix

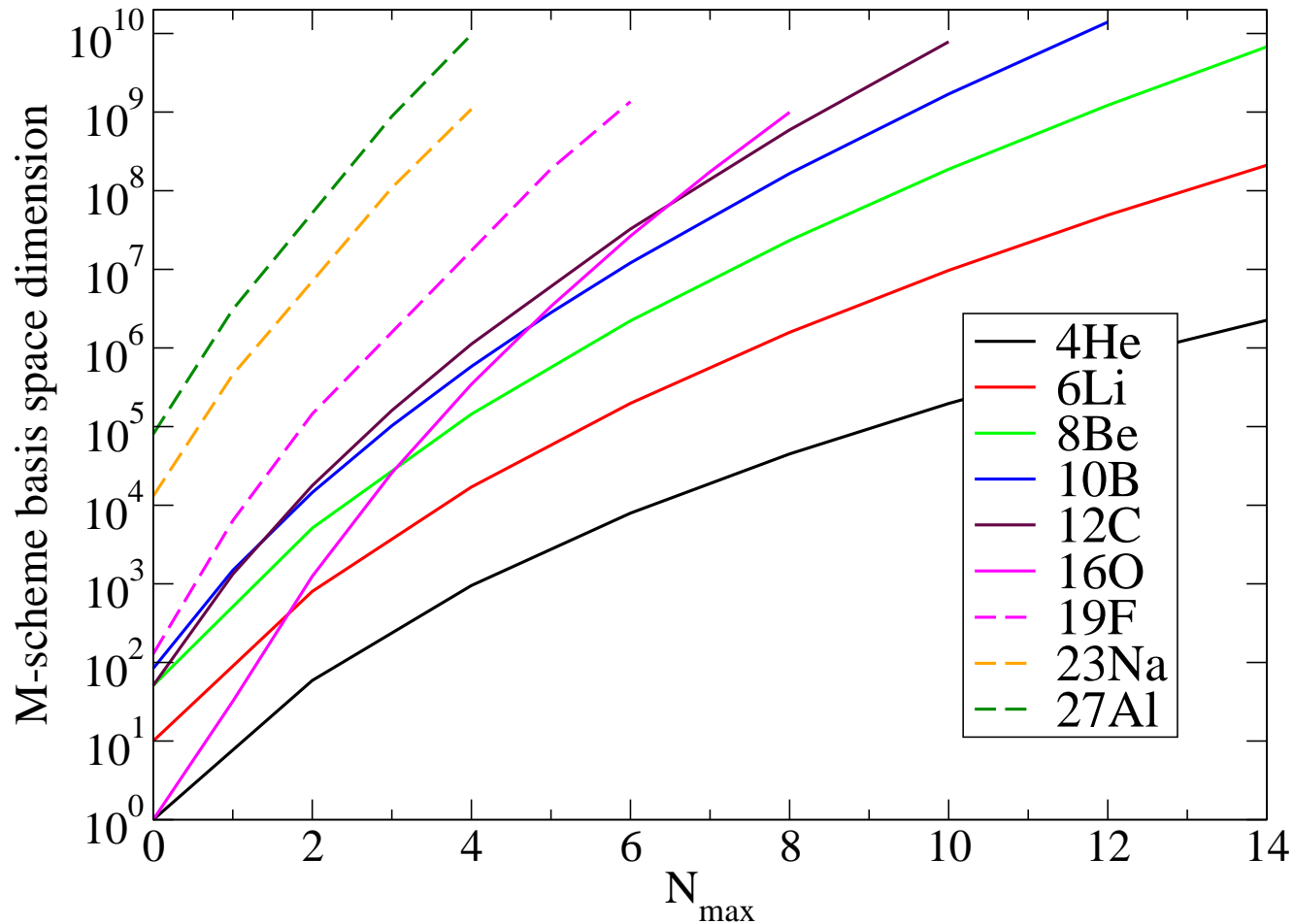
Sparsity Structure for ${}^6\text{Li}$



- Obtain lowest eigenvalues using Lanczos algorithm
 - Eigenvalues: bound state spectrum
 - Eigenvectors: nuclear wavefunctions

CI calculations – main challenge

Single most important computational issue:
exponential increase of dimensionality with increasing H.O. levels



Intermezzo: Extrapolation Techniques

Challenge: achieve numerical convergence for no-core Full Configuration calculations using finite model space calculations

- Perform a series of calculations with increasing N_{\max} truncation (while keeping everything else fixed)
- Extrapolate to infinite model space \longrightarrow exact results
 - binding energy: exponential in N_{\max}

$$E_{\text{binding}}^N = E_{\text{binding}}^{\infty} + a_1 \exp(-a_2 N_{\max})$$

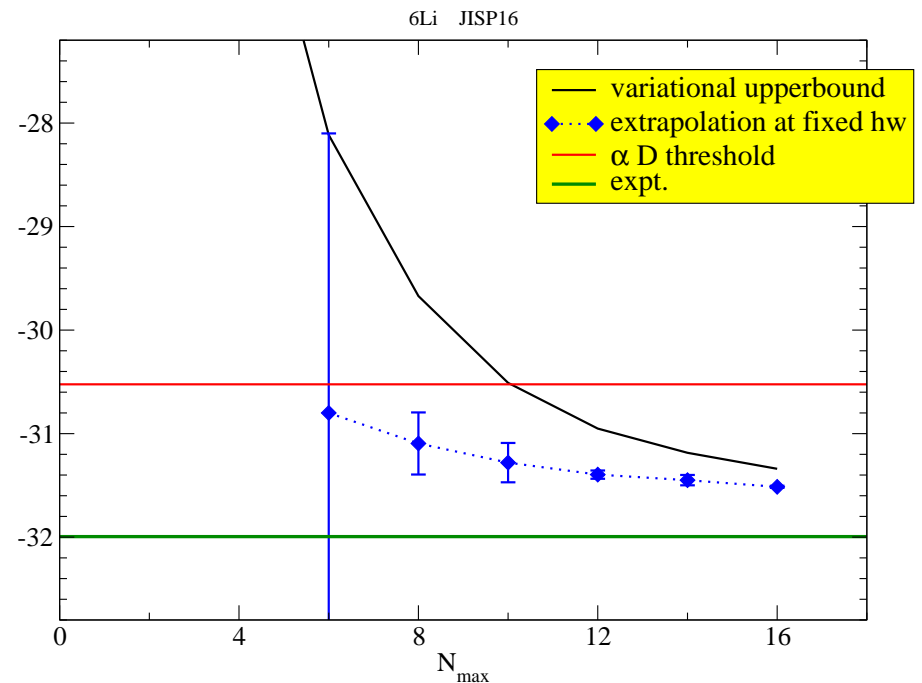
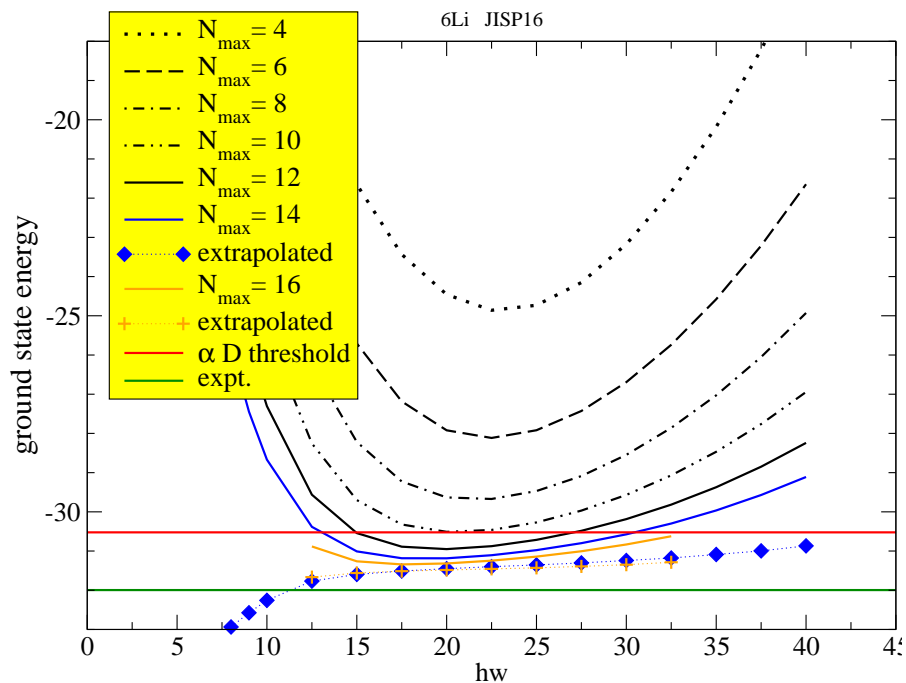
- use 3 or 4 consecutive N_{\max} values to determine $E_{\text{binding}}^{\infty}$
- use $\hbar\omega$ and N_{\max} dependence to estimate numerical error bars

Maris, Shirokov, Vary, Phys. Rev. C79, 014308 (2009)

Intermezzo: Extrapolation Techniques

Challenge: achieve numerical convergence for no-core Full Configuration calculations using finite model space calculations

- Perform a series of calculations with increasing N_{\max} truncation (while keeping everything else fixed)
- Extrapolate to infinite model space \longrightarrow exact results



Intermezzo: Bare vs. Renormalized Hamiltonians

Challenge: achieve numerical convergence for no-core Full Configuration calculations using finite model space calculations

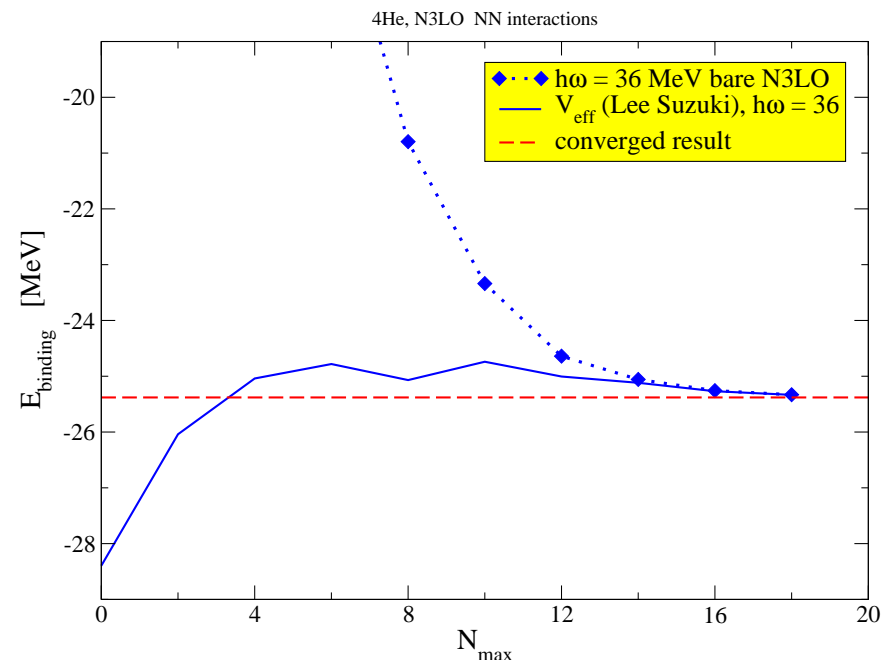
- Extrapolate to infinite model space \longrightarrow exact results
- Renormalize interaction \longrightarrow effective interaction V_{eff}
 - can improve quality of results in small model spaces
- Caveats
 - induces many-body forces
 - induced 3-body forces are often neglected
 - induced 4-, 5-, ..., A -body forces are always neglected
 - variational principle applicable to renormalized Hamiltonian not to original (bare) Hamiltonian
 - often complicates extrapolation to asymptotic values
 - need to renormalize operators as well

Intermezzo: Bare vs. Renormalized Hamiltonians

Challenge: achieve numerical convergence for no-core Full Configuration calculations using finite model space calculations

- Extrapolate to infinite model space \longrightarrow exact results
- Renormalize interaction \longrightarrow effective interaction V_{eff}
- Commonly used in No Core Shell Model calculations: Lee–Suzuki Renormalization

- different V_{eff} for each nucleus and each model space
- disadvantages
 - no variational principle
 - approach to continuum limit non-monotonic



Intermezzo: Bare vs. Renormalized Hamiltonians

Challenge: achieve numerical convergence for no-core Full Configuration calculations using finite model space calculations

- Extrapolate to infinite model space \longrightarrow exact results
- Renormalize interaction \longrightarrow effective interaction V_{eff}
- Commonly used in No Core Shell Model calculations:
Lee–Suzuki Renormalization
- SRG: Similarity transformations applied to the free-space Hamiltonian
 - same renormalized Ham. for all nuclei and model spaces
 - variational principle applicable to renormalized Hamiltonian
 - smooth approach to continuum limit
 - currently working on renormalized operators (OSU group)
- Other methods: $V_{\text{low } k}$, V_{UCOM}

Nuclear interaction

Strong force between nucleons not well known . . .

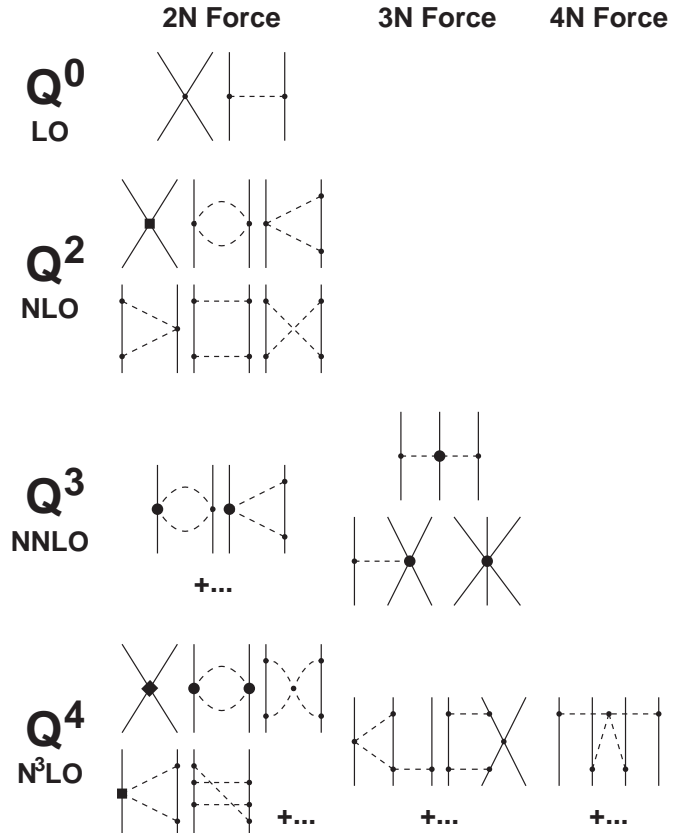
. . . but should be calculable from QCD

- Use **chiral perturbation theory** to obtain effective A -body interaction from QCD Entem and Machleidt, Phys. Rev. C68, 041001 (2003)

- controlled power series expansion in Q/Λ_χ with $\Lambda_\chi \sim 1$ GeV
- natural hierarchy for many-body forces

$$V_{NN} \gg V_{NNN} \gg V_{NNNN}$$

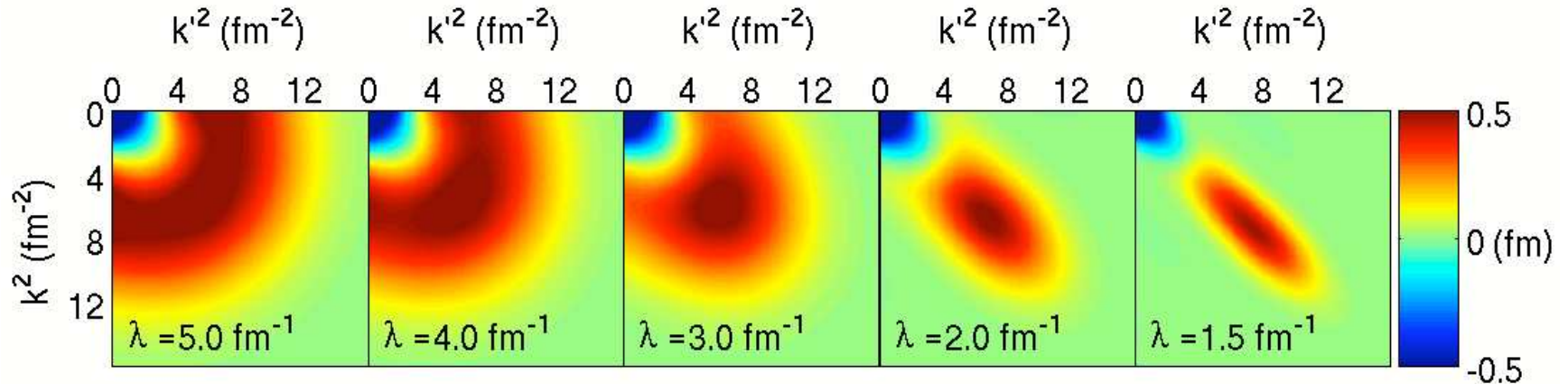
- in principle no free parameters
- in practice a few undetermined parameters
- renormalization necessary



Similarity Renormalization Group – NN interaction

SRG evolution

Bogner, Furnstahl, Perry, PRC 75 (2007) 061001

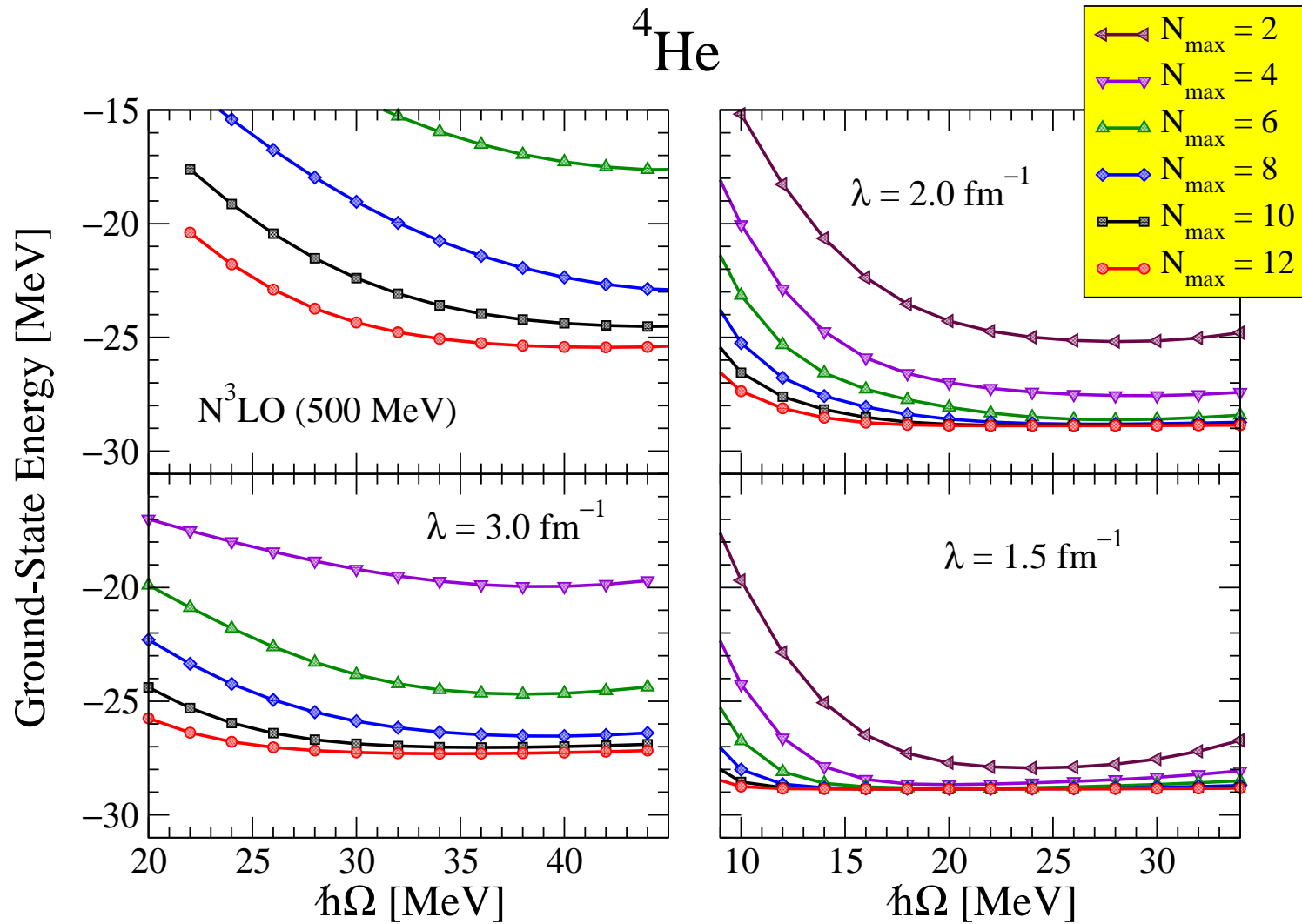


- drives interaction towards band-diagonal structure
 - SRG shifts strength between 2-body and many-body forces
- Initial chiral EFT Hamiltonian
power-counting hierarchy A -body forces

$$V_{NN} \gg V_{NNN} \gg V_{NNNN}$$

- key issue: preserve hierarchy of many-body forces

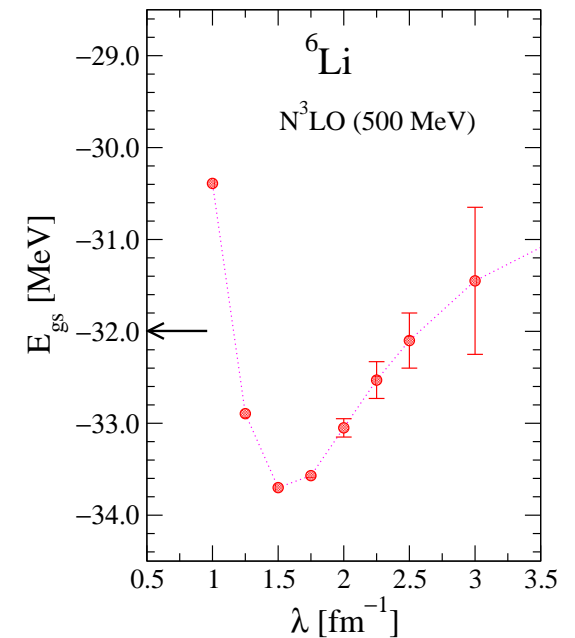
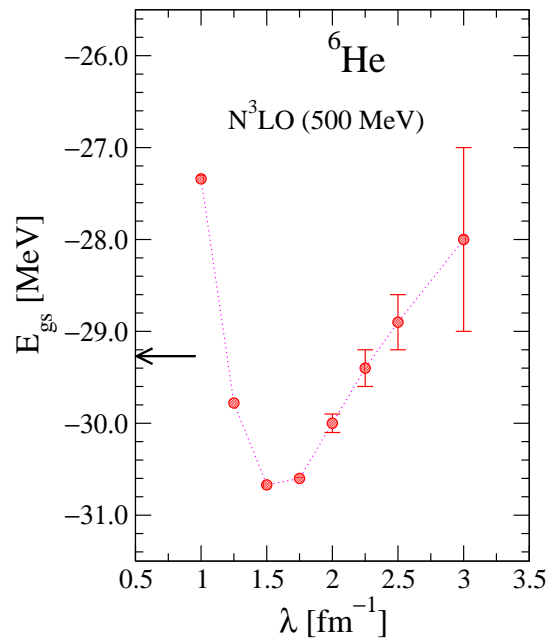
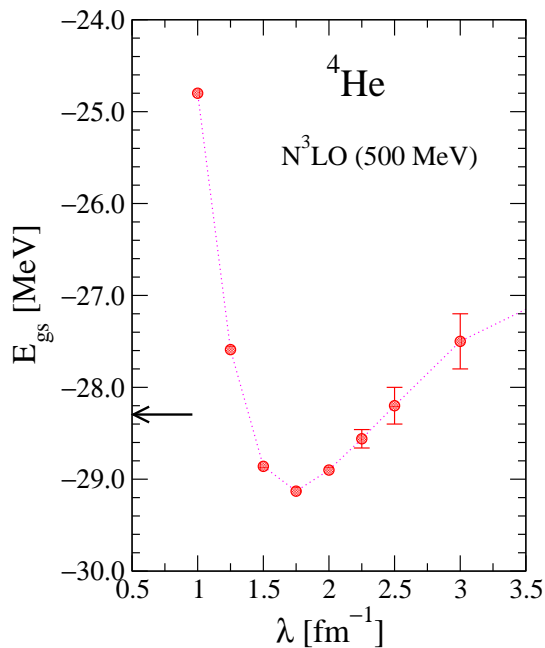
Improve convergence rate by applying SRG to N3LO



(Bogner, Furnstahl, Maris, Perry, Schwenk, Vary, NPA801, 21 (2008), arXiv:0708.3754)

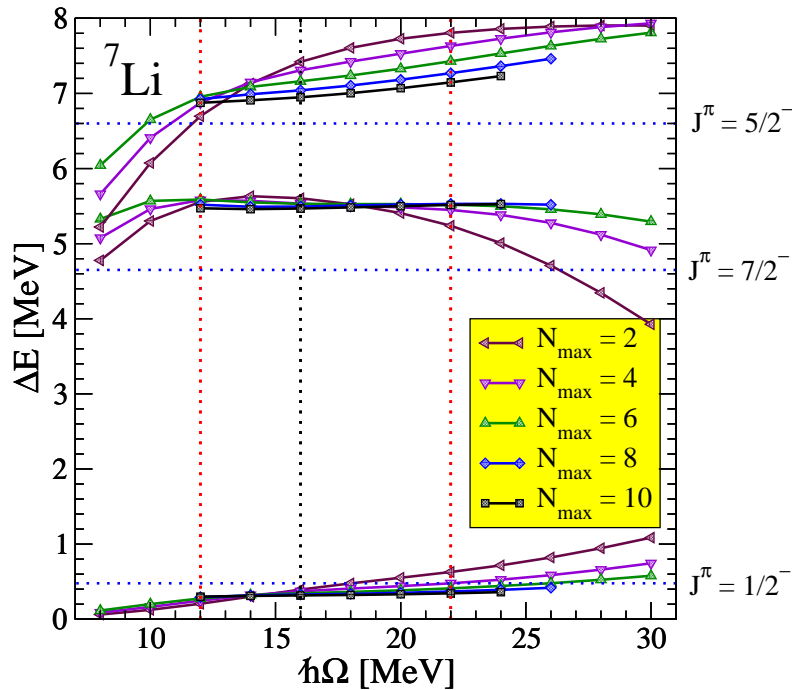
Results with SRG-evolved N3LO NN interaction

- Convergence rate significantly improved by SRG evolution
- Net contribution of many-body body forces remain “natural” down to $\lambda \approx 1.2 \sim 1.5 \text{ fm}^{-1}$

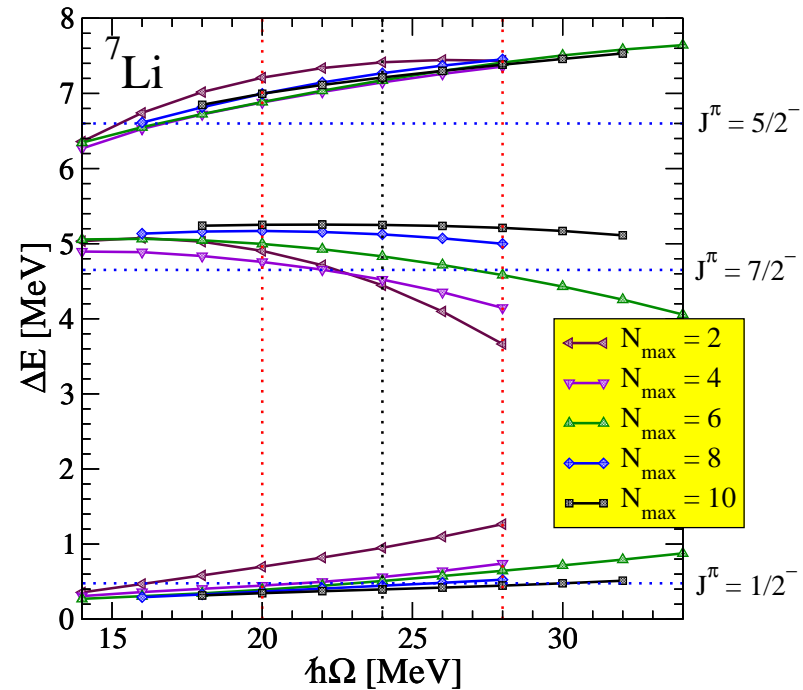


(Bogner, Furnstahl, Maris, Perry, Schwenk, Vary, NPA801, 21 (2008), arXiv:0708.3754)

Spectrum of ${}^7\text{Li}$ with SRG-evolved N3LO NN interaction



$$\lambda = 1.50$$

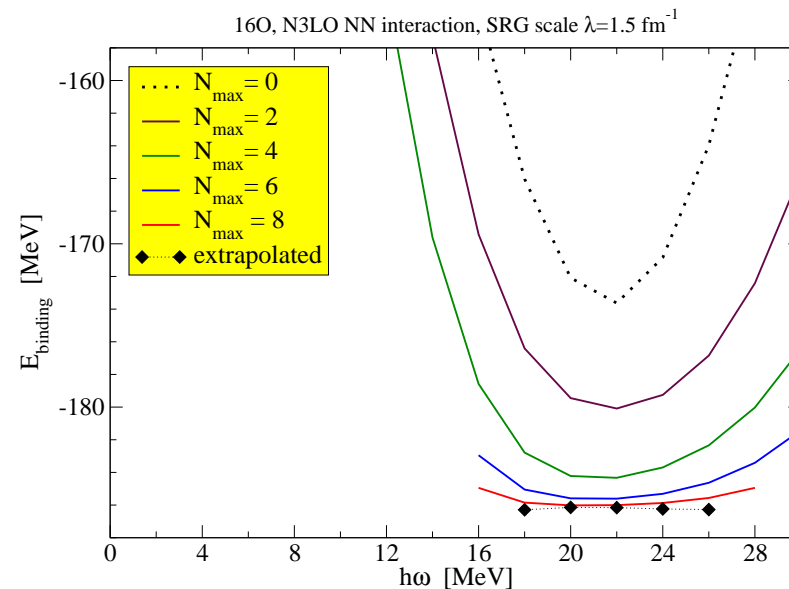
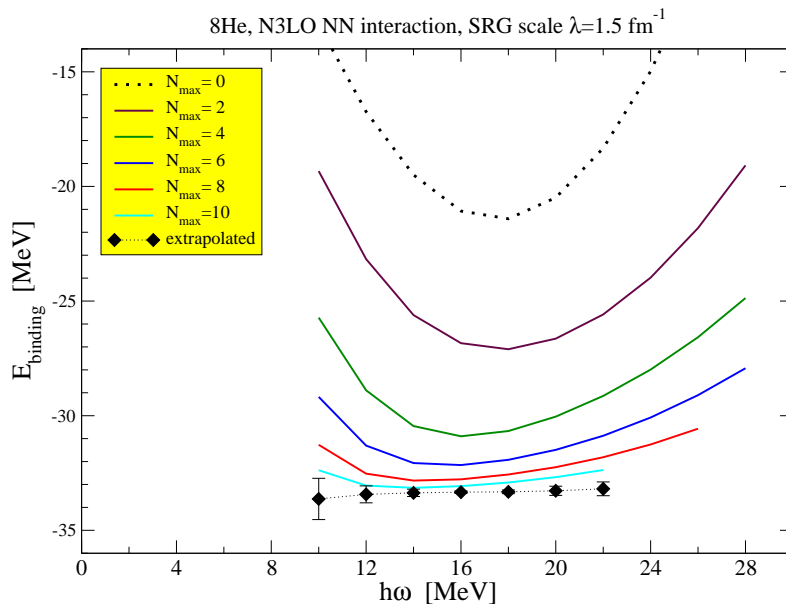


$$\lambda = 2.50$$

- Lowest two excited states well-converged
- Spectrum remarkably insensitive to λ , despite difference in ground state energy of about 3 MeV

More results with SRG-evolved N3LO NN interaction

- Rapid convergence 8He, 12C, 16O, 40Ca at $\lambda = 1.5 \text{ fm}^{-1}$
 - check between our results and Coupled Cluster (Dean *et al*)
- All overbound at $\lambda = 1.5 \text{ fm}^{-1}$

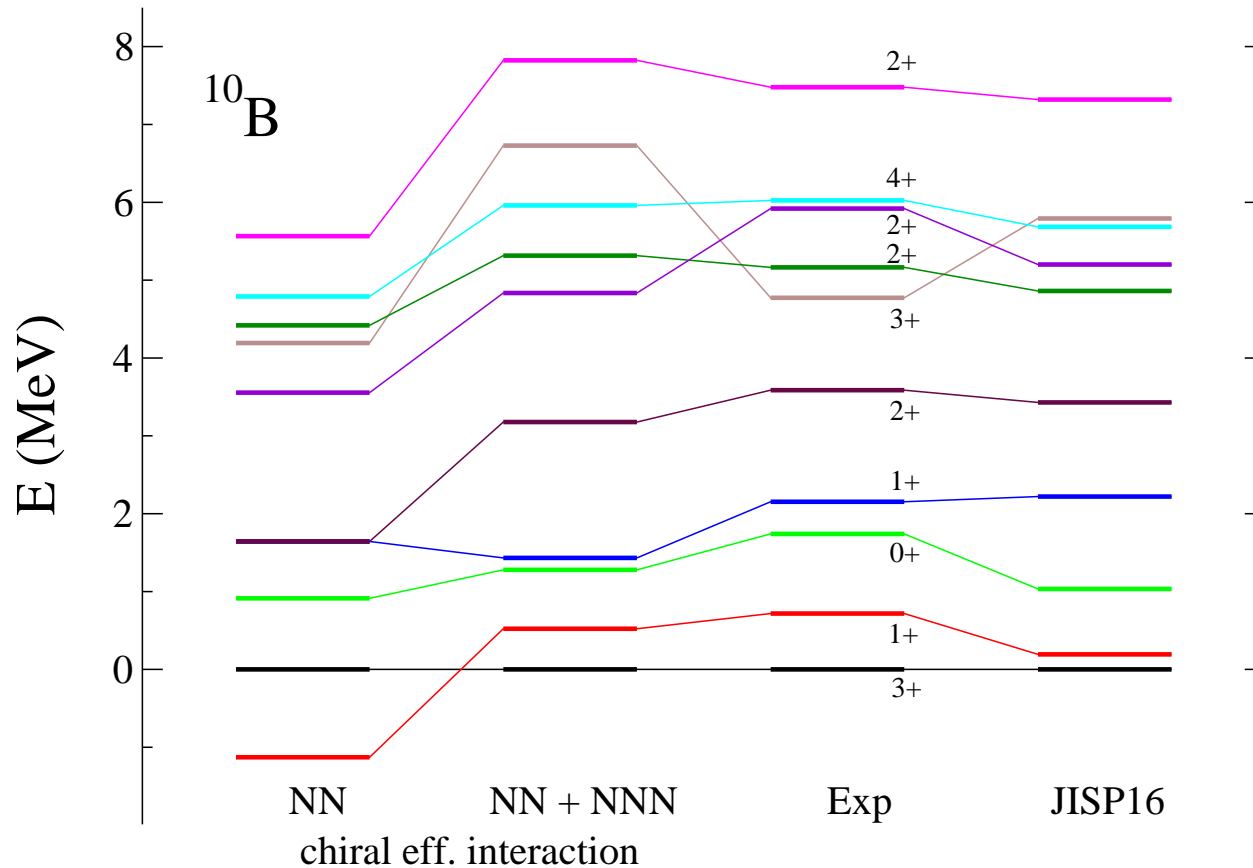


- Also overbound at $\lambda = 2.5 \text{ fm}^{-1}$: 3-body forces necessary
- 3-body forces also necessary for spectra, ...

Navratil, Gueorguiev, Vary, Ormand, Nogga, PRL99, 014315 (2007)

- Next step: include NNN forces (LLNL group)

Results with chiral N3LO (3-body) and JISP16 (2-body)



Spectrum of ^{10}B

with chiral 2- and
3-body forces
at $N_{\text{max}} = 6$

nonlocal 2-body
interaction JISP16
at $N_{\text{max}} = 8$

Vary, Maris, Negoita, Navratil, Gueorguiev, Ormand, Nogga, Shirokov, and Stoica,
in “Exotic Nuclei and Nuclear/Particle Astrophysics (II), Proceedings of the Carpathian
Summer School of Physics 2007, AIP Conference Proceedings 972, 49 (2008)

Phenomenological NN interaction: JISP16

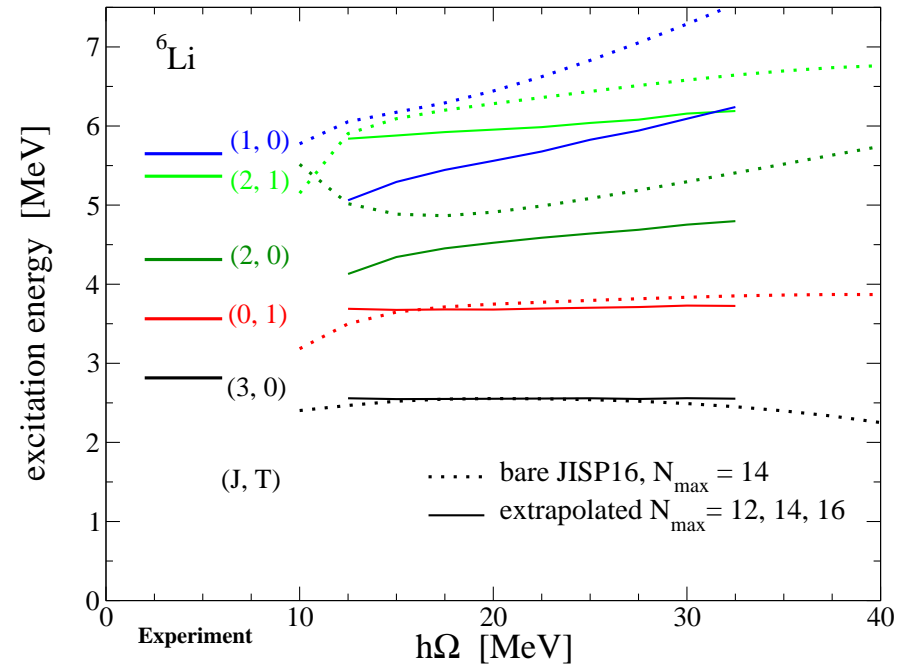
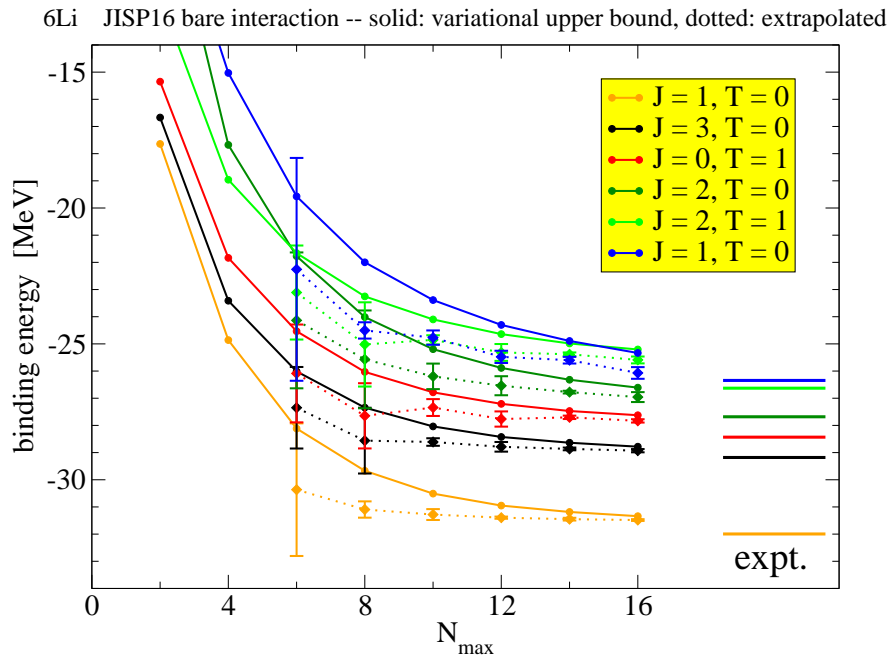
A.M. Shirokov, J.P. Vary, A.I. Mazur, T.A. Weber, PLB 644, 33 (2007)

J-matrix Inverse Scattering Potential tuned up to ^{16}O

- finite rank separable potential in H.O. representation
- fitted to available NN scattering data
- use unitary transformations to tune off-shell interaction to
 - binding energy of ^3He
 - low-lying spectrum of ^6Li (JISP6, precursor to JISP16)
 - binding energy of ^{16}O
- good fit to a range of light nuclear properties
- very soft potential compared to other NN potentials
- nonlocal potential (by construction)
- details available at

<http://nuclear.physics.iastate.edu/>

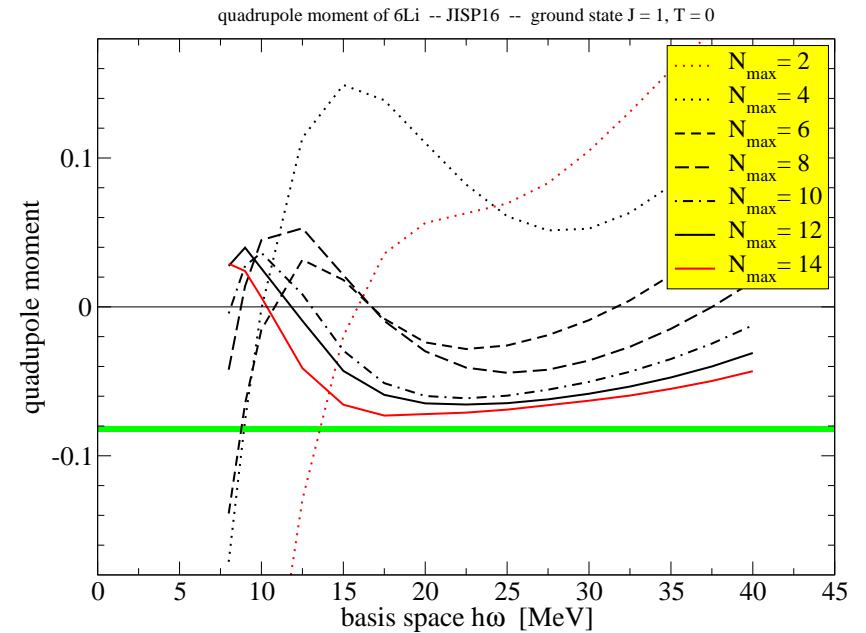
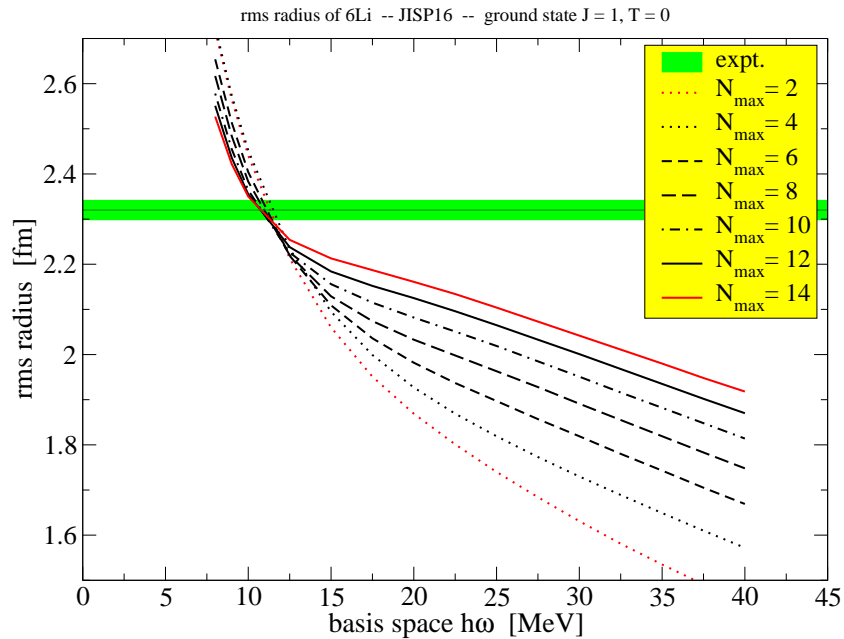
Results with JISP16 for ${}^6\text{Li}$



- Spectrum in good agreement with experimental data
 - in particular for narrow states
- Broad resonances \leftrightarrow strong dependence on $\hbar\Omega$
 - agrees with inverse scattering analysis of α -nucleon states
Shirokov, Mazur, Vary, Mazur, PRD79, 014610 (2009)
 - numerical error under investigation

Results with JISP16 for ${}^6\text{Li}$

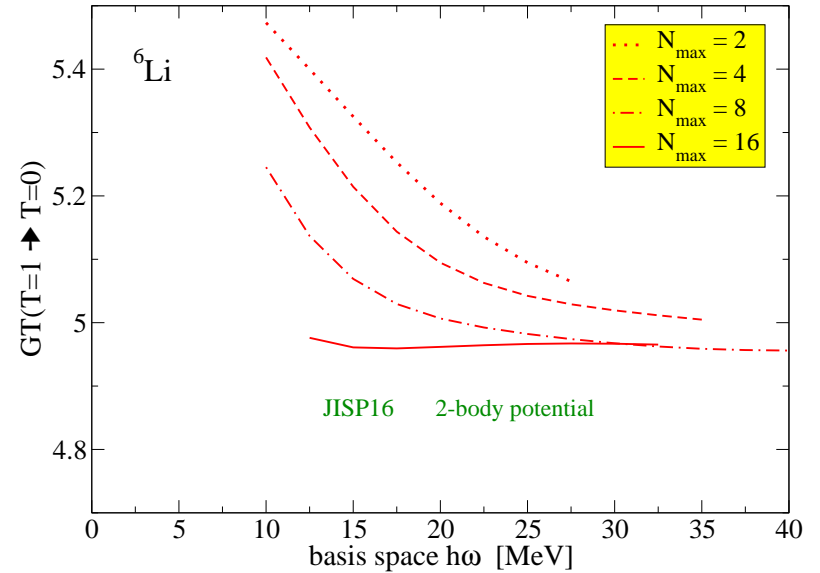
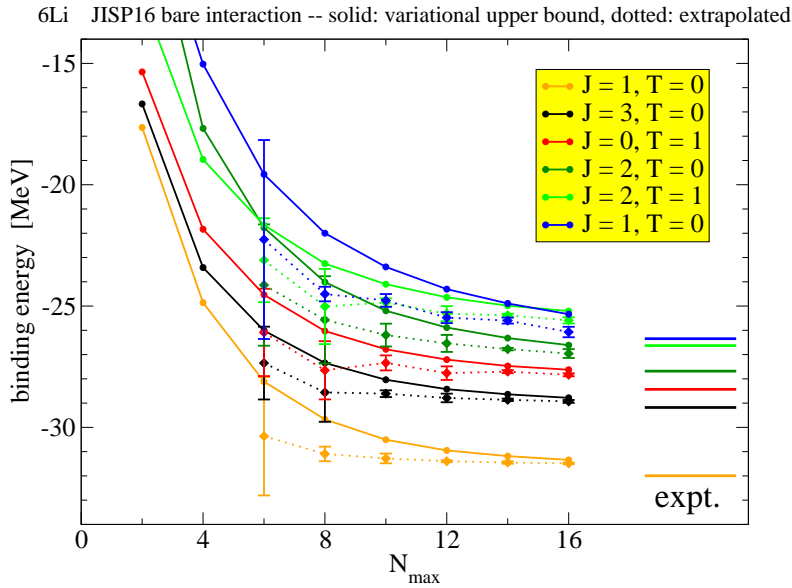
Ground state $\langle r^2 \rangle$ and Q



- Quadrupole moment Q agrees with data within uncertainties
- RMS charge radius in the right ballpark
- Not very well converged:
 - need realistic basis function

Negoita, PhD work in progress

Gamov–Teller transition in ${}^6\text{Li}$ and ${}^6\text{He}$



Calculation within ${}^6\text{Li}$ good approximation for ${}^6\text{He}(0^+, 1) \rightarrow {}^6\text{Li}(1^+, 0)$

- Corrections due to meson-exchange currents small

within ${}^6\text{Li}$	${}^6\text{He}(0^+, 1) \rightarrow {}^6\text{Li}(1^+, 0)$	including MEC corrections	expt.
4.96	4.95	4.83(3)	4.71

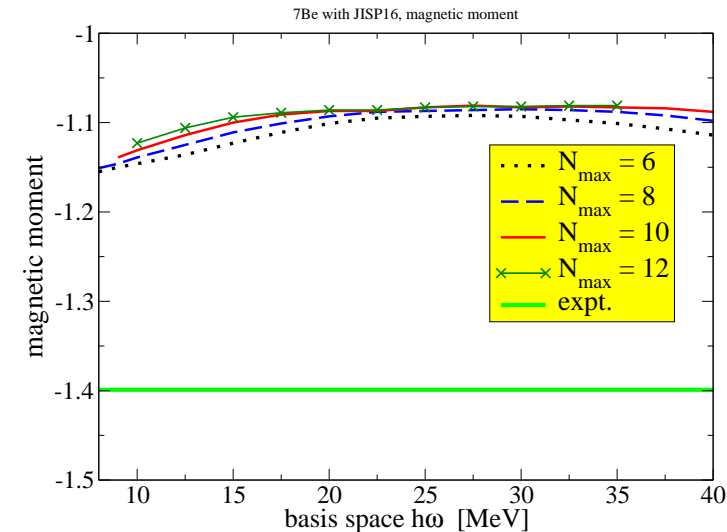
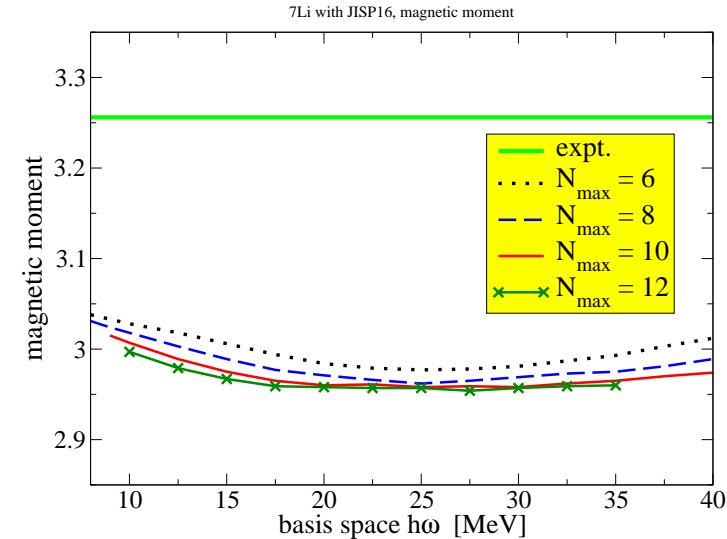
Vaintraub, Barnea, Gazit, arXiv:0903.1048 [nucl-th]

- Excellent agreement between JISP16 results and experiment

Results with JISP16 for ${}^7\text{Li}$ and ${}^7\text{Be}$ – magnetic moments

	expt.	$N_{\text{max}} = 12$ at $\hbar\omega = 20$ MeV	
${}^7\text{Li}$, E	39.270	-38.58(4)	extrapolated
${}^7\text{Li}$, r_p	2.27(2)	2.07	
${}^7\text{Li}$, μ	+3.256	+2.96(1)	converged
${}^7\text{Li}$, Q	-4.00(6)	-2.79	
${}^7\text{Be}$, E	37.600	-36.94(5)	extrapolated
${}^7\text{Be}$, r_p	2.36(2)	2.234	
${}^7\text{Be}$, μ	-1.399	-1.09(1)	converged
${}^7\text{Be}$, Q		-4.695	

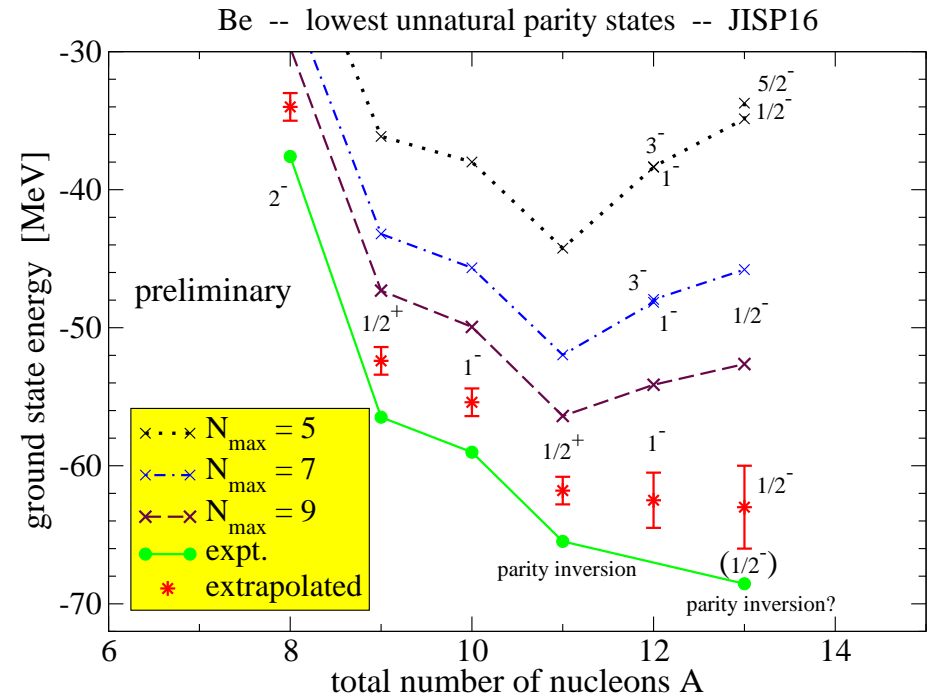
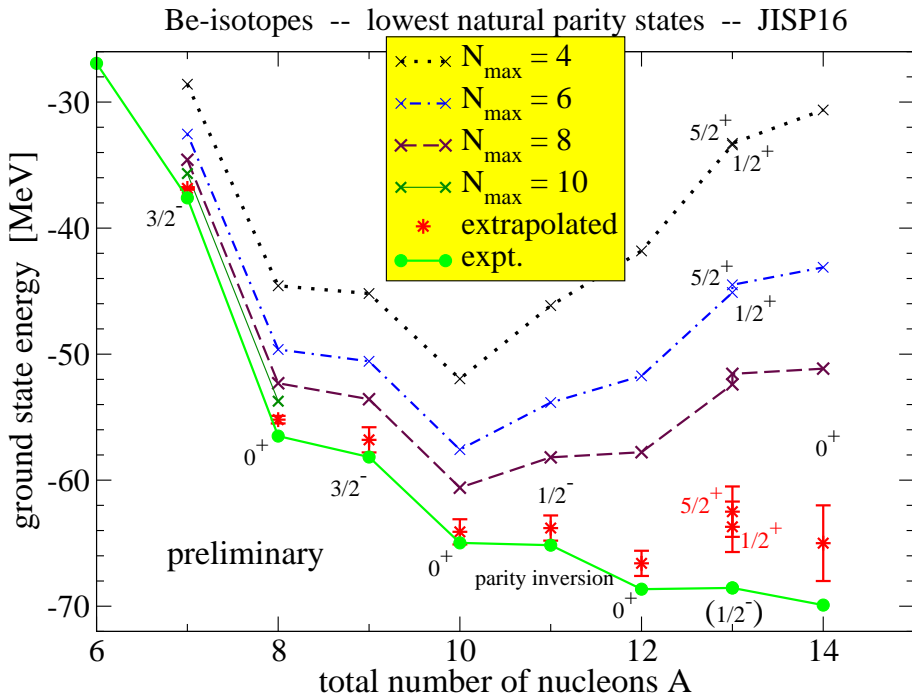
● Discrepancy between expt. and calc. μ
most likely due to two-body currents



Beryllium isotopes

updated from Vary, Maris, Ng, Yang, Sosonkina, arXiv:0907.0209 [nucl-th],

J. Phys. Conf. Ser. 180, 012083 (2009)

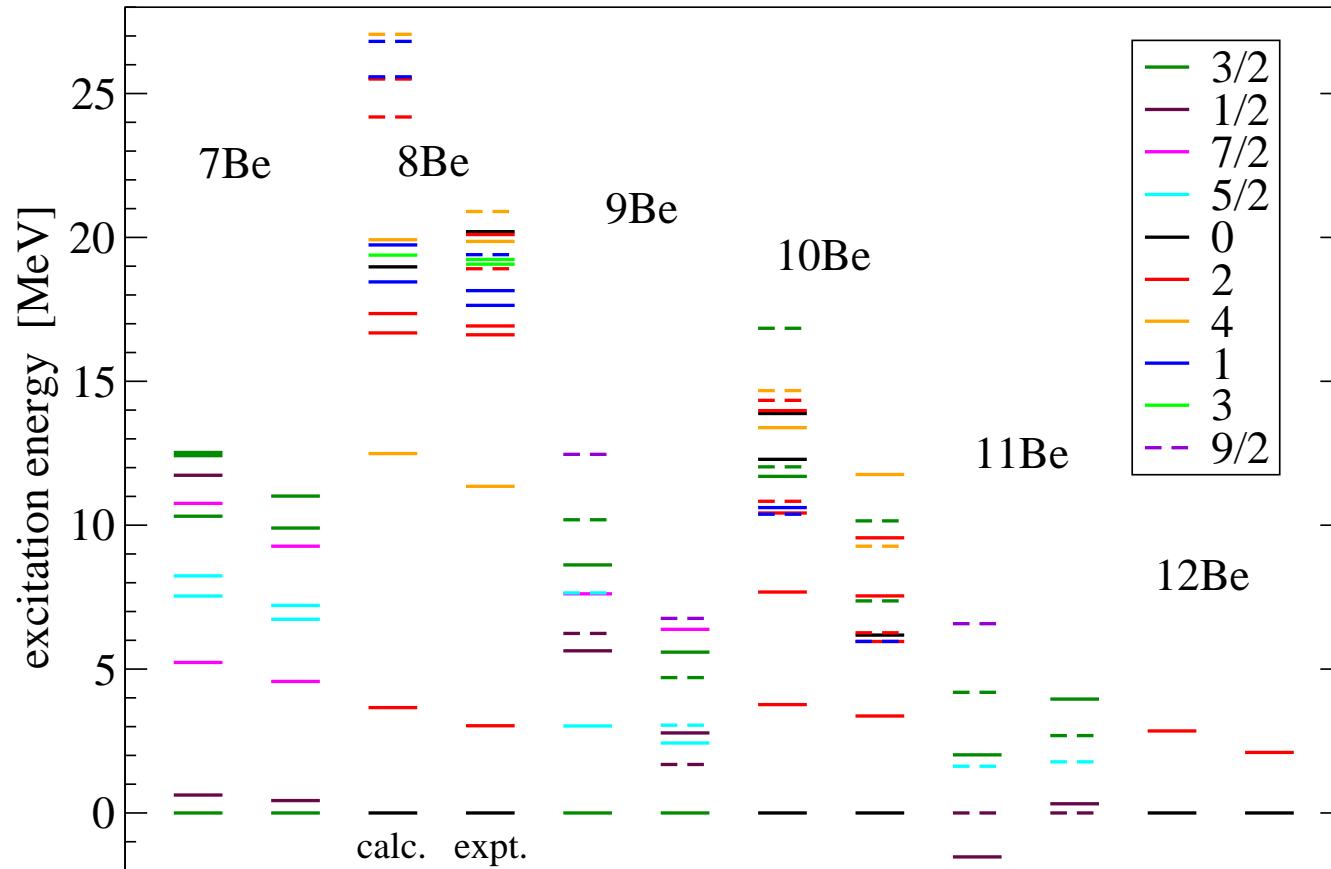


- Exploring physics near the neutron drip line – in progress
- Un-natural parity states systematically underbound with JISP16
- Similar results for He- and Li-isotopes

Results with JISP16 for Be spectra – work in progress

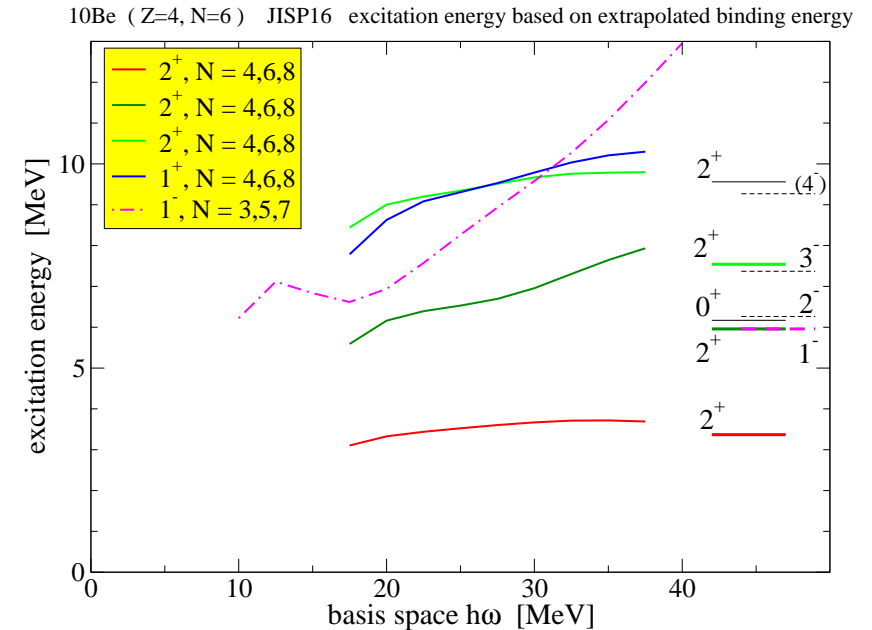
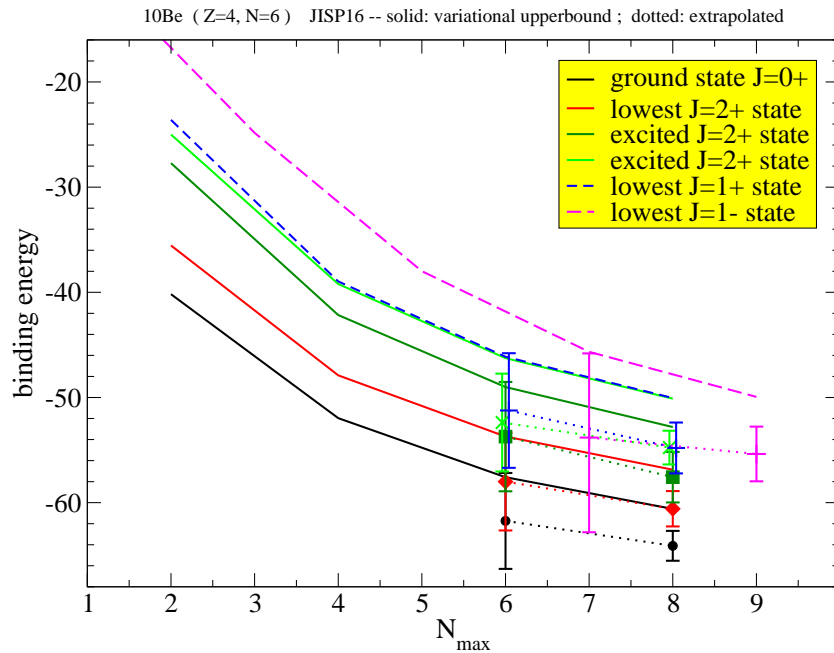
- Excitation energy at $\hbar\omega = 22.5$ MeV in largest model spaces (extrapolation to infinite model space in progress)

solid: natural parity dashed: un-natural parity



- Un-natural parity states systematically underbound with JISP16

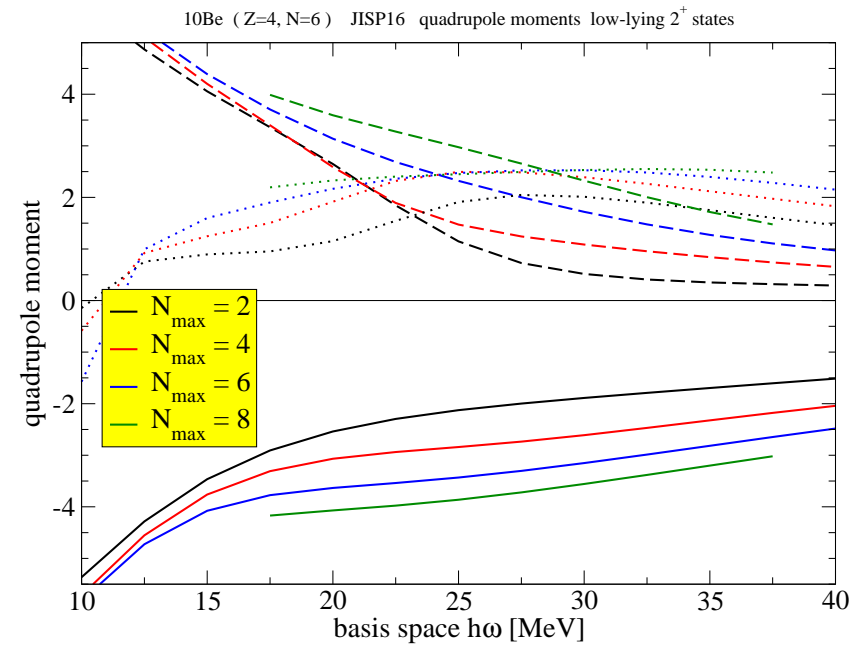
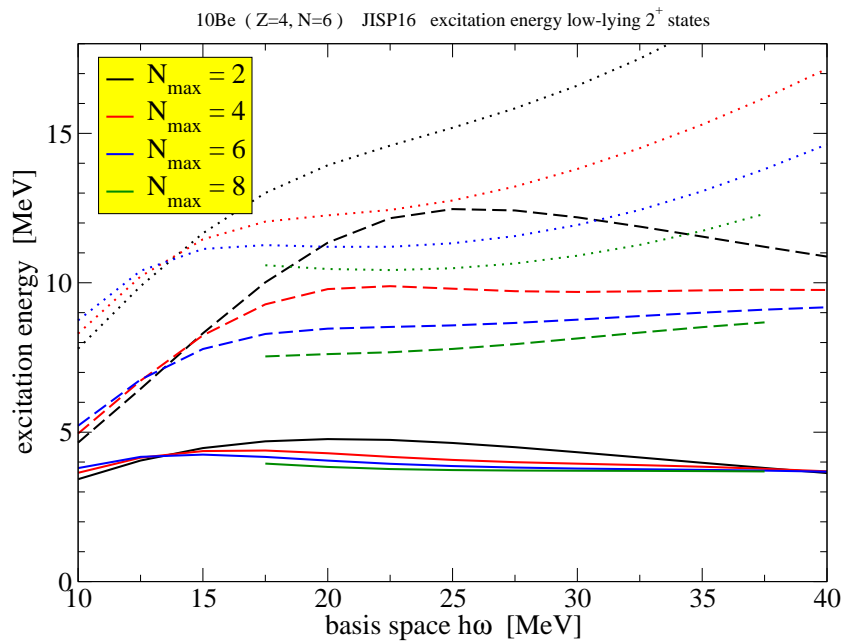
Results with JISP16 for ^{10}Be – work in progress



- Ground state energy in good agreement with data
64.11 \pm 1 MeV vs. 64.977 MeV
- First excited state converged and in good agreement with data
- Higher excited states and neg. parity states not converged yet
- Extrapolated spectrum in reasonable agreement with data

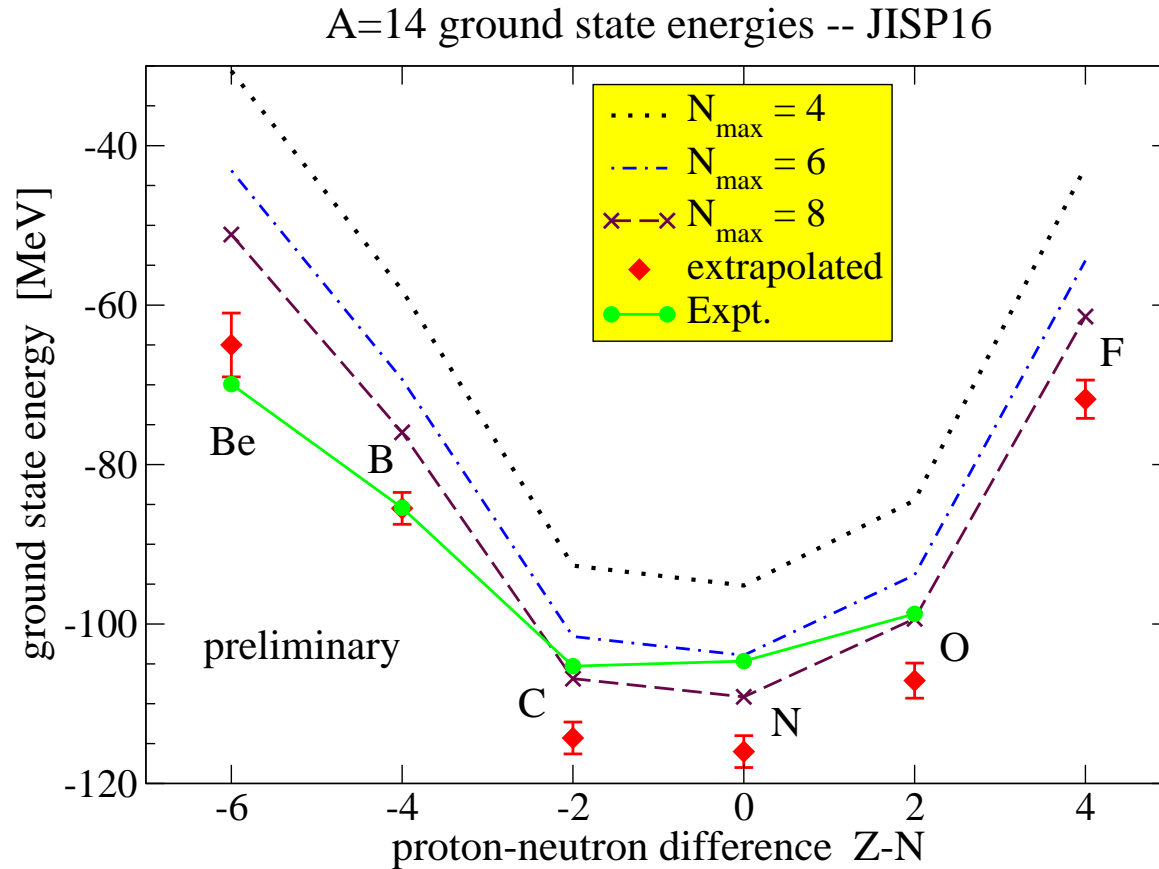
Quadrupole moments of ^{10}Be with JISP16 – work in progress

- Motivated by GFMC results
 - without 3-body forces: Q of lowest 2^+ state is positive but with 3-body forces: Q of lowest 2^+ state is negative and vice-versa for first excited 2^+ state



- JISP16 agrees qualitatively with AV18 plus 3-body forces

Valley of stability for $A = 14$ Nuclei



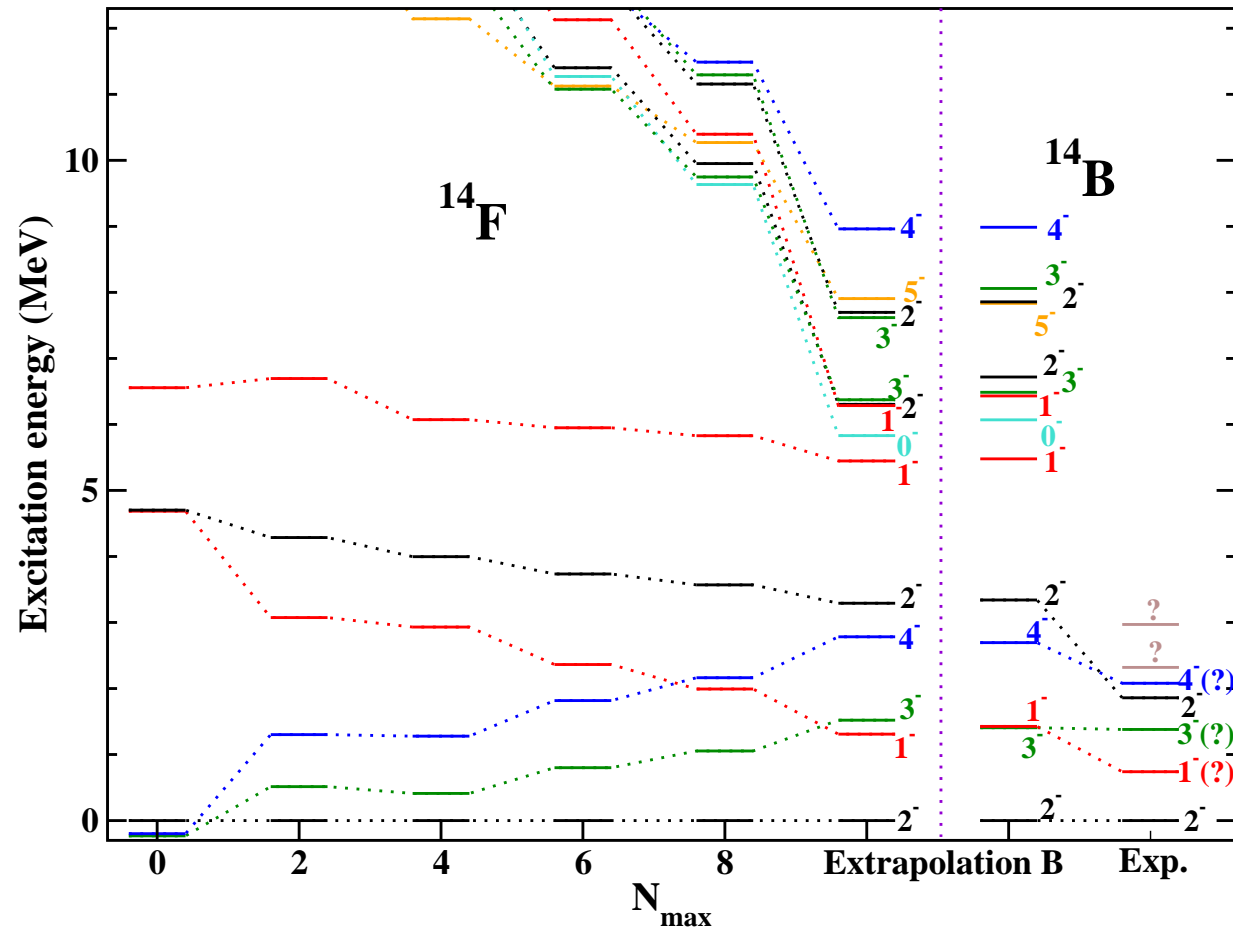
● Preliminary results

- up to 10% overbinding for ^{14}C , ^{14}N , ^{14}O

Maris, Vary, Shirokov, arXiv:0808.3420 [nucl-th], Phys. Rev. C79, 014308 (2009)

● Need improved (charge-dependent?) interaction

Spectrum of ^{14}F – submitted for SciDAC Highlights



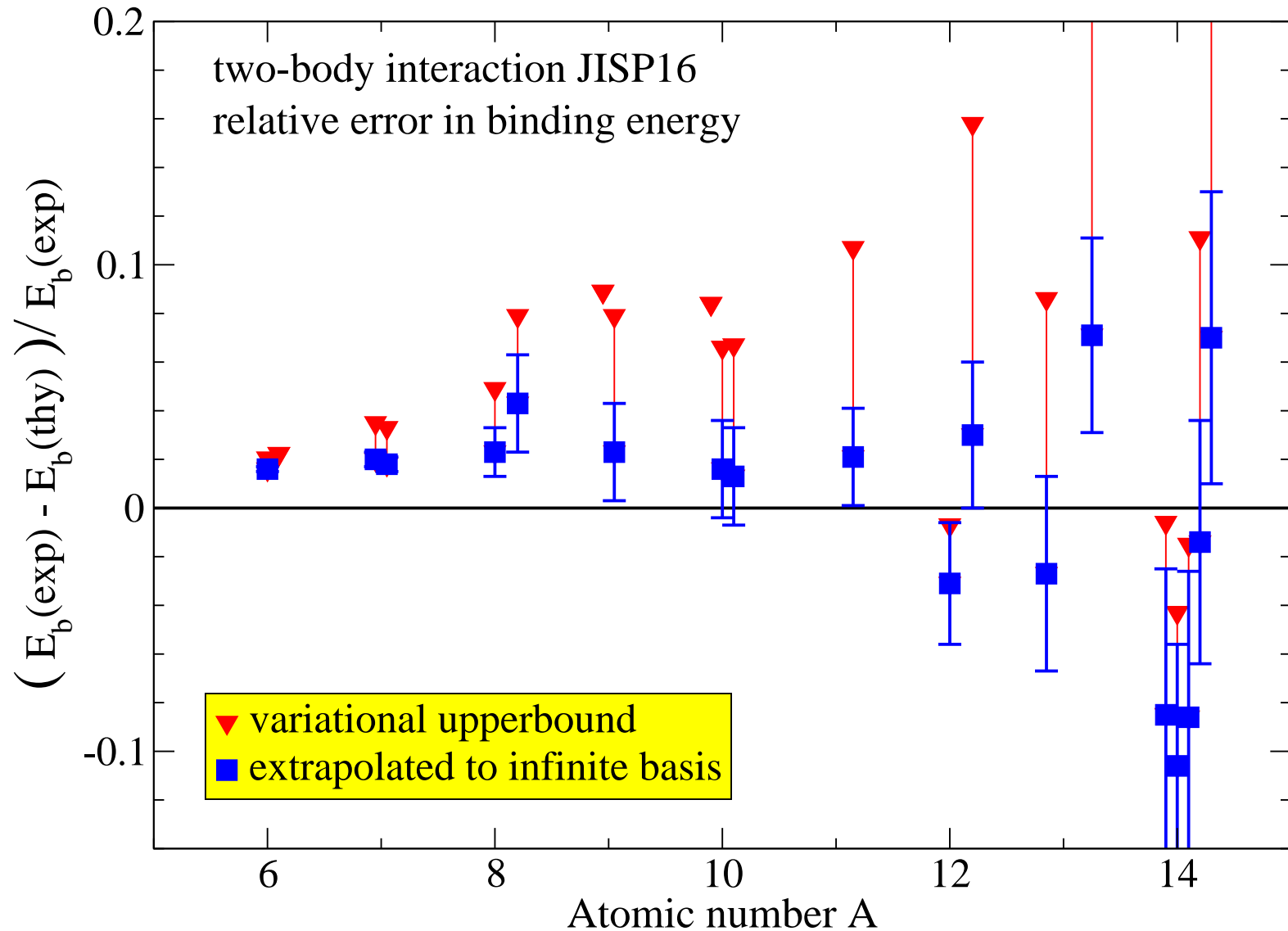
dimension $2 \cdot 10^9$
 # nonzero m.e. $2 \cdot 10^{12}$
 runtime 2 to 3 hours on
 7,626 quad-core nodes
 on Jaguar (XT4)
 (summer 2009)

Maris, Shirokov, Vary,
 arXiv:0911.2281 [nucl-th]
 Phys. Rev. C81, 021301(R) (2010)

- Predicted ground state energy: 72 ± 4 MeV (unstable)
- Mirror nucleus ^{14}B : 86 ± 4 MeV agrees with experiment 85.423 MeV
- Experiments in progress at Texas A&M

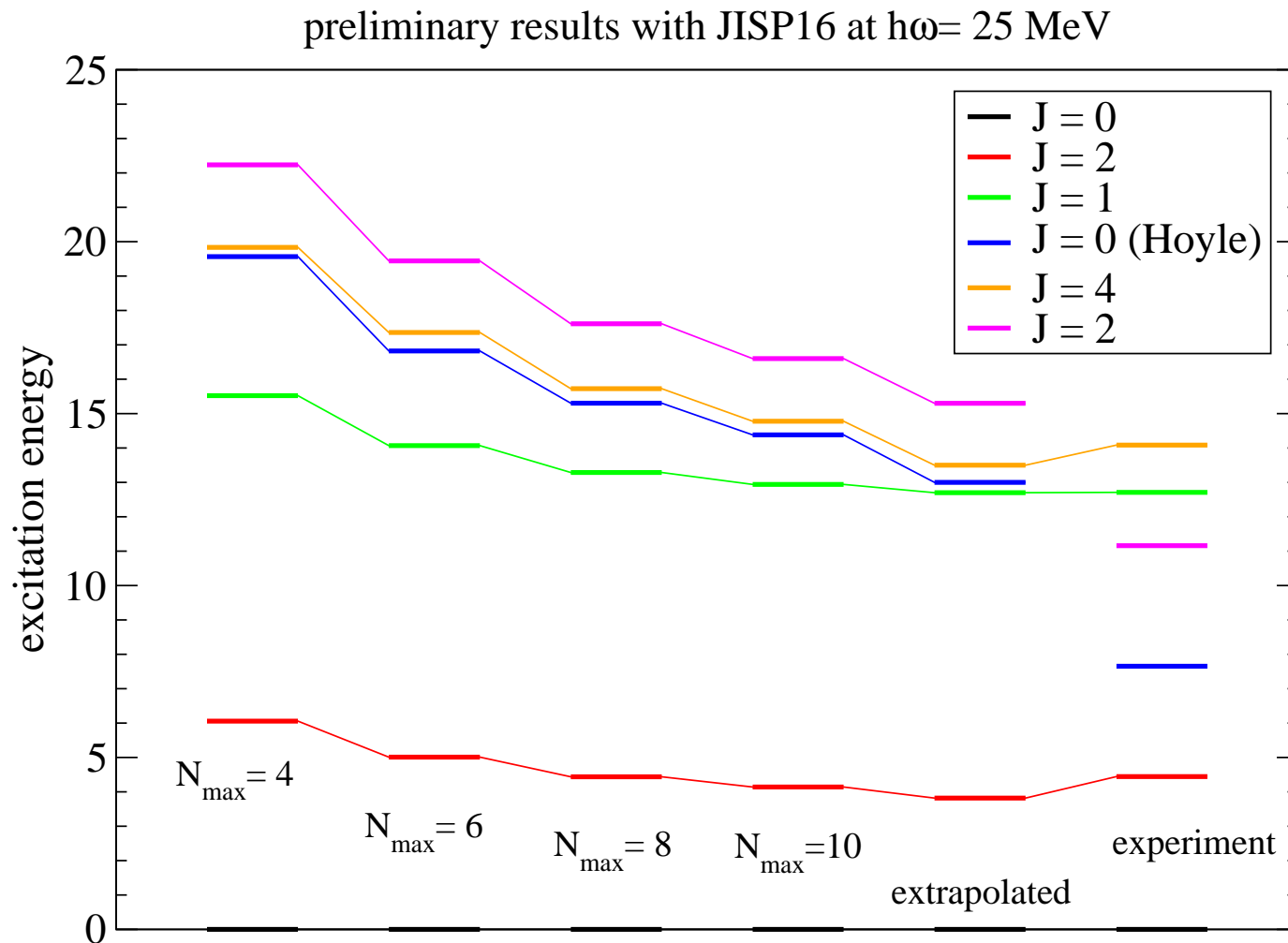
Summary of NCFC results with JISP16

Comparison of data and theory for natural parity ground state energies



12C spectrum with JISP16 – work in progress

- Excitation energy at $\hbar\omega = 25$ MeV up through $N = 10$



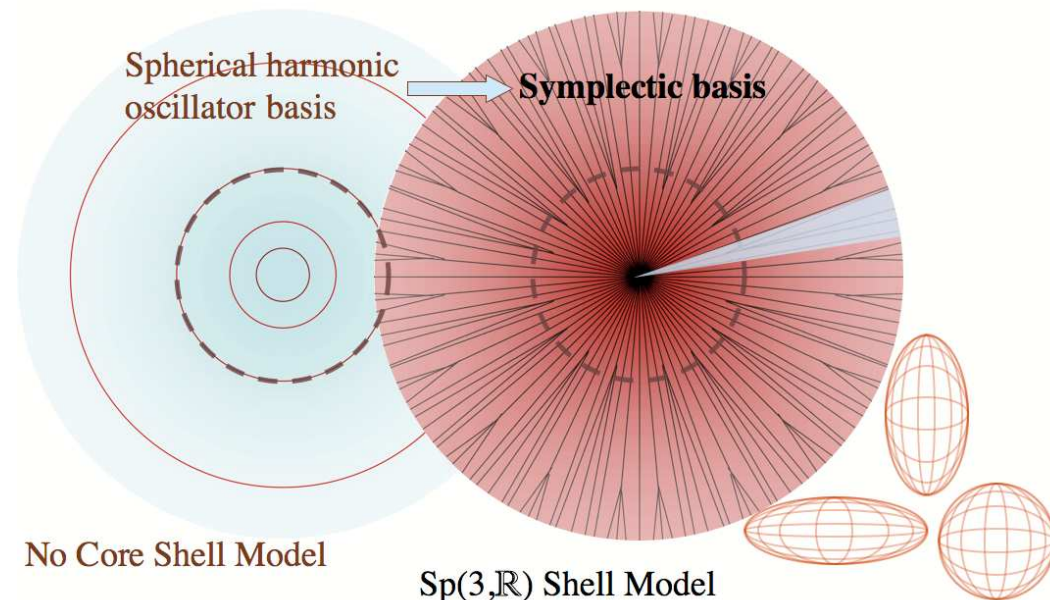
- Hoyle state not converged and/or far from experiment w. JISP16

Taming the scale explosion in nuclear structure calculations

PetaApps award (2009)

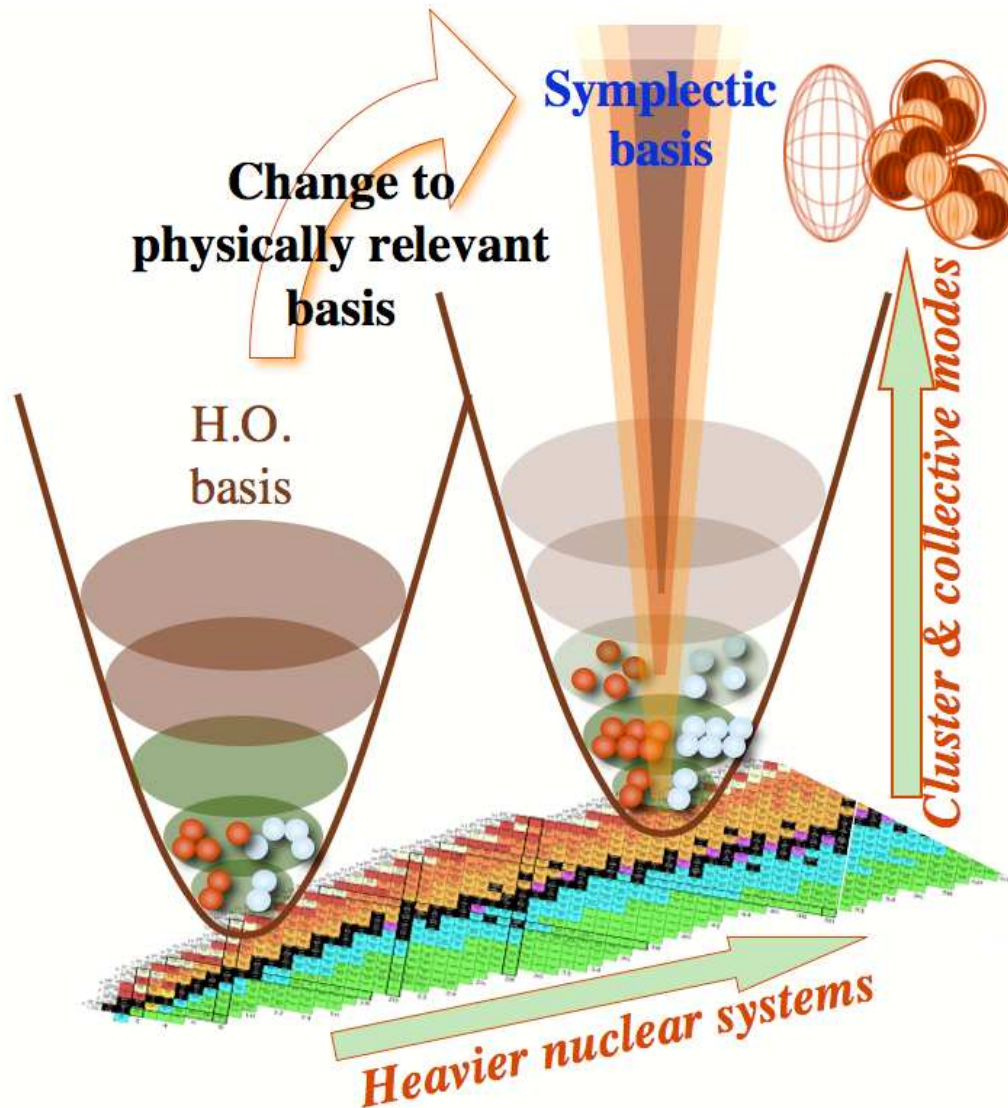
PIs: Jerry Draayer (LSU), Umit Catalyurek (OSU),
Masha Sosonkina, James Vary (ISU)

- Harmonic Oscillator basis: spherically symmetric
- Symplectic $Sp(3, \mathbb{R})$ basis: exploit symmetries of nuclear physics
 - highly deformed states (oblong, oblate)
 - α -cluster states
- Unitary transformation from H.O. to $Sp(3, \mathbb{R})$ basis: reducing the exponential increase of the basis space dimension with N_{\max} while incorporating important physics associated with clustering and deformations



Taming the scale explosion in nuclear structure calculations

PetaApps award (2009)



- Allows for ab initio calculations of
 - cluster states
 - deformed nuclei
 - nuclei in sd -shell (beyond ^{16}O)
- Requires innovative loadbalancing techniques
- Astrophysical applications: Hoyle state in ^{12}C (3 α -cluster state) crucial for nucleosynthesis

Conclusions

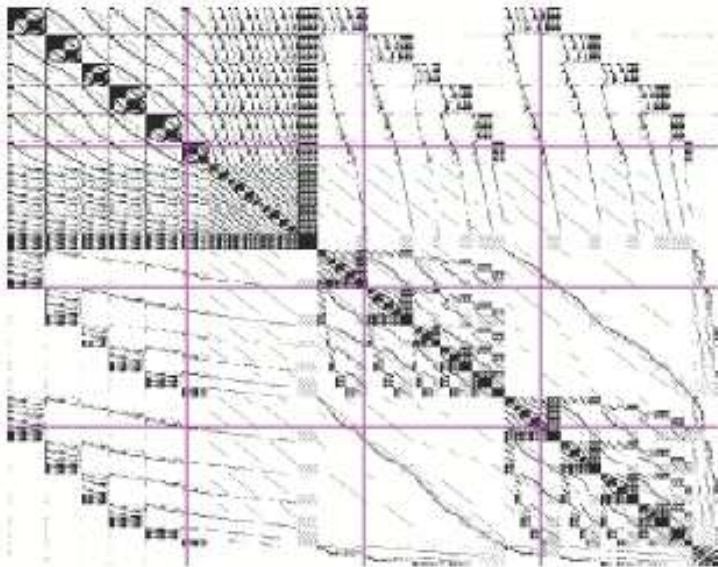
- **No-core Full Configuration** nuclear structure calculations
 - **Binding energy**, spectrum
 - $\langle r^2 \rangle$, μ , Q , transitions, wfns, one-body densities
- **Numerical convergence**
 - **Extrapolation** to infinite basis space for (absolute) energies
 - Other observables:
convergence rate depends on observable
- **JISP16**
 - Nonlocal phenomenological 2-body interaction
 - Good description of a range of light nuclei
 - Rapid convergence
- **Chiral effective interactions**
 - True ab-initio – nuclear interaction obtained from QCD
 - 3-body forces necessary for description of light nuclei
 - Renormalization techniques necessary

and Outlook

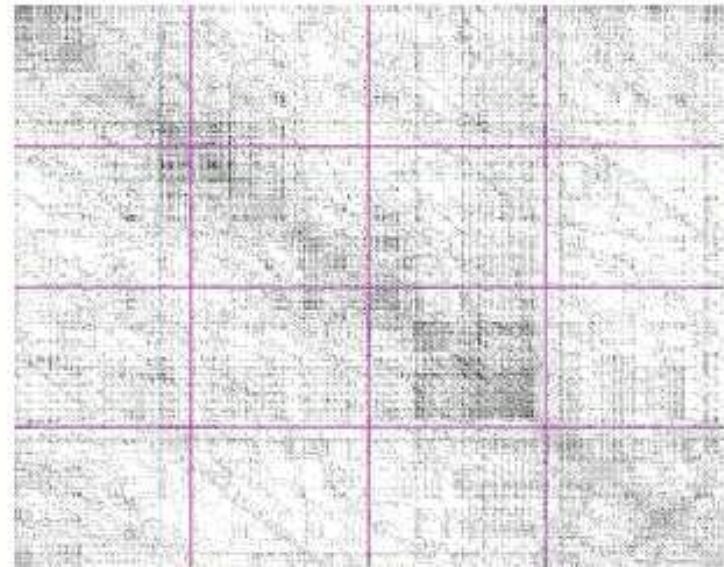
- Need realistic basis function to improve convergence $\langle r^2 \rangle$, Q Negoita, PhD work in progress
- Need improved (charge-dependent) nonlocal NN -interaction Shirokov, Mazur, work in progress
 - tuned to selected light nuclei and Nuclear Matter?
- Need SRG evolution of (OSU/LLNL), work in progress
 - chiral 3-body forces
 - operators
- Symplectic basis Draayer *et al.*, PetaApps award
 - significantly larger model spaces
- Quantum field theory Honkanen *et al.*
 - QED and QCD in lightfront dynamics

Many Fermion Dynamics – nuclear physics

- Platform-independent F90 code with both MPI and OpenMP
- Scalable (has run successfully on 45k+ cores on JaguarPF)
- Load-balanced



on single processor

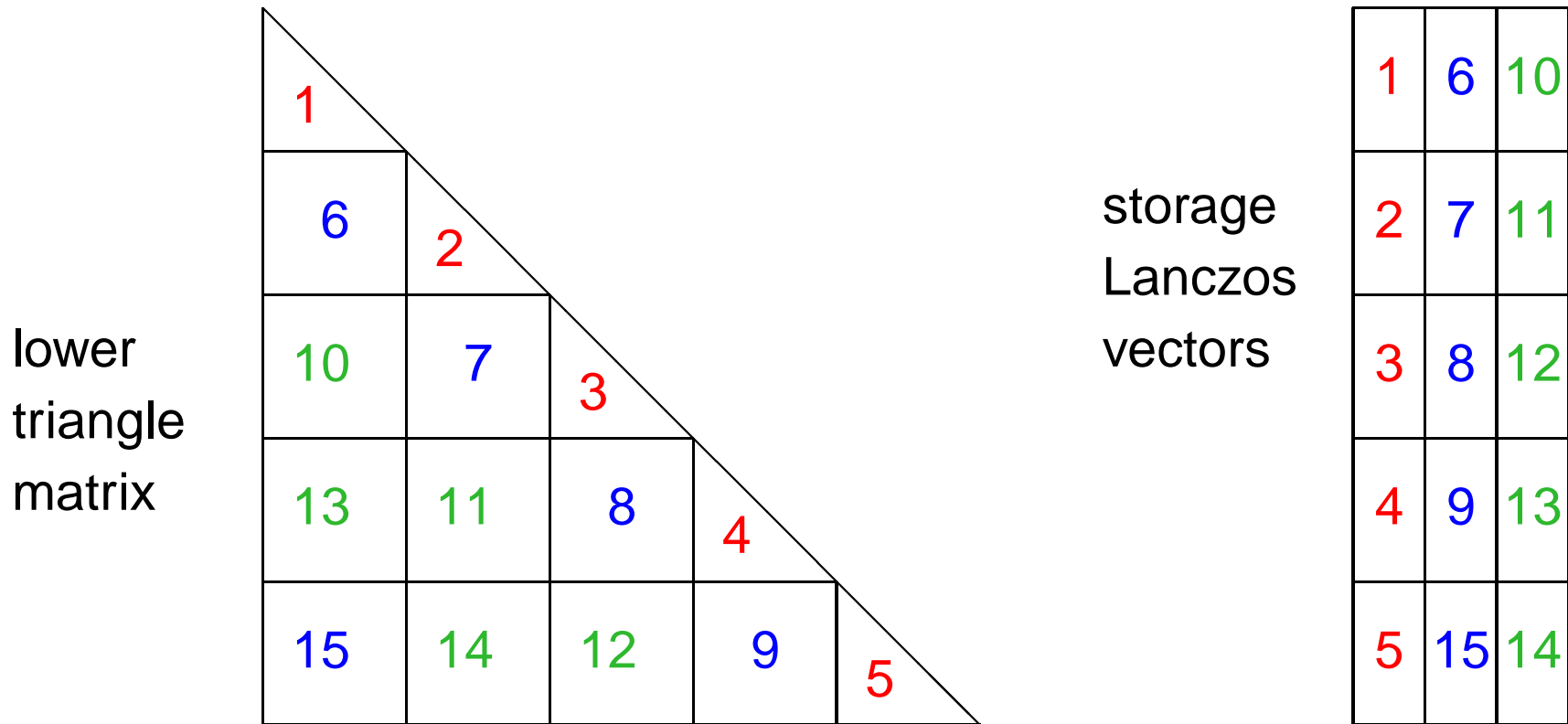


on 10 processors

- round-robin distribution of many-body states over d procs
- however, no (apparent) structure in sparse matrix

Distribution of matrix and vectors

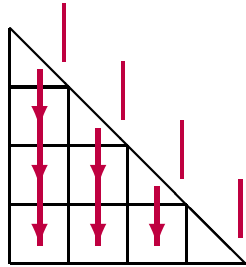
- Store lower half of symmetric matrix, distributed over $n = d \cdot (d + 1)/2$ processors with d “diagonal” proc’s
- Store Lanczos vectors on one of $(d + 1)/2$ groups of d proc’s



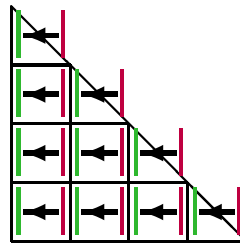
Distribution of matrix and vectors

- Store lower half of symmetric matrix, distributed over $n = d \cdot (d + 1)/2$ processors with d “diagonal” proc’s
- Store Lanczos vectors on one of $(d + 1)/2$ groups of d proc’s
- Communication pattern matrix-vector multiplication

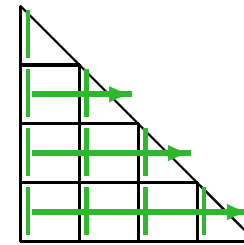
lower triangle



BCast(x)

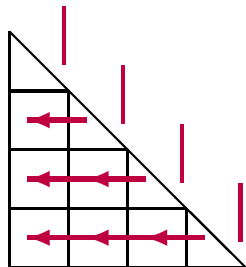


$y \leftarrow Ax$

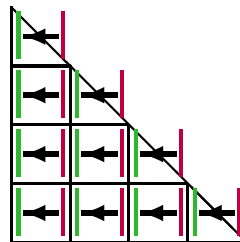


Reduce(y)

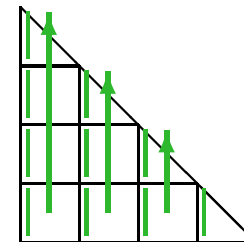
upper triangle



BCast(x)



$y \leftarrow A^T x$

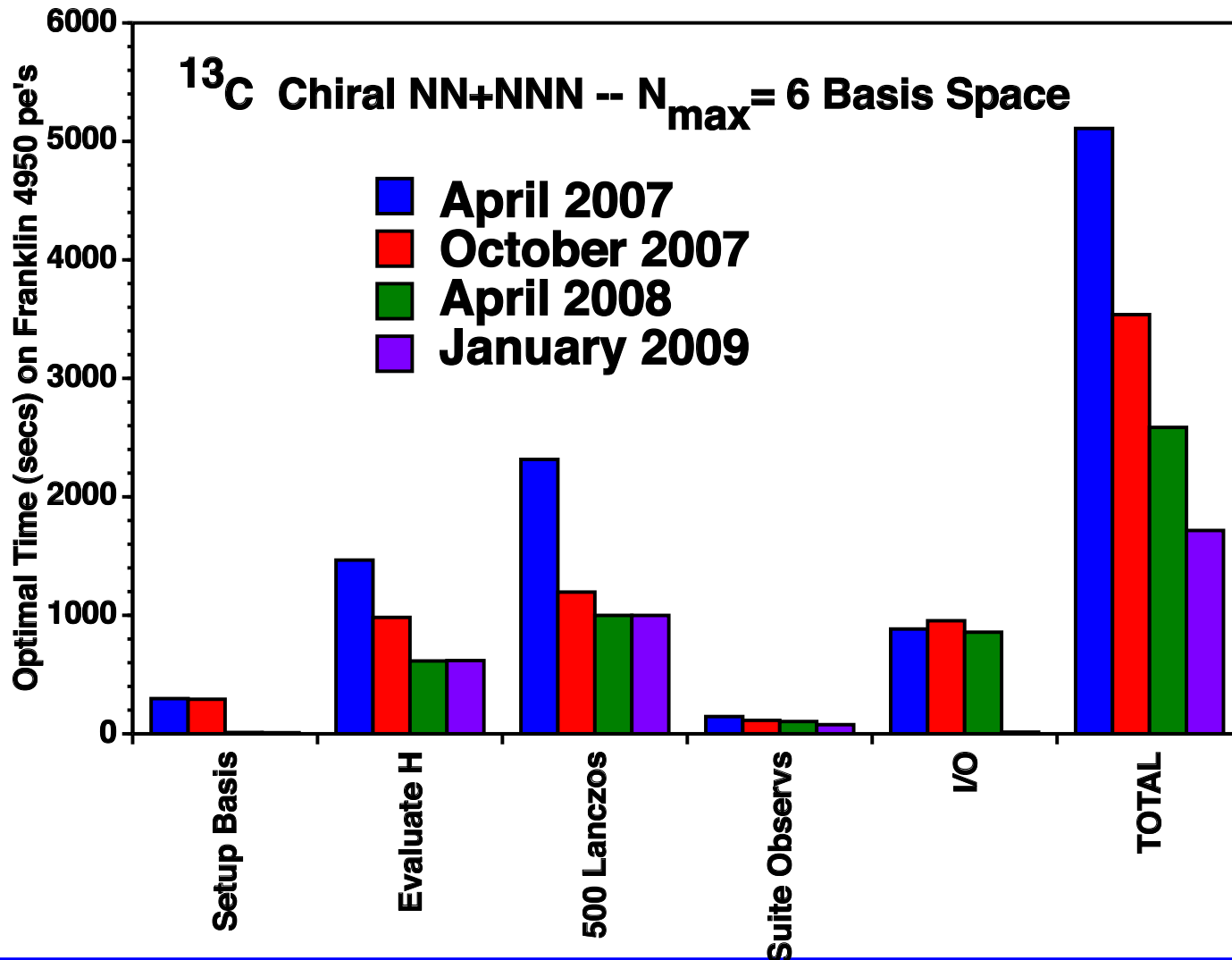


Reduce(y)

Recent performance improvements

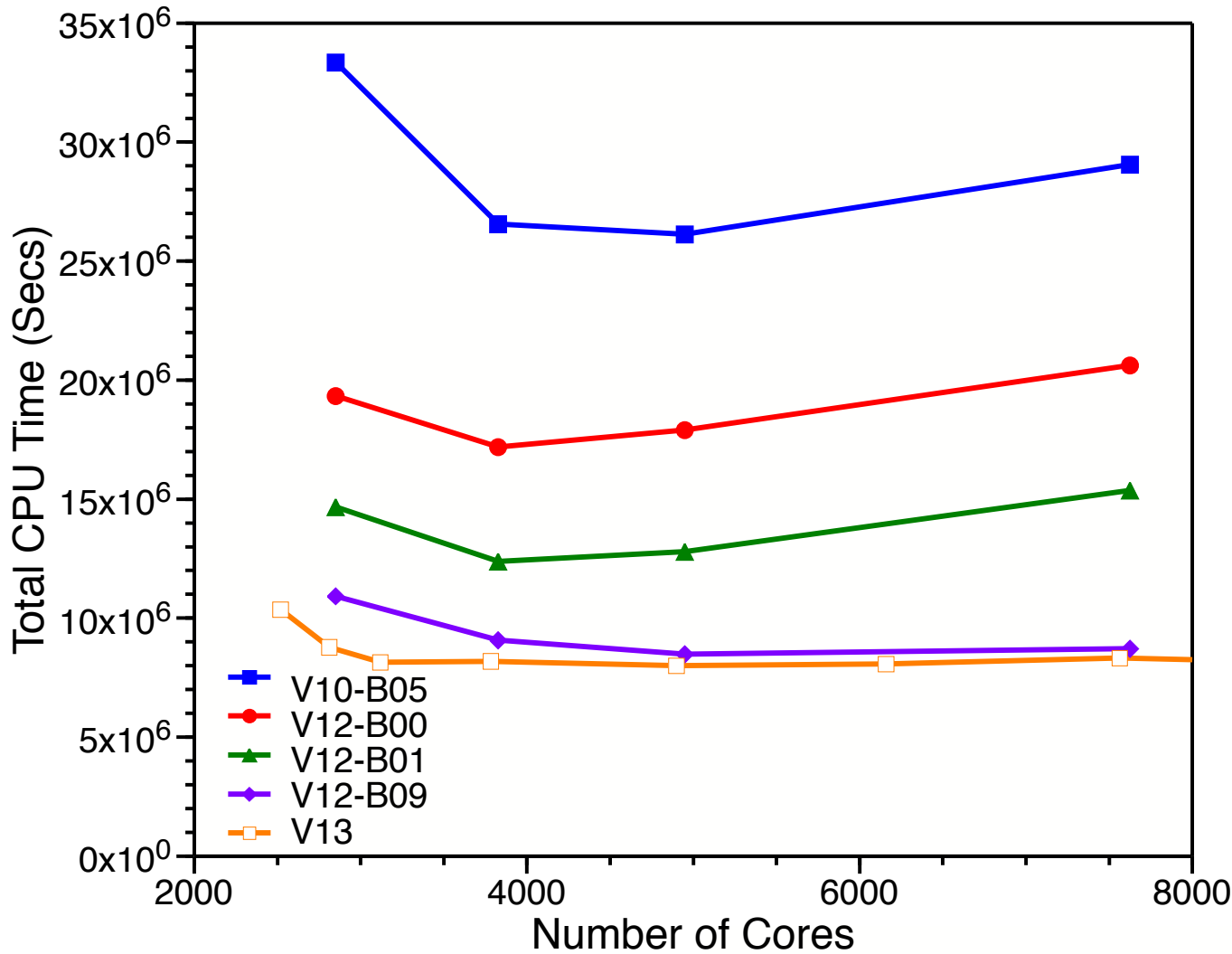
updated from Sternberg, Ng, Yang, Maris, Vary, Sosonkina, Le,

Accelerating Configuration Interaction calculations for nuclear structure, presented at SuperComputing08



dimension $38 \cdot 10^6$
nonzero m.e. $56 \cdot 10^{10}$
size input 3 GB

Strong Scaling on Cray XT4



^{13}C chiral N3LO
2- and 3-body
interactions

dimension $38 \cdot 10^6$
nonzero m.e. $56 \cdot 10^{10}$
size input 3 GB

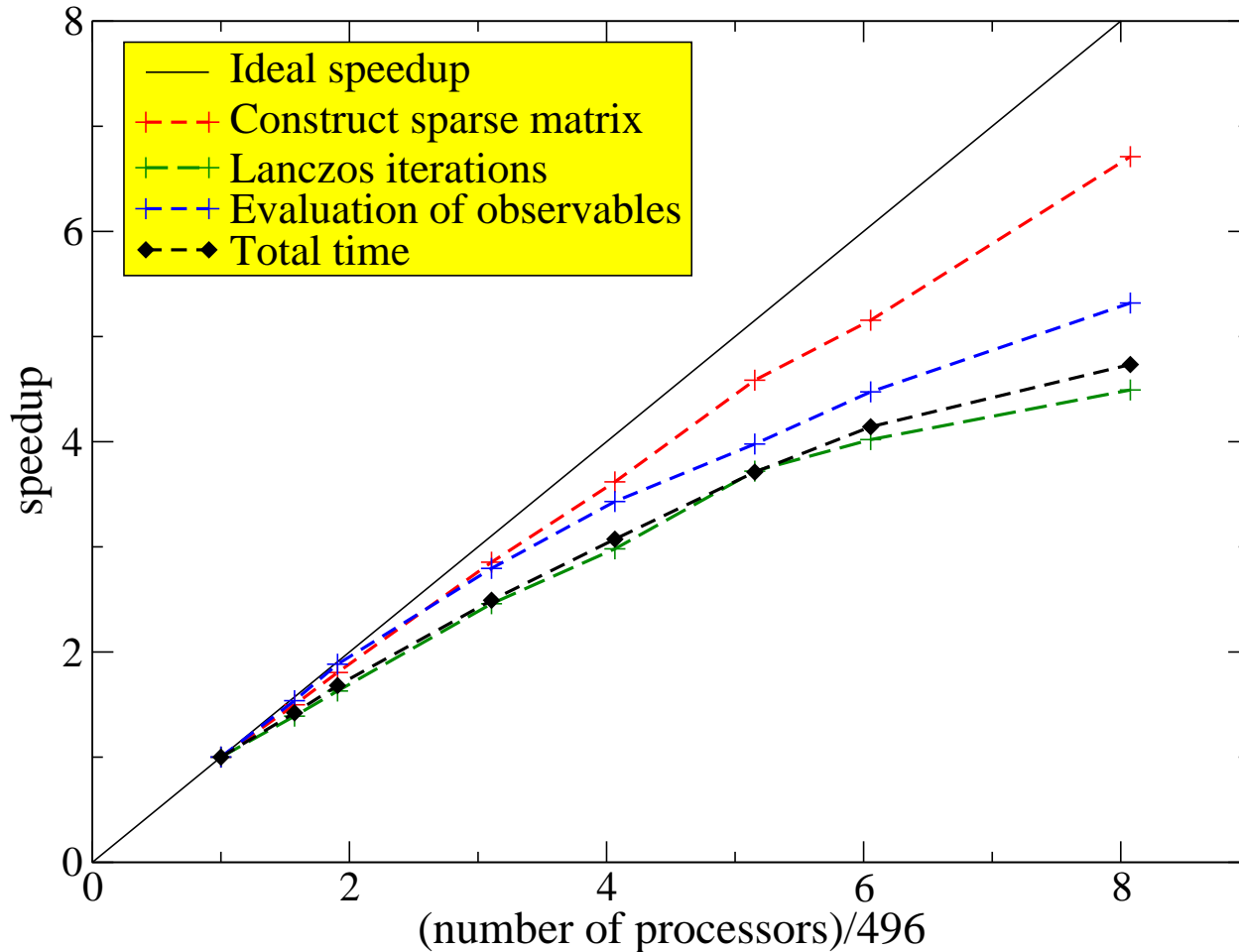
Version 13:
hybrid MPI
and OpenMP
(Aug. 2009)

Strong Scaling of MFDn on Cray XT4

adapted from Maris, Sosonkina, Vary, Ng, Yang,

submitted to ICCS 2010

Scaling of ab-initio nuclear physics calculations on multicore computer architectures

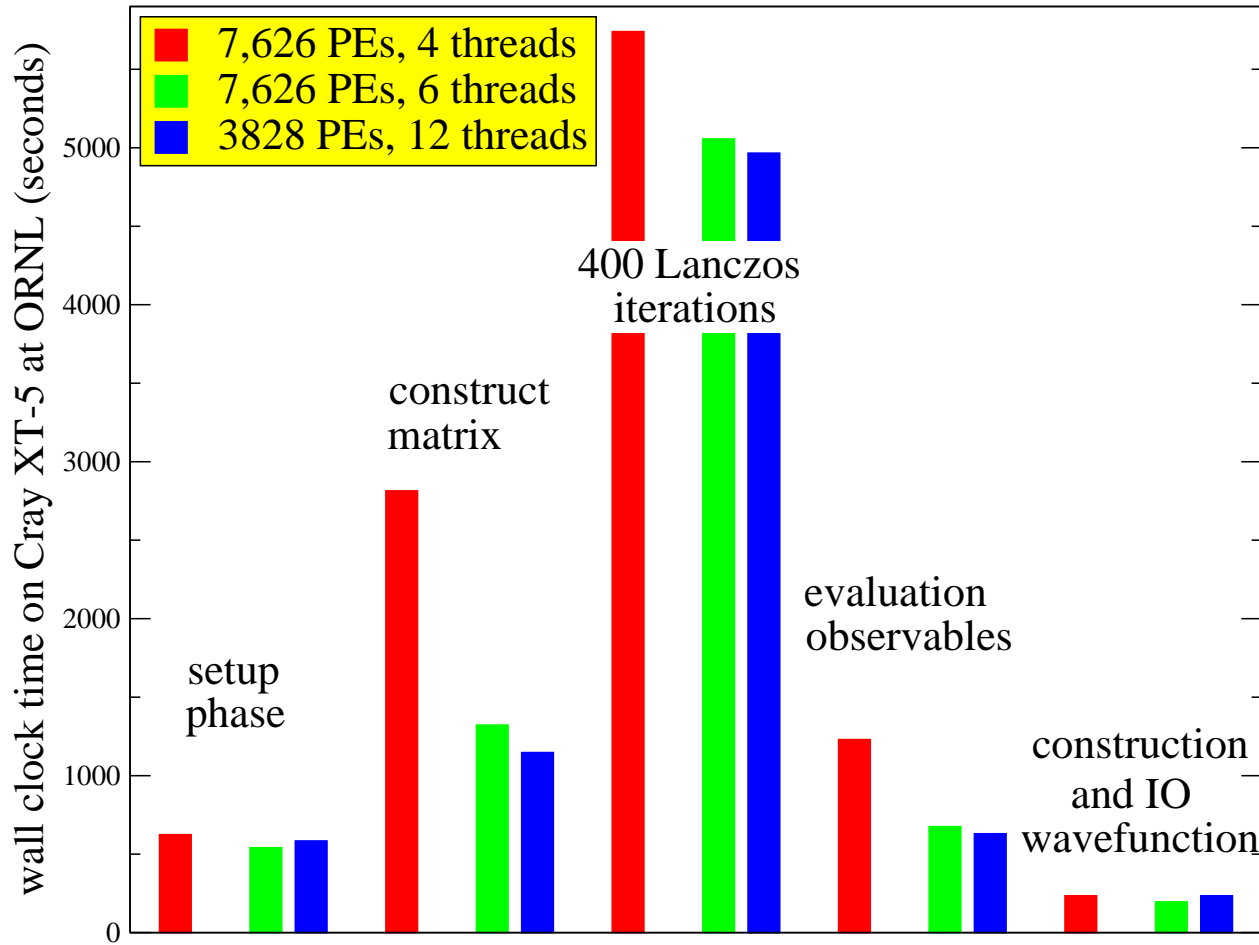


${}^6\text{Li}$, $N_{\text{max}} = 12$,
2-body interactions
strong scaling
from 496 PE's
to over 4,000 PE's

$$\text{speedup} = \frac{T_{496 \text{ PE's}}}{T_{\# \text{ PE's}}}$$

dimension $49 \cdot 10^6$
nonzero m.e. $74 \cdot 10^9$
memory for storing matrix:
600 GB

Quad core vs. Hex core performance on Cray XT5



^{14}Be
2-body interaction
dimension $2.8 \cdot 10^9$
nonzero m.e. $2.8 \cdot 10^{12}$