

Observables of the High-Density Equation of State in Supernovae

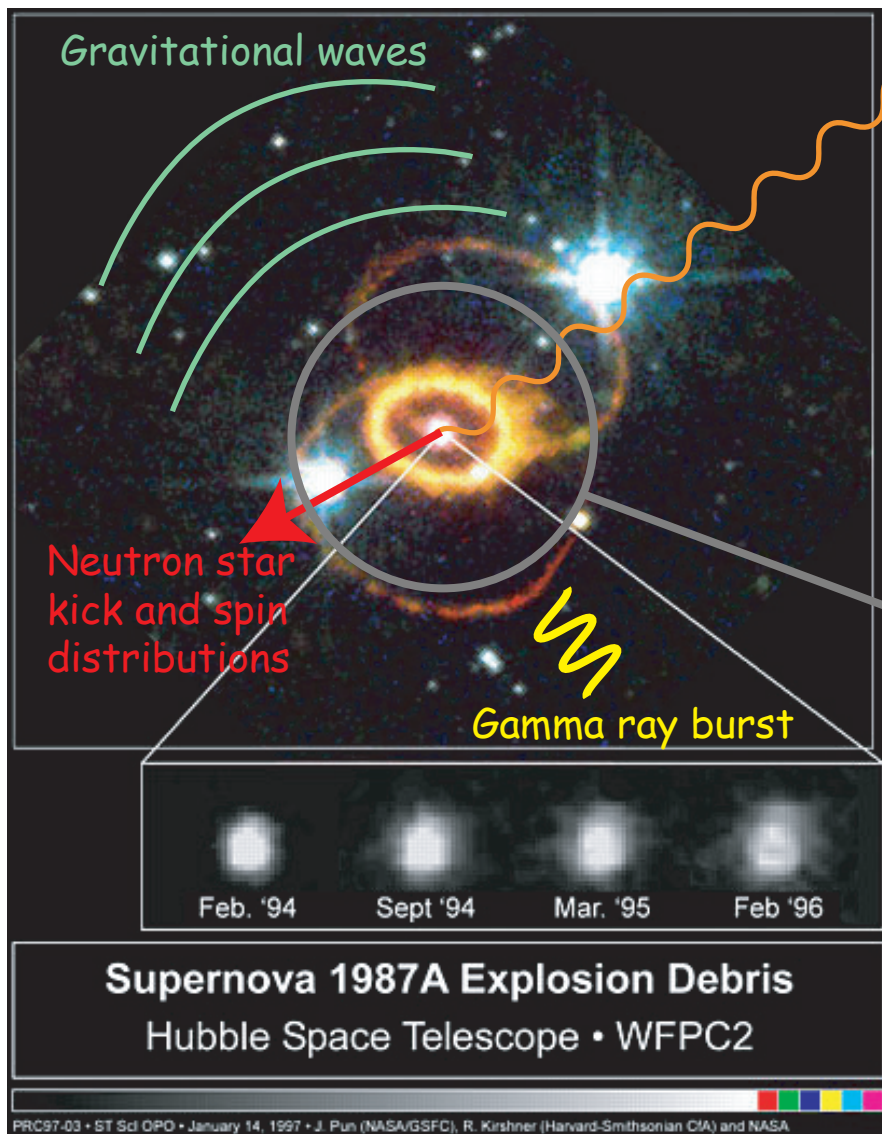
M. Liebendörfer
University of Basel

- The core-collapse supernova neutrino signature
- Imprints of the compressibility of the neutron star and an early QCD phase transition to quark matter
- Gravitational waves from stellar core collapse with different equation of state incompressibilities

with

- T. Fischer
- M. Hempel
- A. Mezzacappa
- G. Pagliara
- I. Sagert
- J. Schaffner-Bielich
- S. Scheidegger
- F.-K. Thielemann
- S. Whitehouse

Supernova Observables



neutrino signal
from interior

direct ejecta:
• composition
• velocity
(spectra)

• asymmetry
(polarization)

indirect ejecta
• mixing with ISM
• new star formation
• contamination of
metal-poor stars

Cosmology

Galactic evolution

Stellar evolution

Supernova
theory

Nuclear Physics
Hydrodynamics
Radiative transfer

Make extreme
conditions of matter
observable...

Core collapse supernova



JANUARY 15, 1934

PHYSICAL REVIEW

VOLUME 45

Proceedings
of the
American Physical Society

38. Supernovae and Cosmic Rays. W. BAADE, *Mt. Wilson Observatory*, AND F. ZWICKY, *California Institute of Technology*.—Supernovae flare up in every stellar system (nebula) once in several centuries. The lifetime of a supernova is about twenty days and its absolute brightness at maximum may be as high as $M_{\text{vis}} = -14^M$. The visible radiation L_v of a supernova is about 10^8 times the radiation of our sun, that is, $L_v = 3.78 \times 10^{41}$ ergs/sec. Calculations indicate that the total radiation, visible and invisible, is of the order $L_r = 10^7 L_v = 3.78 \times 10^{48}$ ergs/sec. The supernova therefore emits during its life a total energy $E_r \geq 10^8 L_r = 3.78 \times 10^{63}$ ergs. If supernovae initially are quite ordinary stars of mass $M < 10^{34}$ g, E_r/c^2 is of the same order as M itself. In the *supernova process mass in bulk is annihilated*. In addition the hypothesis suggests itself that *cosmic rays are produced by supernovae*. Assuming that in every nebula one supernova occurs every thousand years, the intensity of the cosmic rays to be observed on the earth should be of the order $\sigma = 2 \times 10^{-8}$ erg/cm² sec. The observational values are about $\sigma = 3 \times 10^{-8}$ erg/cm² sec. (Millikan, Regener). With all reserve we advance the view that supernovae represent the transitions from ordinary stars into *neutron stars*, which in their final stages consist of extremely closely packed neutrons.

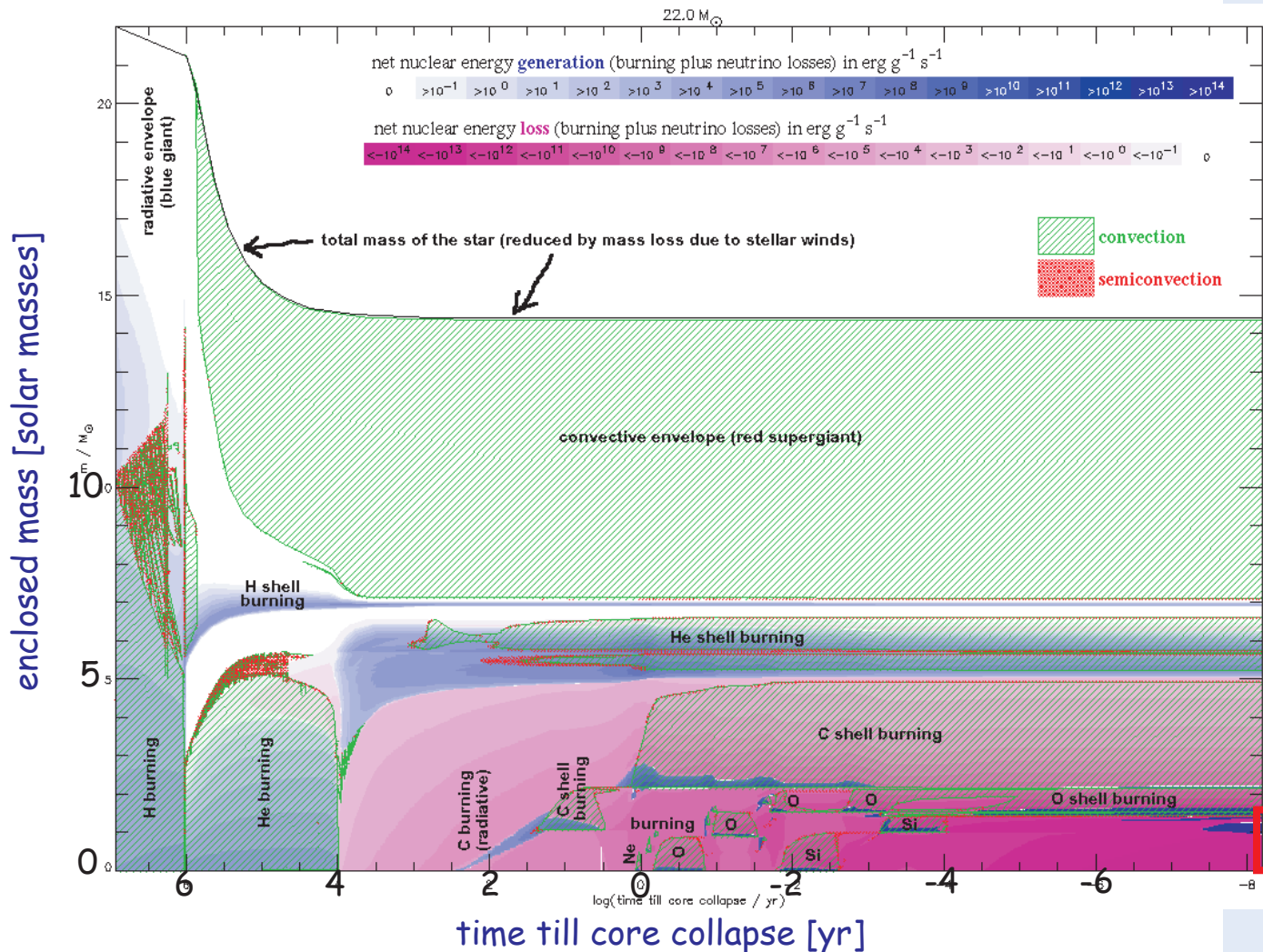
Huge Energies

- neutrinos:
~1e+53 erg
- mechanical:
~1e+51 erg
- electro-magn.:
~1e+48 erg elmag
- visible:
~1e+41 erg visible

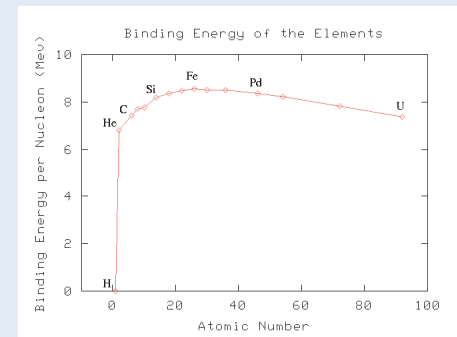
56Ni -> 56Co -> 56Fe
~6d ~110d

Iron core collapse

Overview of burning phases in stellar evolution



- Fusion in core reaches maximum binding energy per baryon



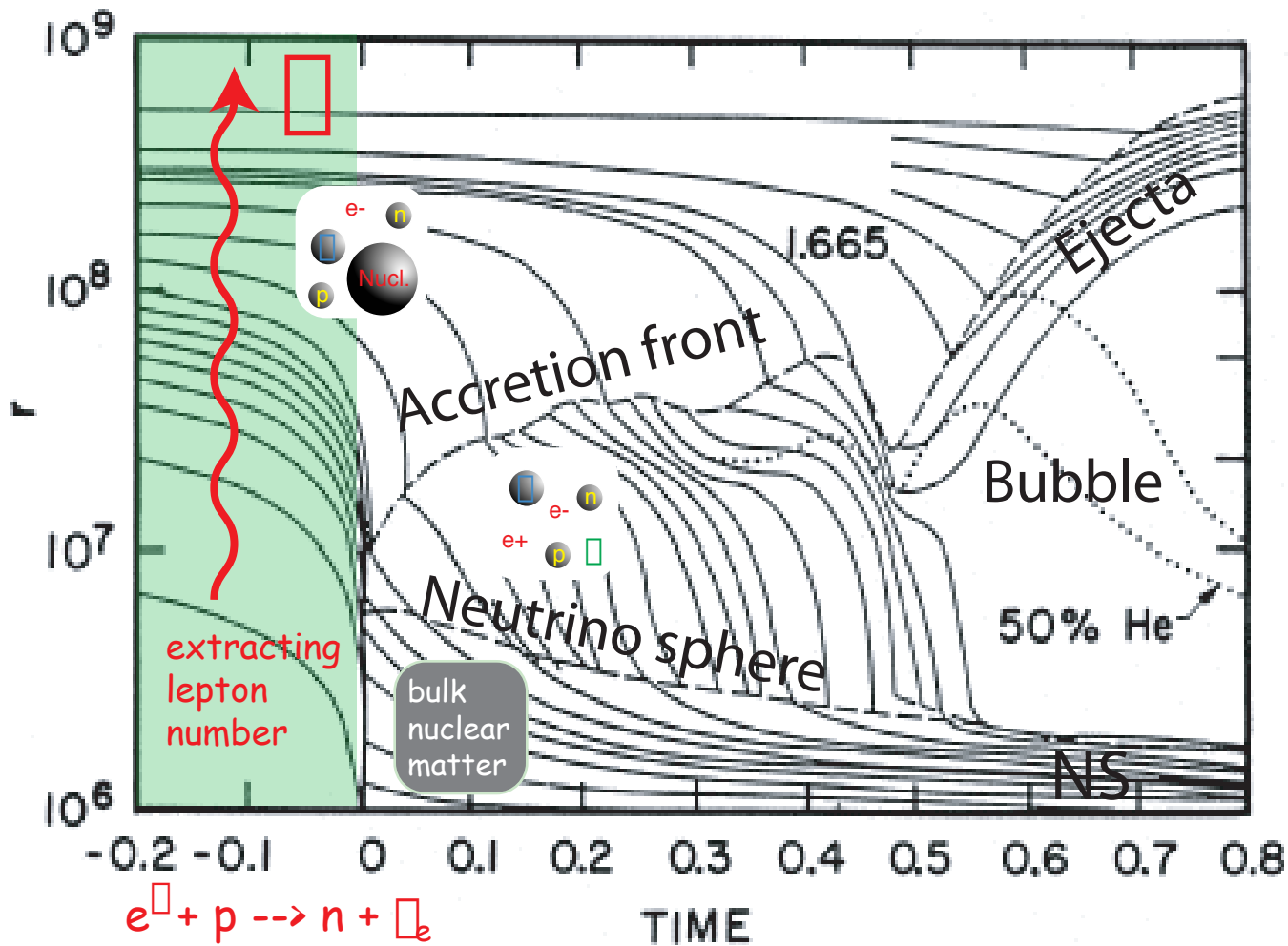
- There is a maximum stable mass: Chandrasekhar mass

stellar core collapse
 <-- happens here!

(Heger & Woosley 2002, see also Hirschi, Meynet, Maeder 2005)

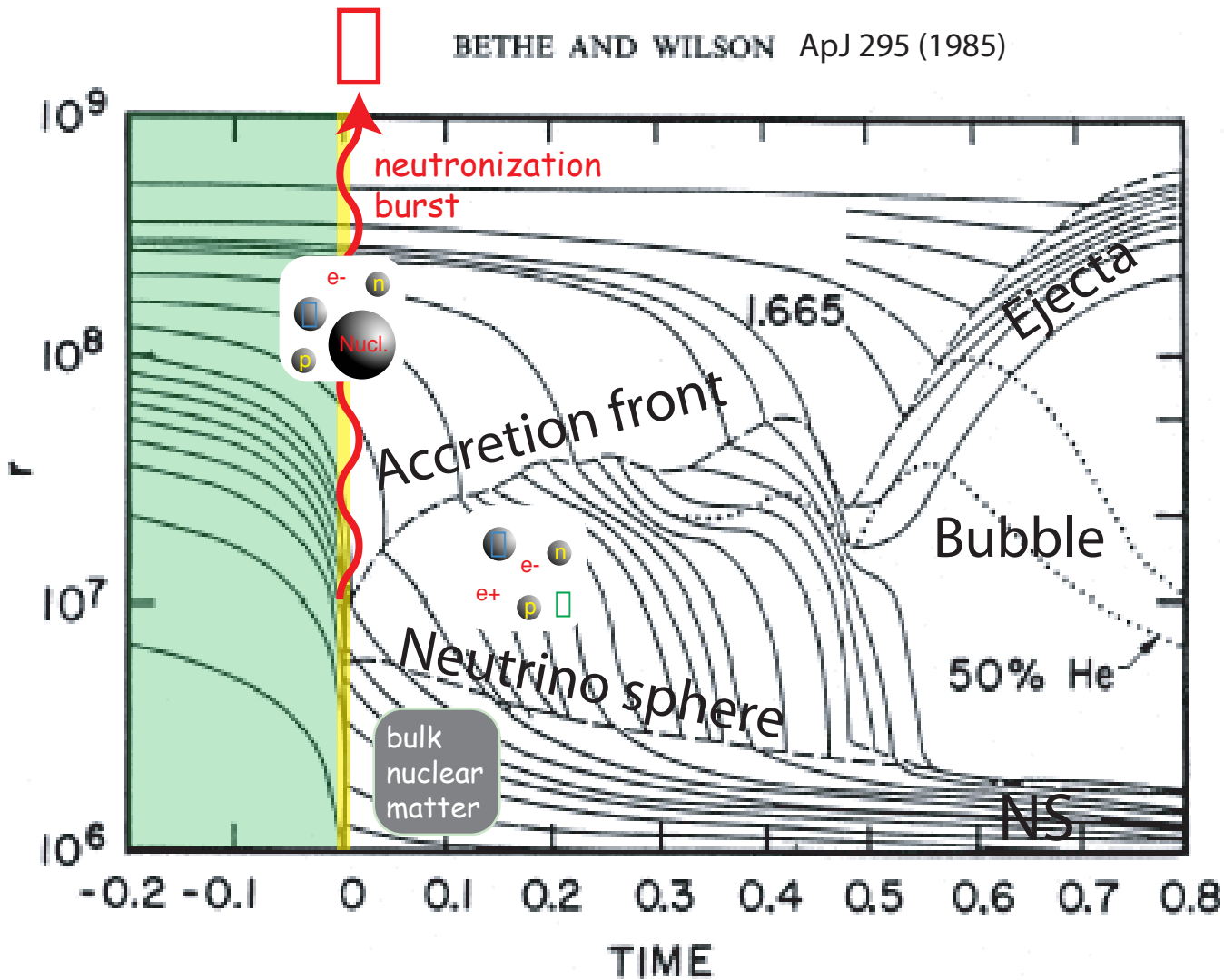
Delayed explosion: 4 phases

BETHE AND WILSON ApJ 295 (1985)



1) Collapse

Delayed explosion: 4 phases

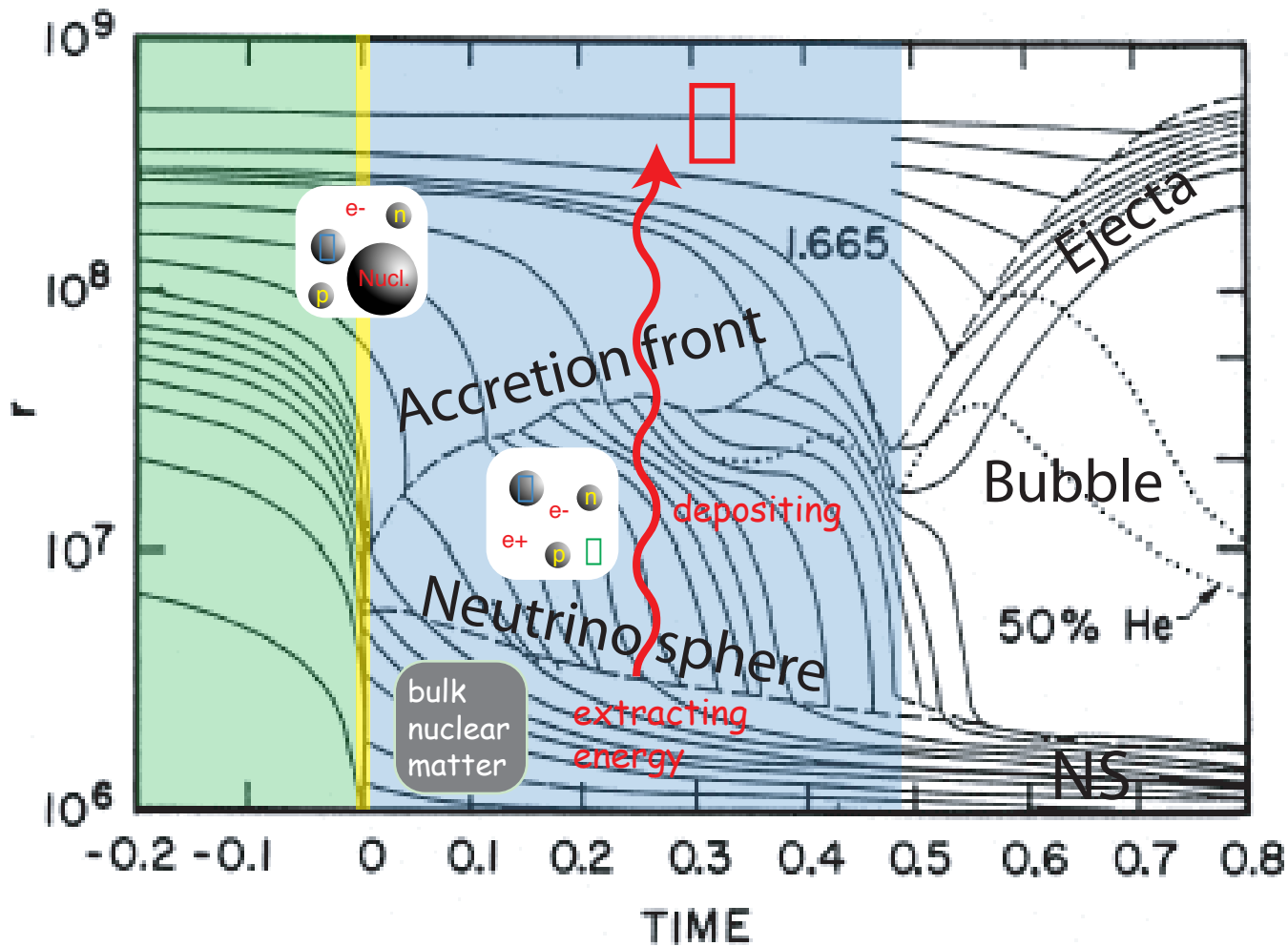


1) Collapse

2) Bounce

Delayed explosion: 4 phases

BETHE AND WILSON ApJ 295 (1985)



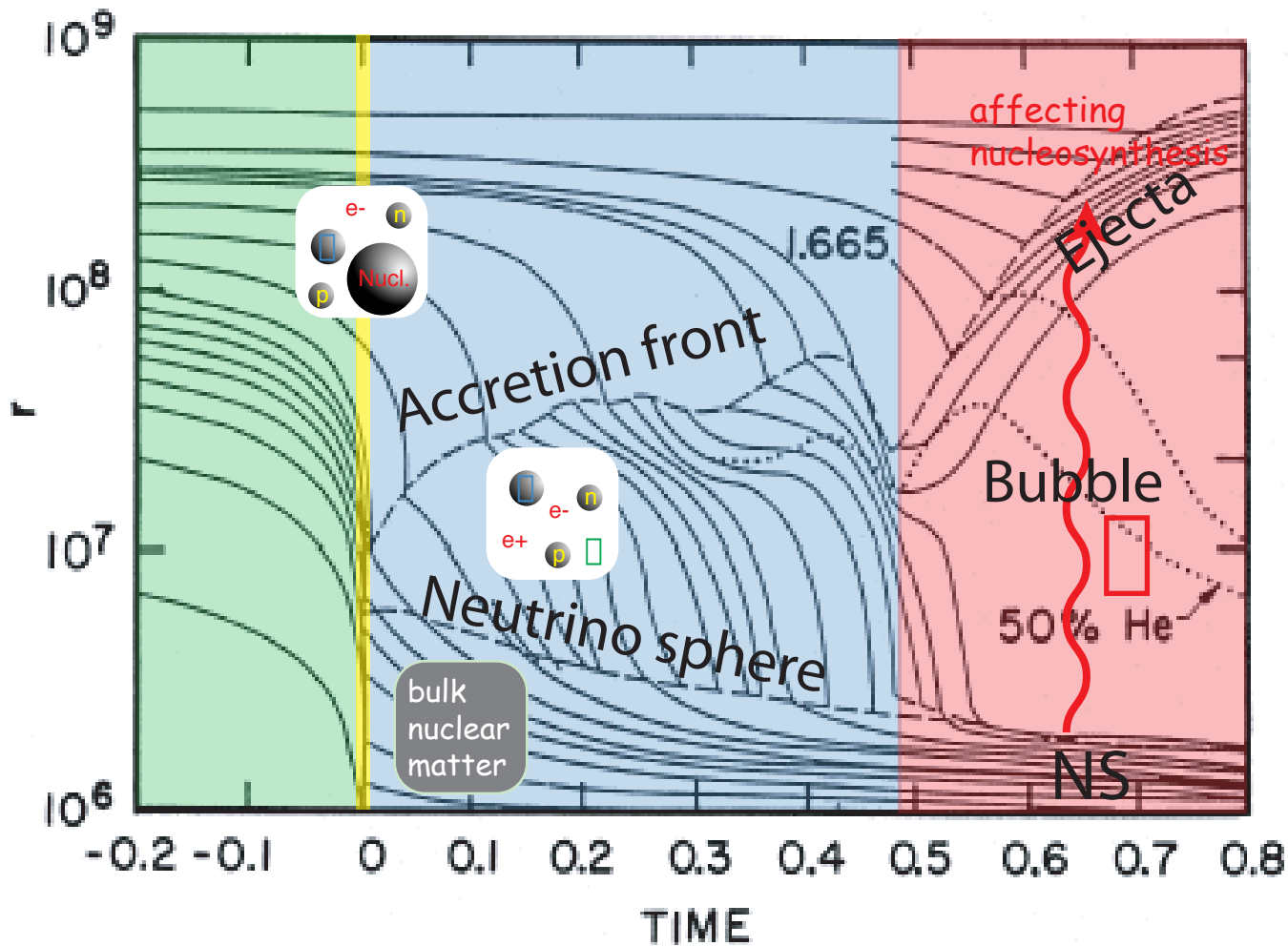
1) Collapse

2) Bounce

3) Accretion

Delayed explosion: 4 phases

BETHE AND WILSON ApJ 295 (1985)



Colgate & White, ApJ 143 (1966)

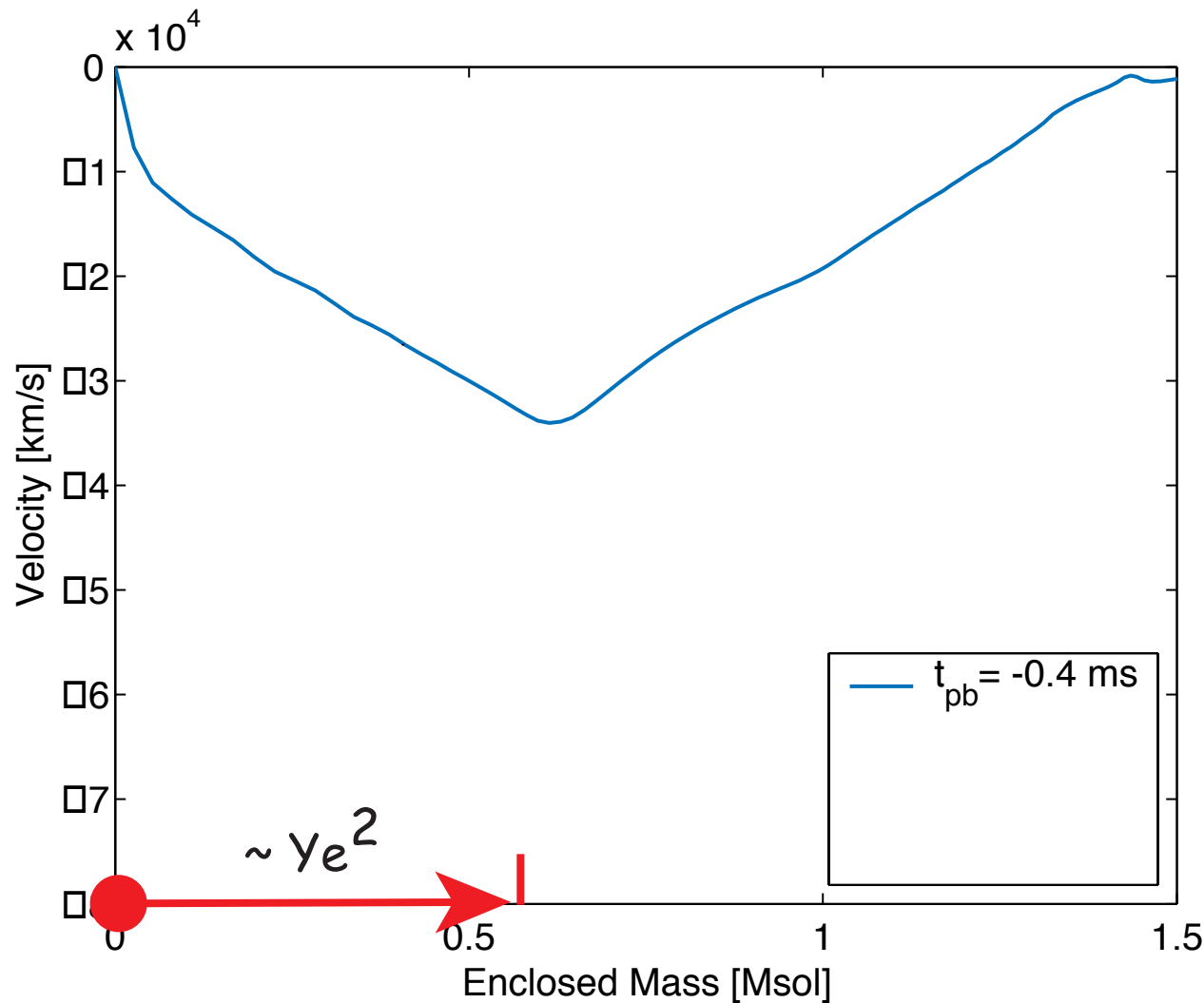
1) Collapse

2) Bounce

3) Accretion

4) Explosion

Neutronization during collapse



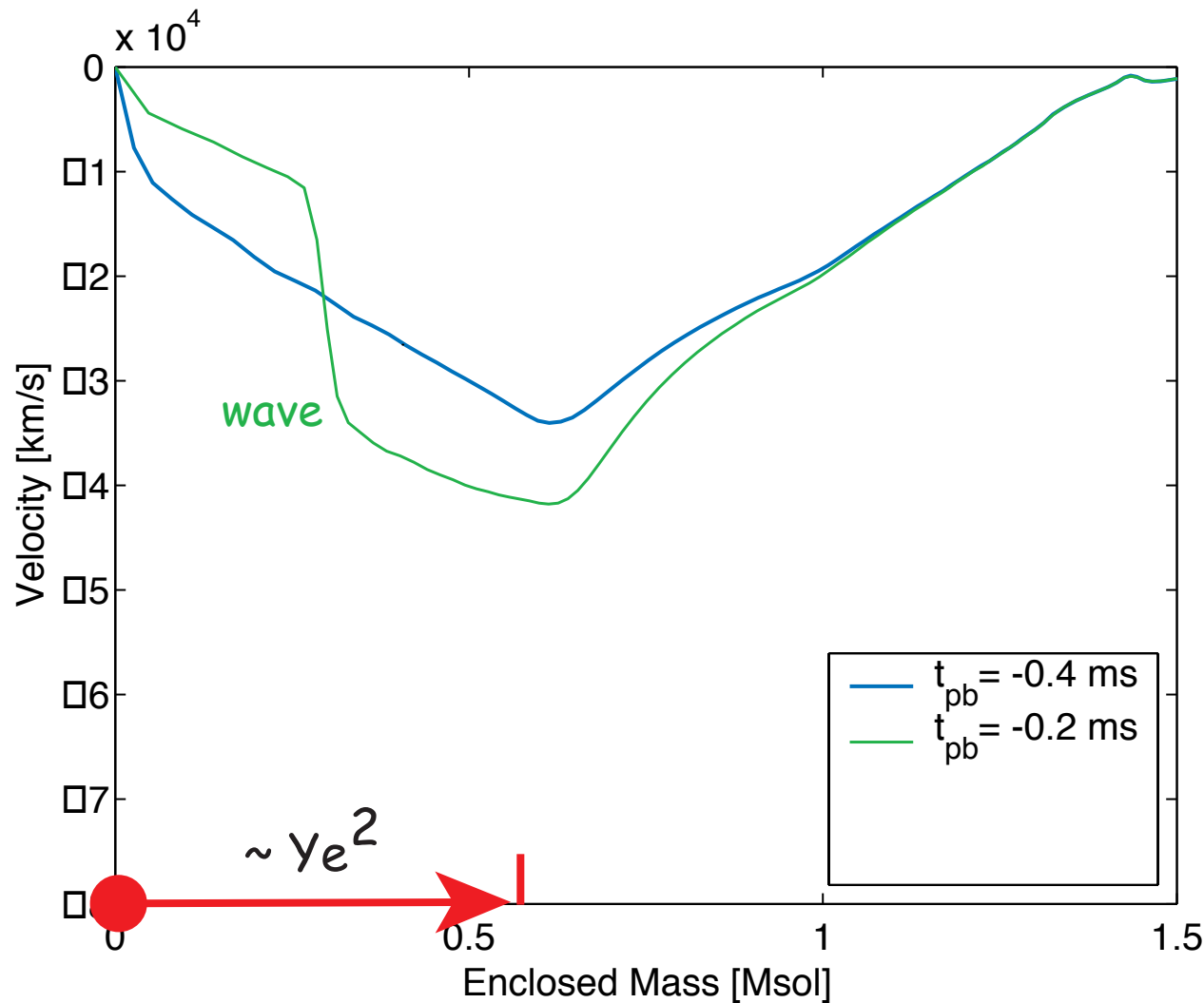
- causally connected collapse of iron core
- lower density in outer layers --> lower sound speed
- retarded infall of outer layers

$$M_{ic} \simeq (\kappa/\kappa_0)^{3/2} M_0,$$

$$\kappa = \frac{\ell_{ic}}{4} (3\pi^2)^{1/3} \left(\frac{Y_e}{\mu_B}\right)^{4/3}$$

(Goldreich & Weber 1980)

Neutronization during collapse



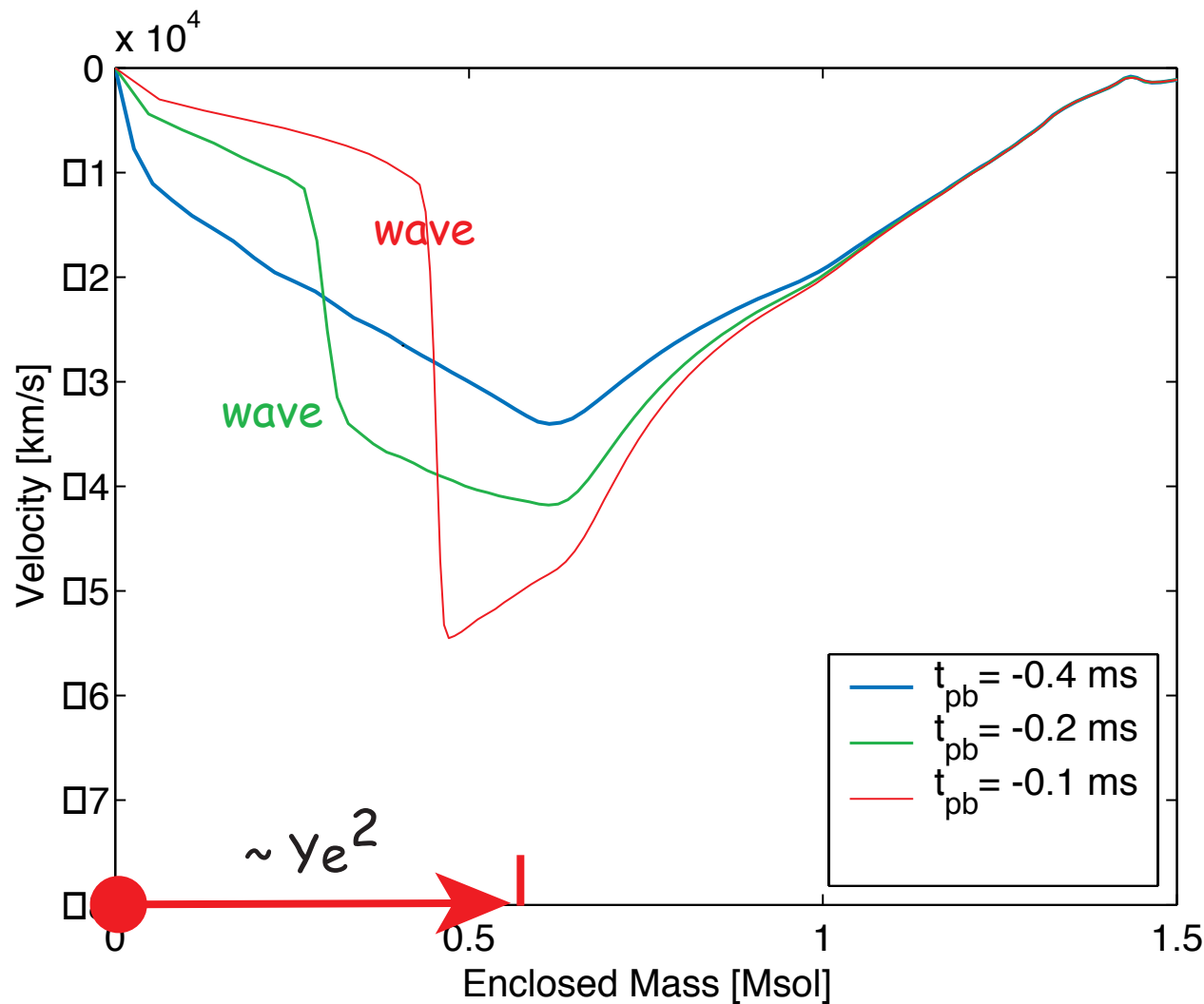
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Neutronization during collapse



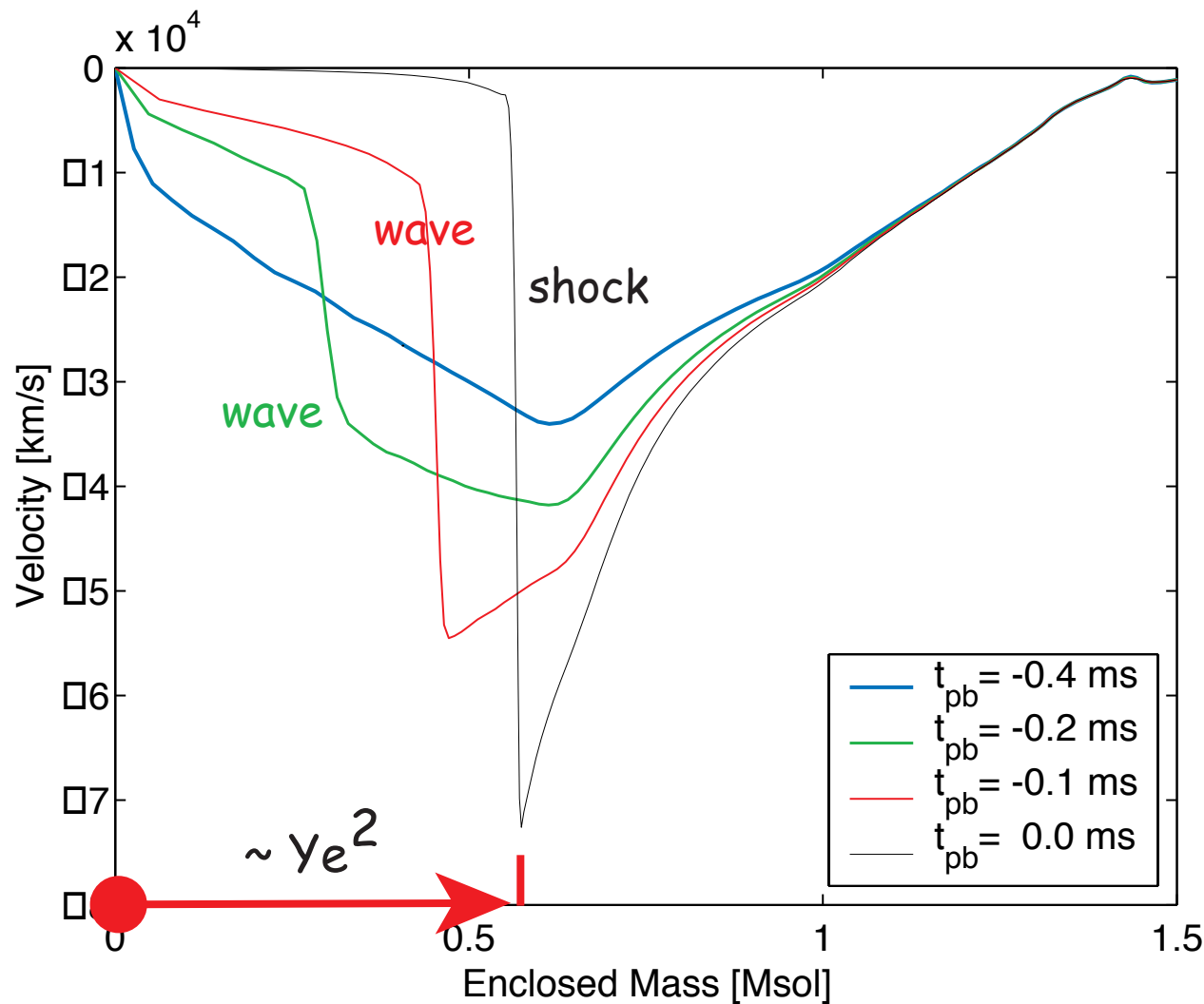
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Neutronization during collapse



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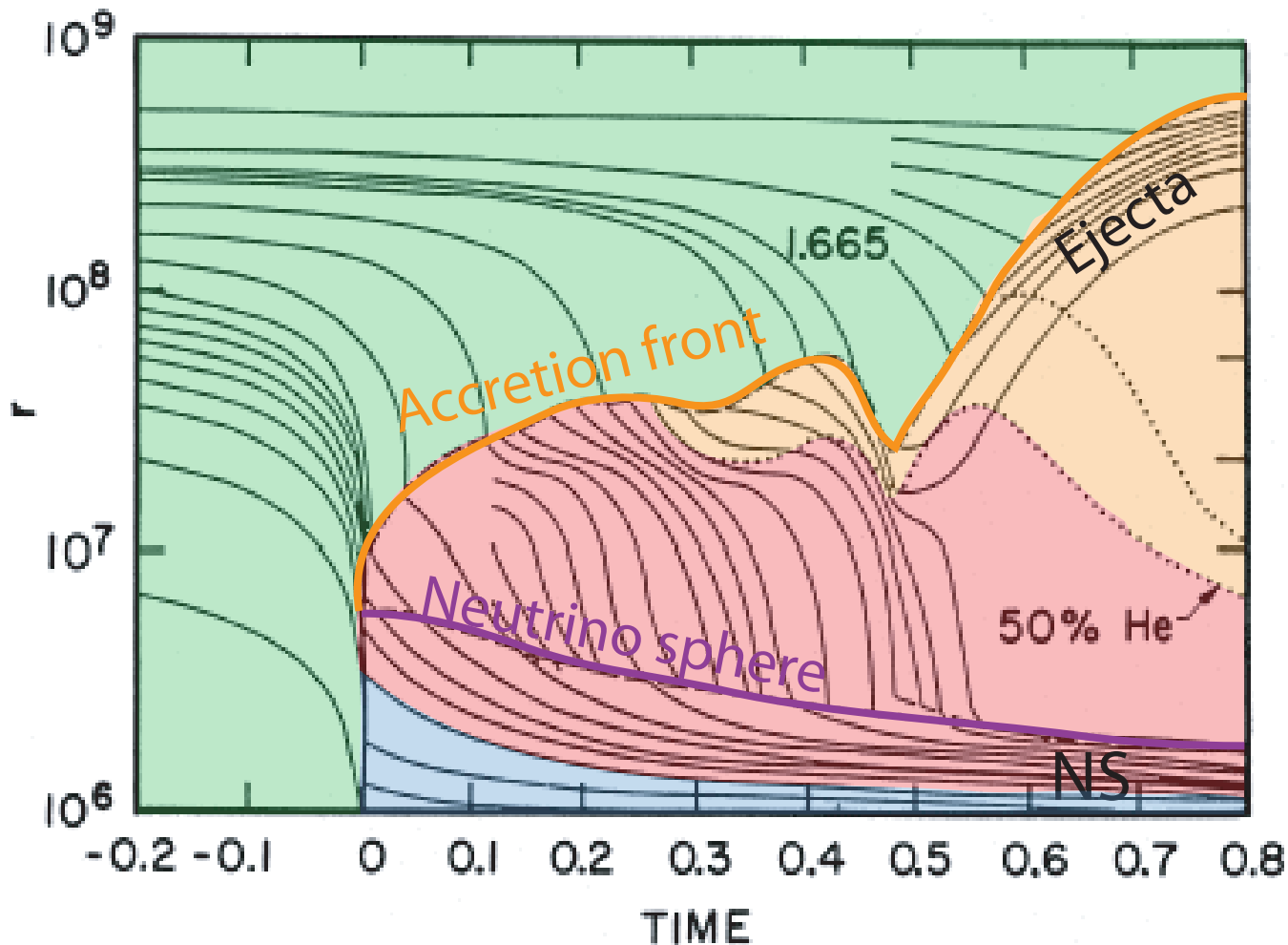
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Overview matter conditions

BETHE AND WILSON ApJ 295 (1985)



Colgate & White, ApJ 143 (1966)

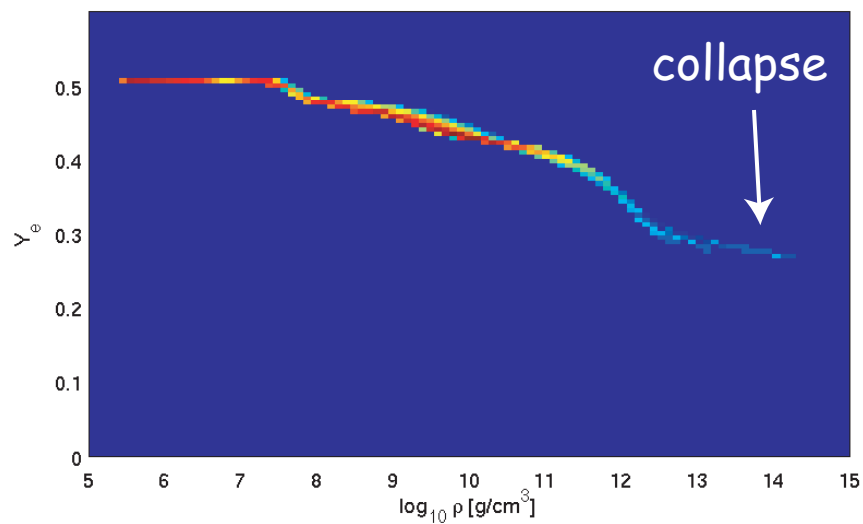
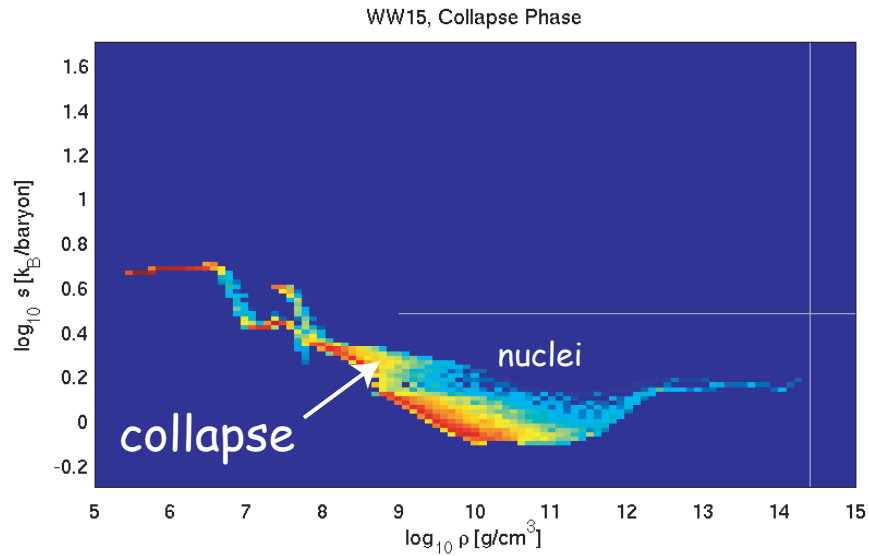
1) Ensemble
of nuclei

2) Cool bulk
nuclear matter

3) Hot dissociated

4) Freeze-out
of nuclei

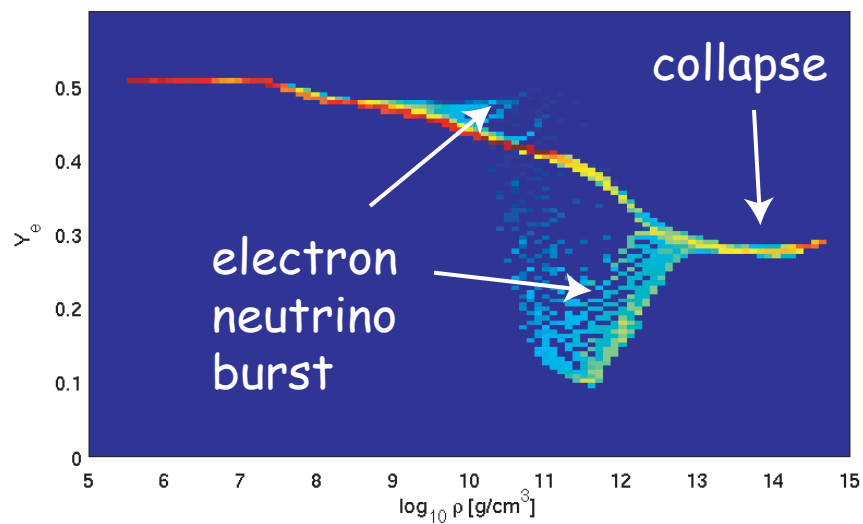
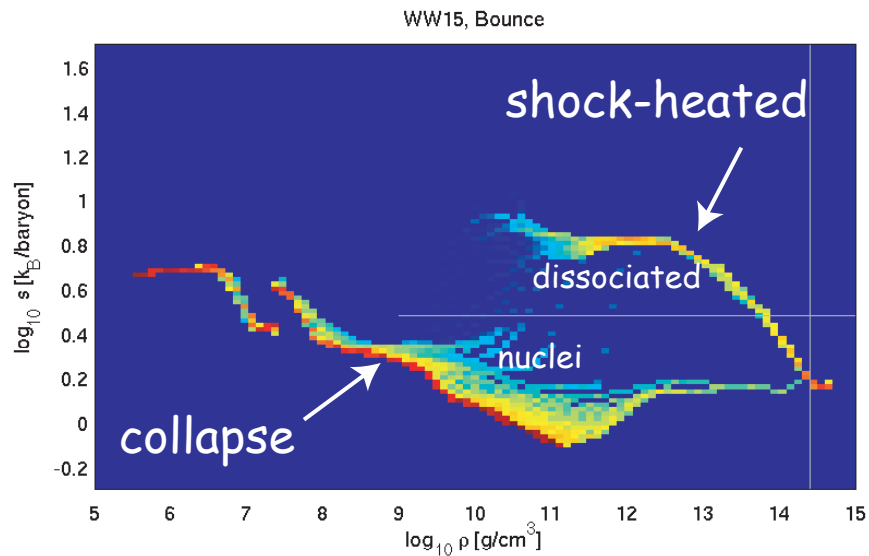
Conditions in (ρ, s, Y_e) -space



Collapse Phase:

- deelectronization along narrow trajectory
- slight entropy increase

Conditions in (ρ, s, Y_e) -space



Collapse Phase:

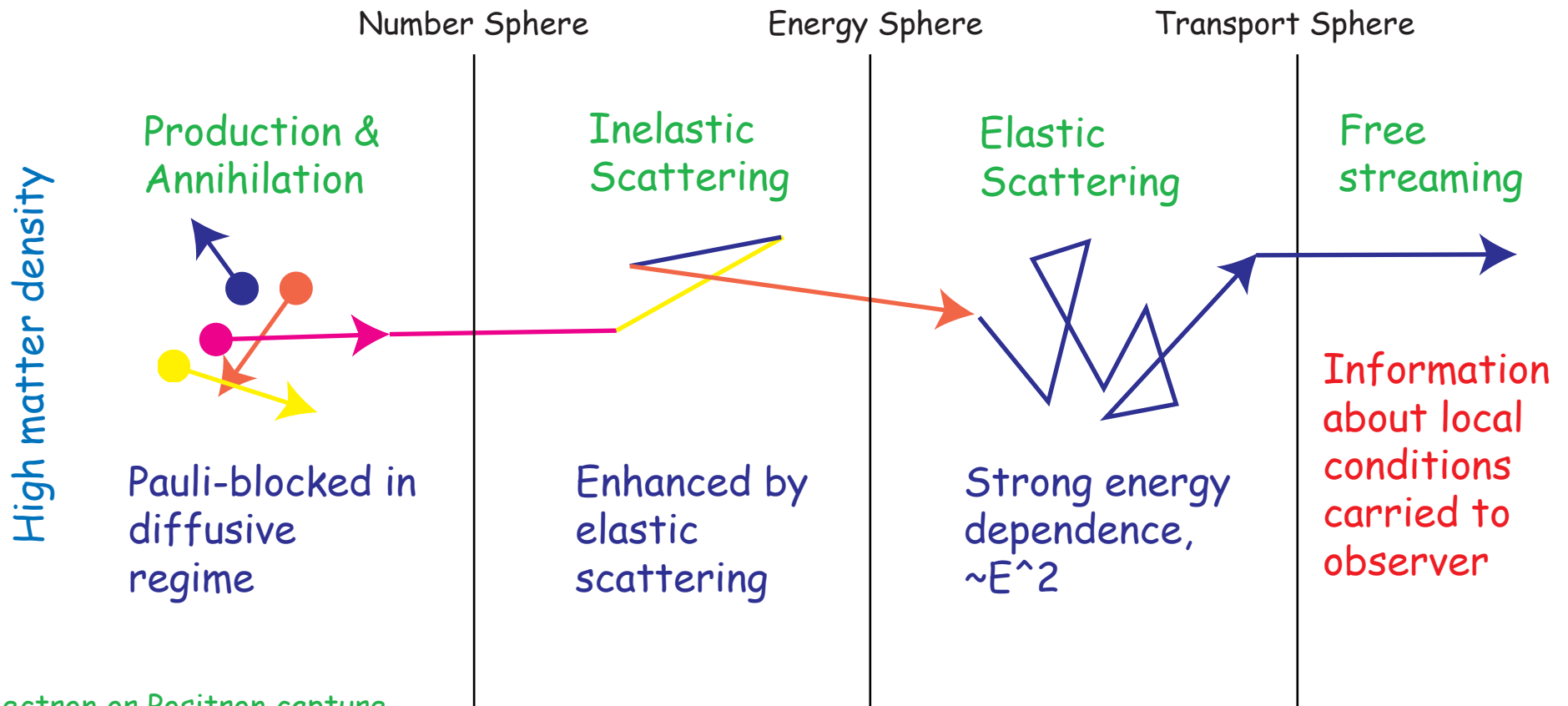
- deelectronization along narrow trajectory
- slight entropy increase

Postbounce Phase:

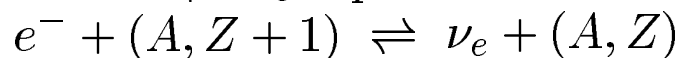
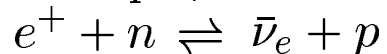
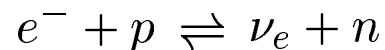
- jump to dissociated nucleon plasma

Neutrino-matter interactions

Bruenn (1985)
Raffelt (2001)



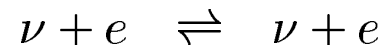
Electron or Positron capture



NN bremsstrahlung (Thompson et al. 2002)

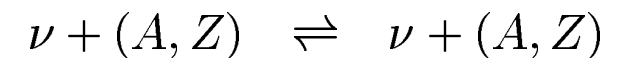
$\square e$ pair \rightarrow $\square \square$ pair (Buras et al. 2003)

Neutrino-electron scattering

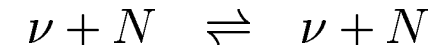


Inelastic neutrino-nucleus scattering

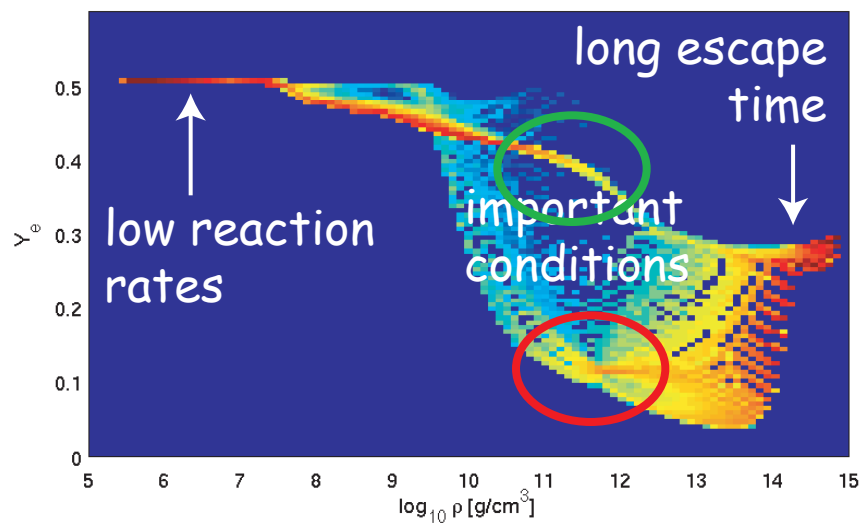
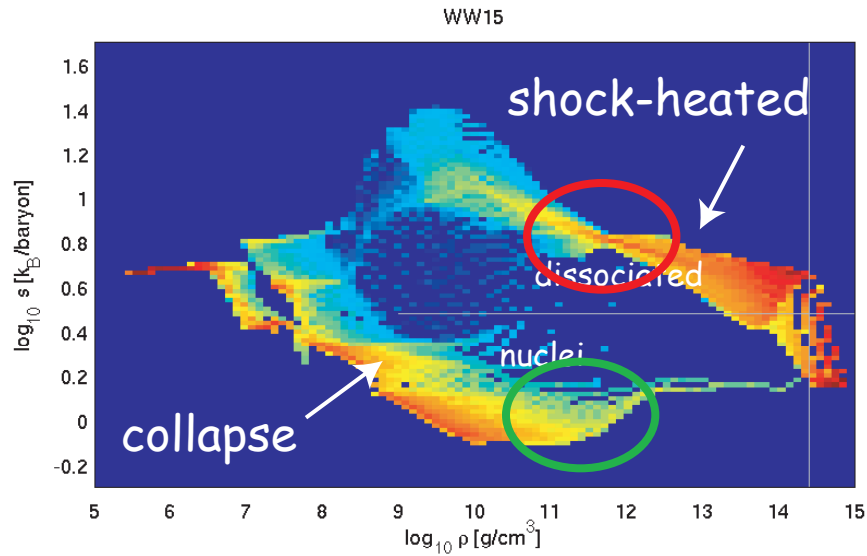
Elastic coherent scattering of neutrinos on nuclei



Neutrino-nucleon scattering



Relevant μ -matter interactions



The conditions around the neutrino spheres are marked in

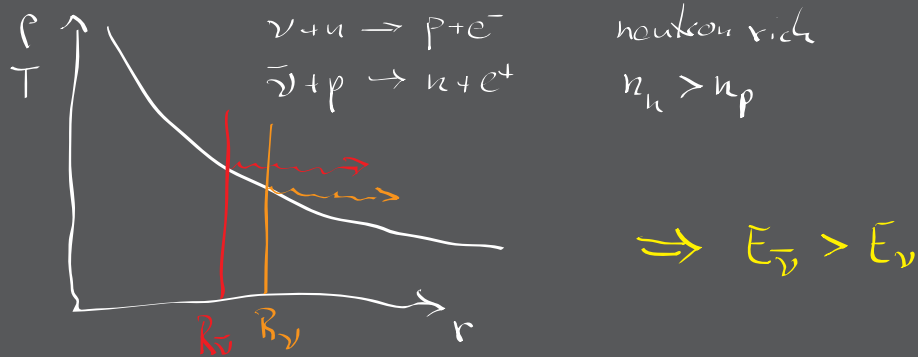
green ... collapse

red ... postbounce

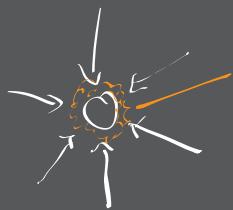
Different luminosity contributions

• Luminosity composed of two parts:

2) neutrinos of cooling protoneutron star



1) neutrinos from accretion flow



compression
of degenerate
electron gas



- Classical hierarchy among neutrino energies reflects temperature at neutrinospheres

- large accretion rate

--> $L_{\nu} \sim L_{\text{neutrino}}$

Solving the Boltzmann equation



$$\begin{aligned}
 & \frac{\partial F}{\alpha c \partial t} + \frac{\partial (4\pi r^2 \alpha \rho \mu F)}{\alpha \partial m} + \Gamma \left(\frac{1}{r} - \frac{\partial \alpha}{\alpha \partial r} \right) \frac{\partial [(1 - \mu^2) F]}{\partial \mu} \\
 & + \left(\frac{\partial \ln \rho}{\alpha c \partial t} + \frac{3u}{r c} \right) \frac{\partial [\mu (1 - \mu^2) F]}{\partial \mu} \\
 & + \left[\mu^2 \left(\frac{\partial \ln \rho}{\alpha c \partial t} + \frac{3u}{r c} \right) - \frac{1u}{r c} - \mu \Gamma \frac{\partial \alpha}{\alpha \partial r} \right] \frac{1}{E^2} \frac{\partial (E^3 F)}{\partial E} \\
 & = \frac{j}{\rho} - \tilde{\chi} F + \frac{1}{h^3 c^4} E^2 \int d\mu' R_{is}(\mu, \mu', E) F(\mu', E) \\
 & - \frac{1}{h^3 c^4} E^2 F \int d\mu' R_{is}(\mu, \mu', E) \\
 & + \frac{1}{h^3 c^4} \left[\frac{1}{\rho} - F(\mu, E) \right] \int E'^2 dE' d\mu' \tilde{R}_{nes}^{in}(\mu, \mu', E, E') F(\mu', E) \\
 & - \frac{1}{h^3 c^4} F(\mu, E) \int E'^2 dE' d\mu' \tilde{R}_{nes}^{out}(\mu, \mu', E, E') \left[\frac{1}{\rho} - F(\mu', E') \right]
 \end{aligned}$$

$$\frac{\partial Y_e}{\partial t} = -\frac{2\pi m_B}{h^3 c^2} \int E^2 dE d\mu \left(\frac{j}{\rho} - \tilde{\chi} F \right) \quad \frac{\partial e}{\partial t} = \dots \quad \frac{\partial u}{\partial t} = \dots$$

(Mezzacappa & Bruenn 1993, Liebendörfer 2000, Liebendörfer et al. 2004)

Evolution of specific neutrino distr. function:

$$F(t, m, \square, E) = f(t, r, \square, E) / \square$$

=> 3D implicit problem

Comoving metric:

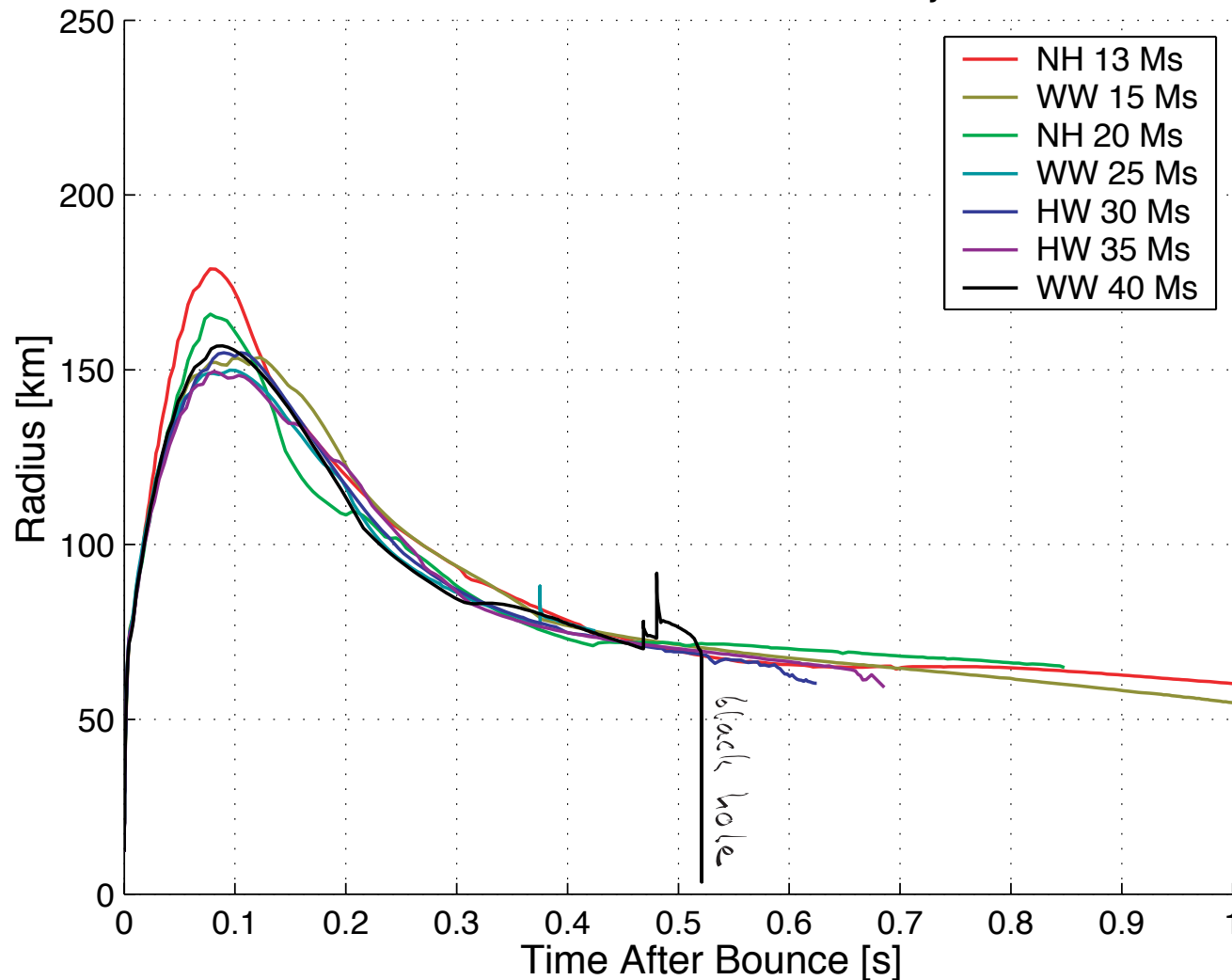
$$\begin{aligned}
 ds^2 &= -\alpha^2 dt^2 + \left(\frac{1}{\Gamma} \frac{\partial r}{\partial a} \right)^2 \\
 &+ r^2 (d\vartheta^2 + \sin^2 \vartheta d\varphi^2)
 \end{aligned}$$

Stress-energy tensor:

$$\begin{aligned}
 T^{tt} &= \rho(1 + e + J) \\
 T^{ta} = T^{at} &= \rho H \\
 T^{aa} &= p + \rho K \\
 T^{\vartheta\vartheta} = T^{\varphi\varphi} &= p + \frac{1}{2} \rho (J - K)
 \end{aligned}$$

Modeling in spherical symmetry

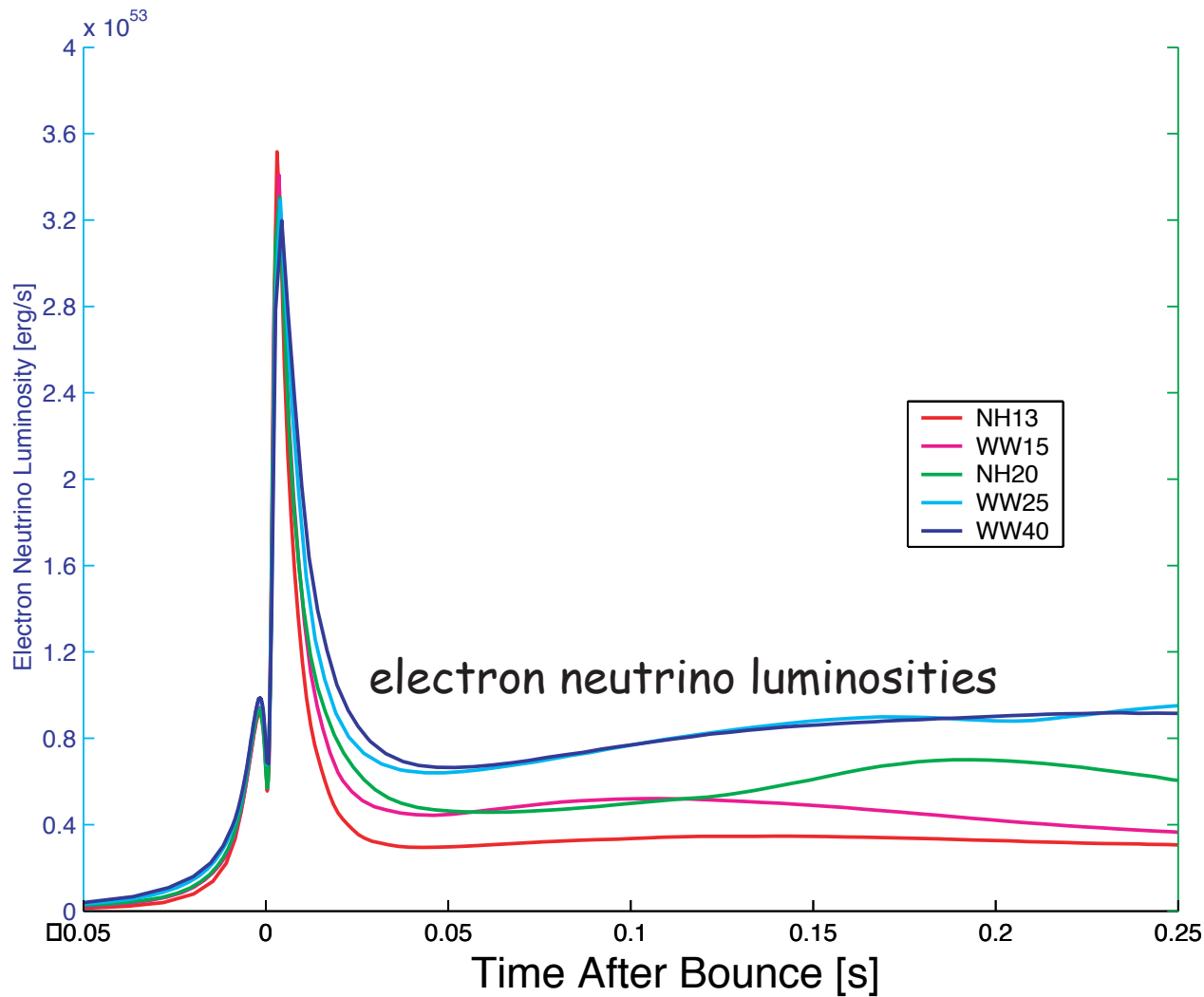
General Relativistic Shock Trajectories



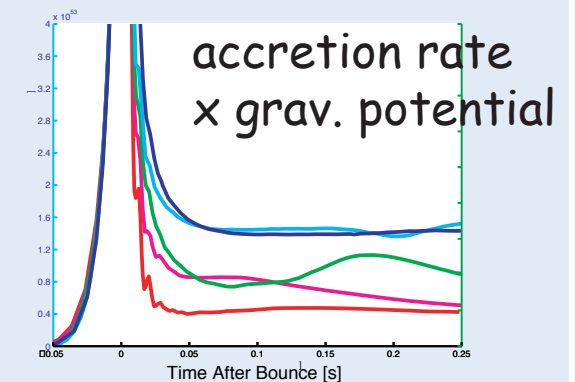
(Liebendörfer, Mezzacappa, Messer, Hix, Martinez-Pinedo, Thielemann, 2003)

- Trajectories of the accretion front for different progenitor stars $13M_{\text{sol}} < M < 40M_{\text{sol}}$
- calculated with Agile-Boltztran
- 40 Msol model forms a black hole
- 13 Msol model more optimistic
- no explosions obtained

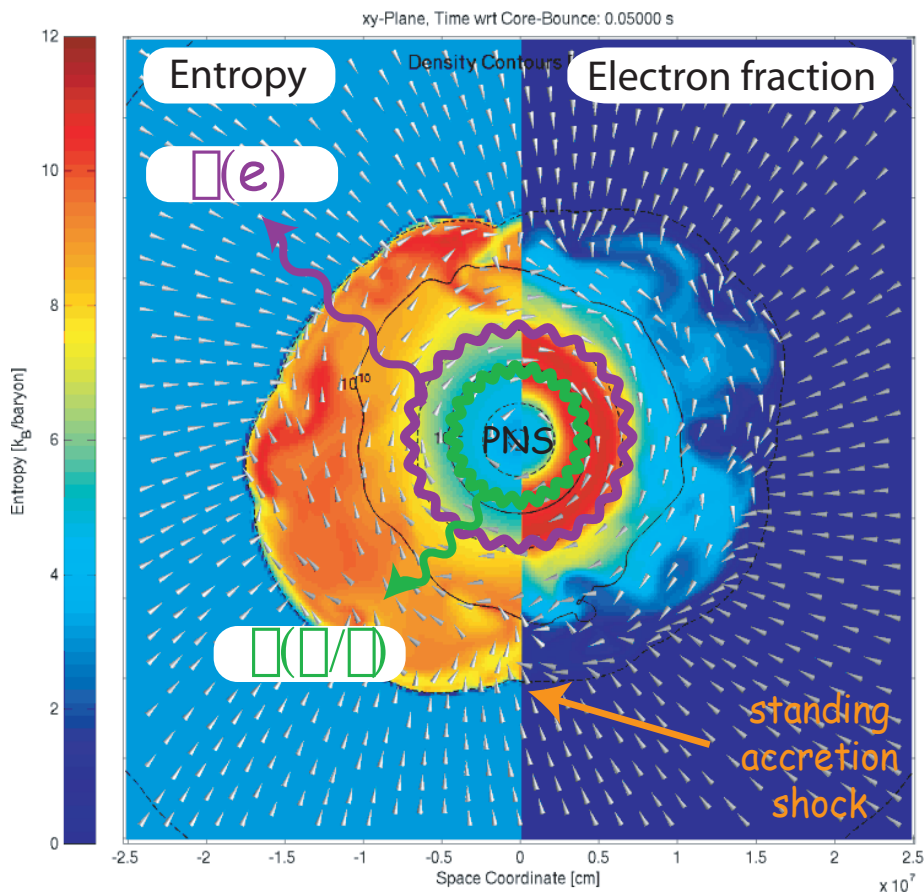
Neutrino signal



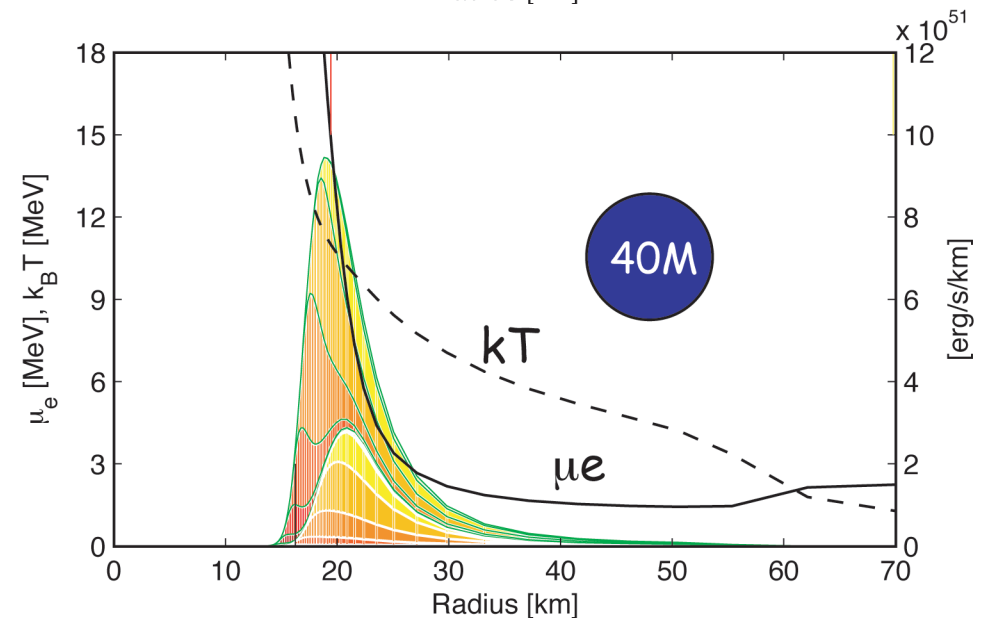
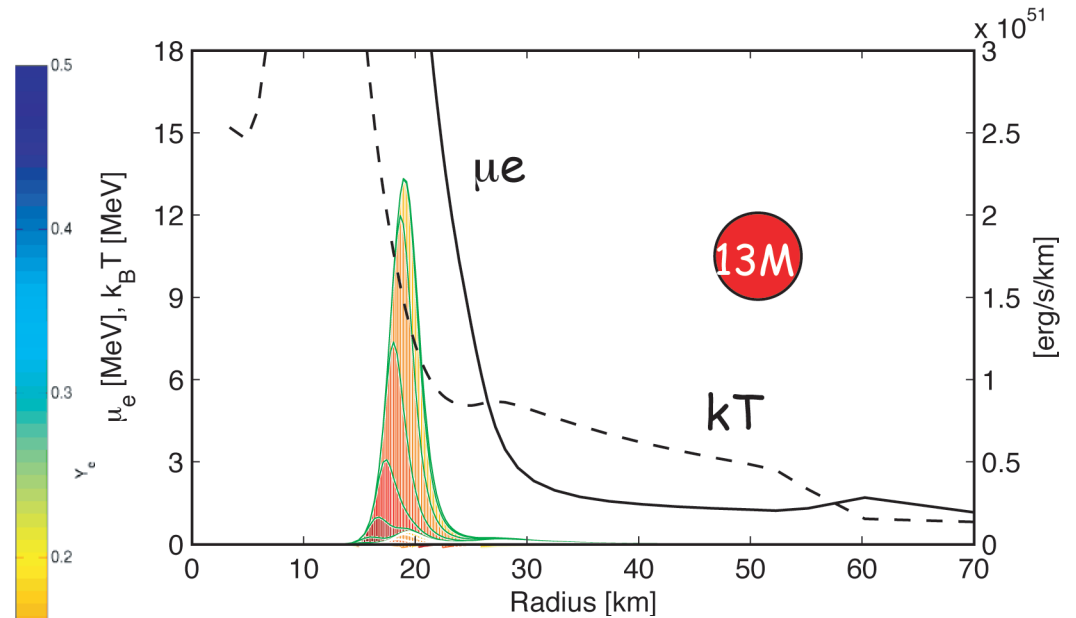
- initially similar luminosities
- differences appear in accretion phase
- >50% accretion lumin.
- density profiles in outer progenitor layers very different



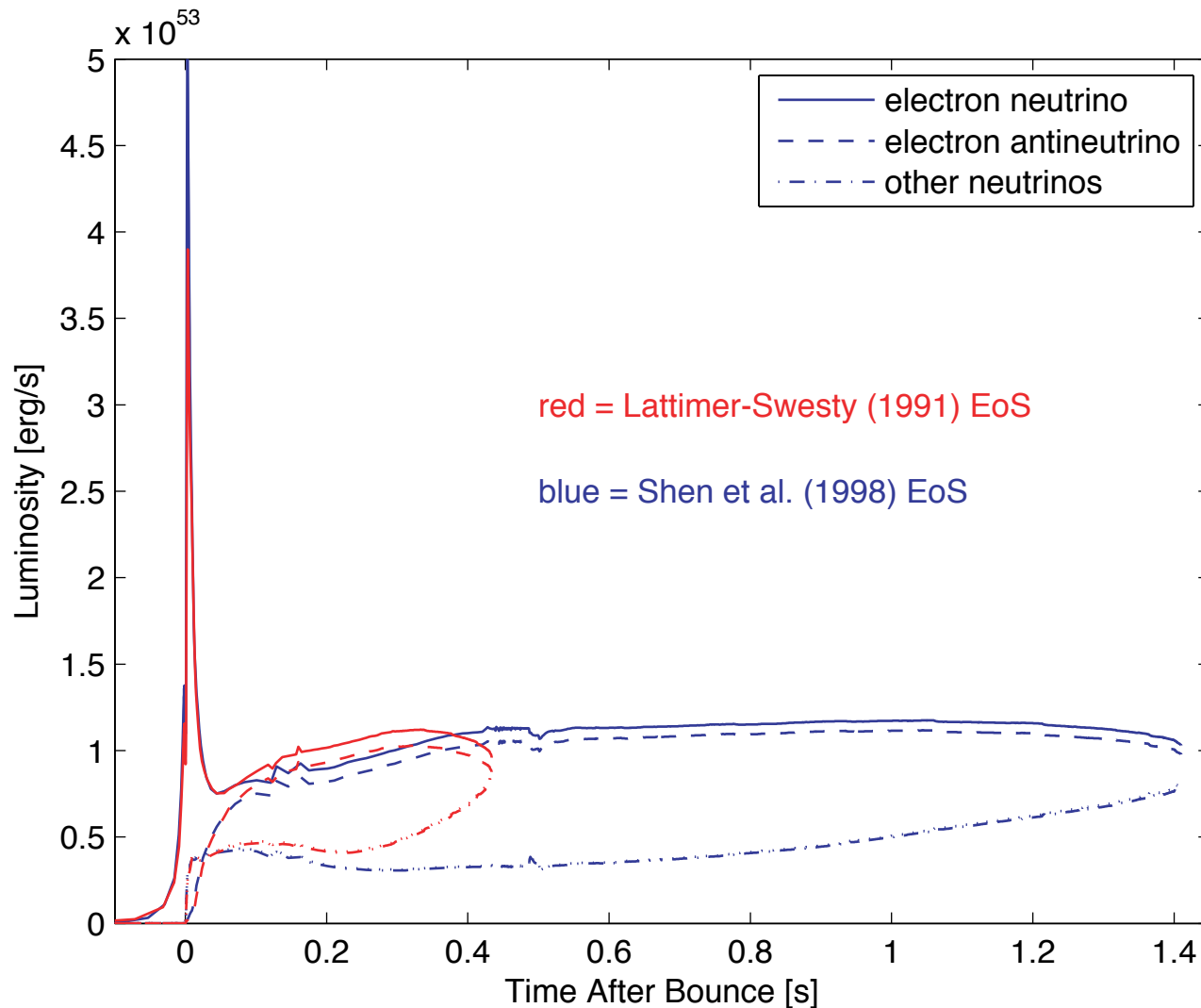
$\square(\square/\square)$ signal from PNS evolution



- low mass proto-neutron star (PNS) --> incompressible accretion
- PNS close to maximum mass --> hot layers pushed inward



Sensitivity with respect to EoS

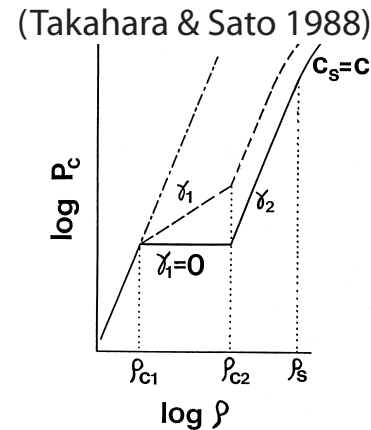


- Collapse, bounce, and postbounce evolution til black hole formation
- The quasi-static compression of the protoneutron star is reflected in mu/tau neutrino luminosities
- The different stiffness of the EoS causes very different delay times until BH formation

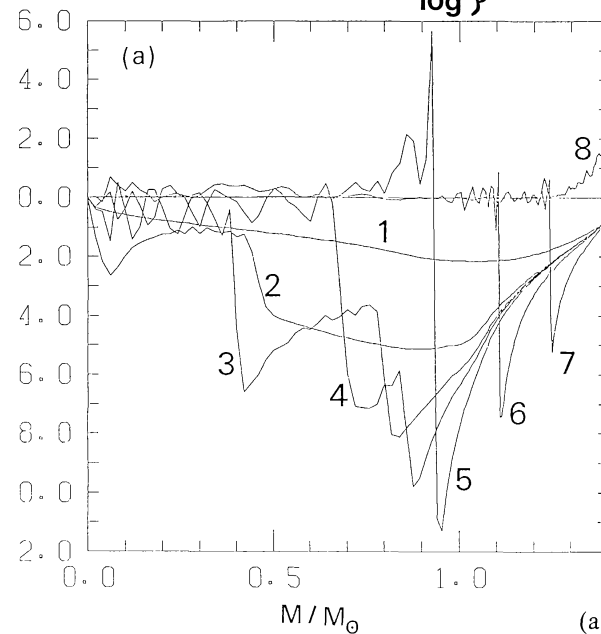
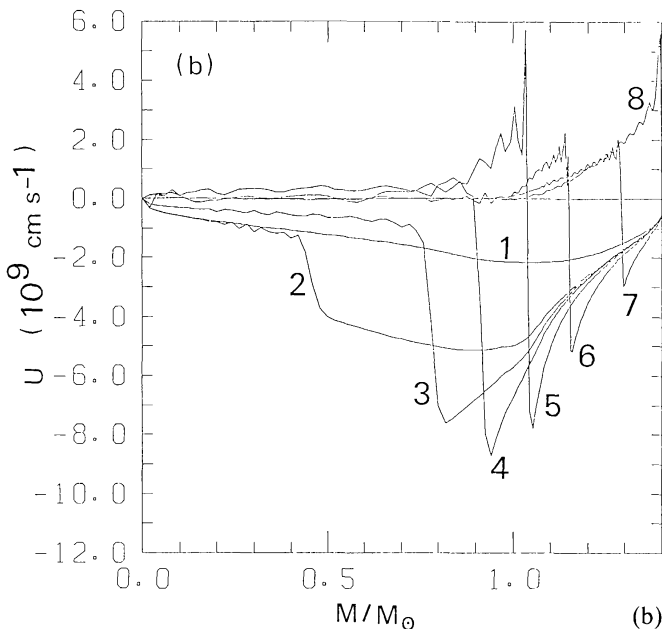
(Fischer et al. 2008, similar Sumiyoshi et al. 2007)

Signals of QCD phase transition?

- early discussion, revived by SN1987A neutrinos (e.g. Migdal et al. 1979, Takahara & Sato 1985-88)
- investigations with parameterised equations of state and GR hydrodynamics
- more realistic EoS's and GR hydrodynamics (Gentile et al. 1993)



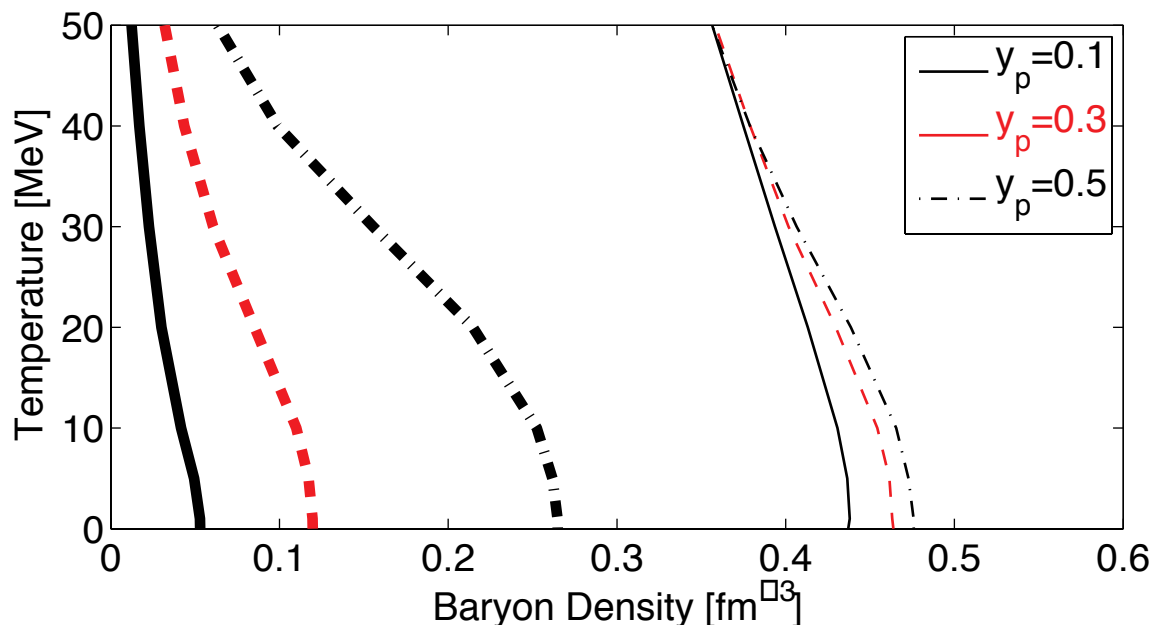
(Takahara & Sato 1986)



- select phase transition at or immediately after core bounce
- a second shock forms
- catches up with first shock
- > Is this observable?
- weak interactions and neutrinos neglected
- simulations only to few ms postbounce

Simple model for phase transition

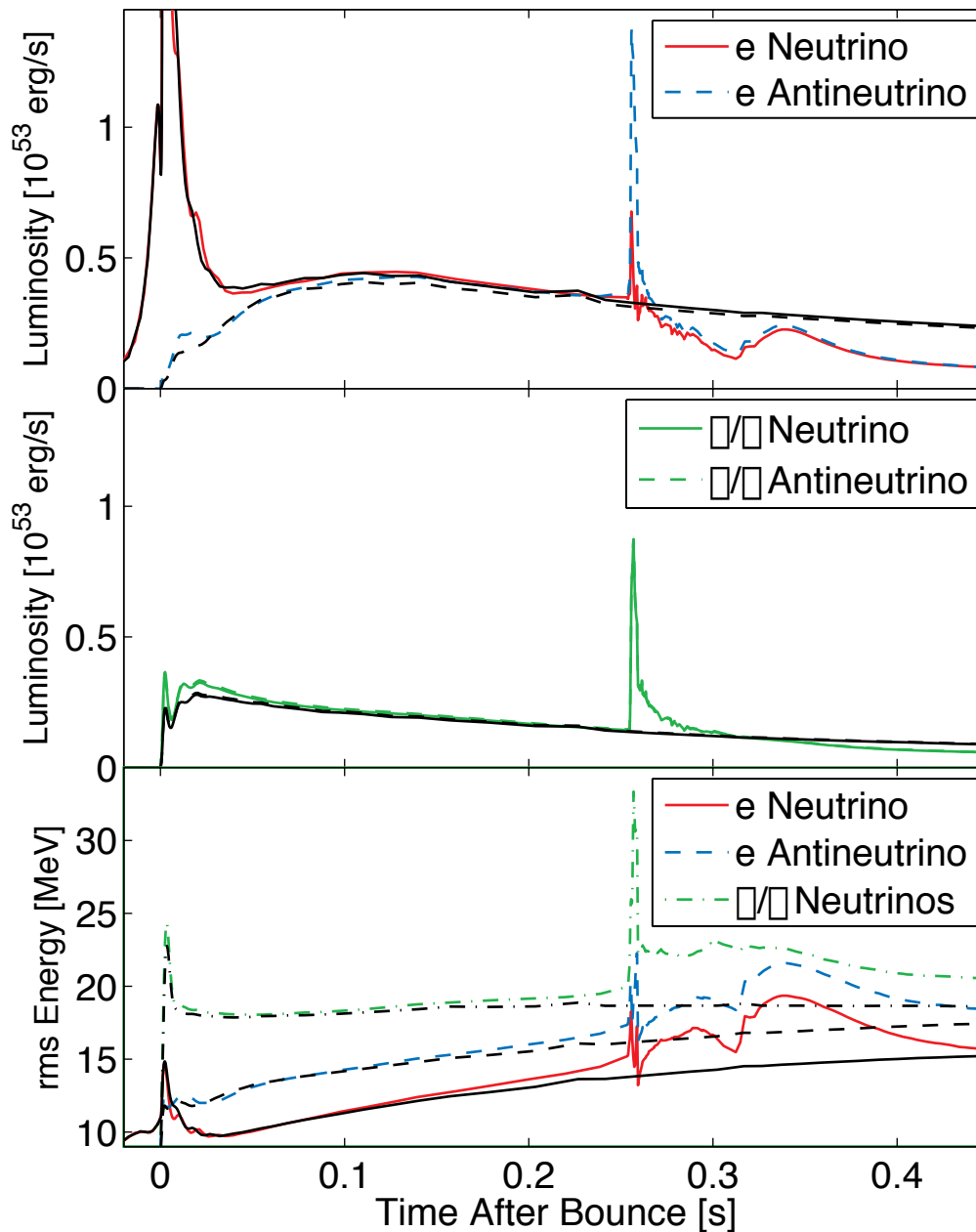
- GR Boltzmann neutrino transport as discussed above
- Shen et al. 1998 equation of state for hadronic phase
- MIT bag model for quark phase, choosing parameters for early phase transition: $B=162-165$ MeV, $m_s=100$ MeV
- Mixed phase according to Gibbs construction (mechanical and chemical equilibrium, μ 's trapped)



(I. Sagert et al, T. Fischer et al. 2008, submitted)

- compatible with heavy ion data
 - isospin-asymmetric
 - weak equilibrium allows for strange quarks
- 'just' compatible with neutron star data:
 - 162 supports 1.56 Ms
 - 165 supports 1.50 Ms

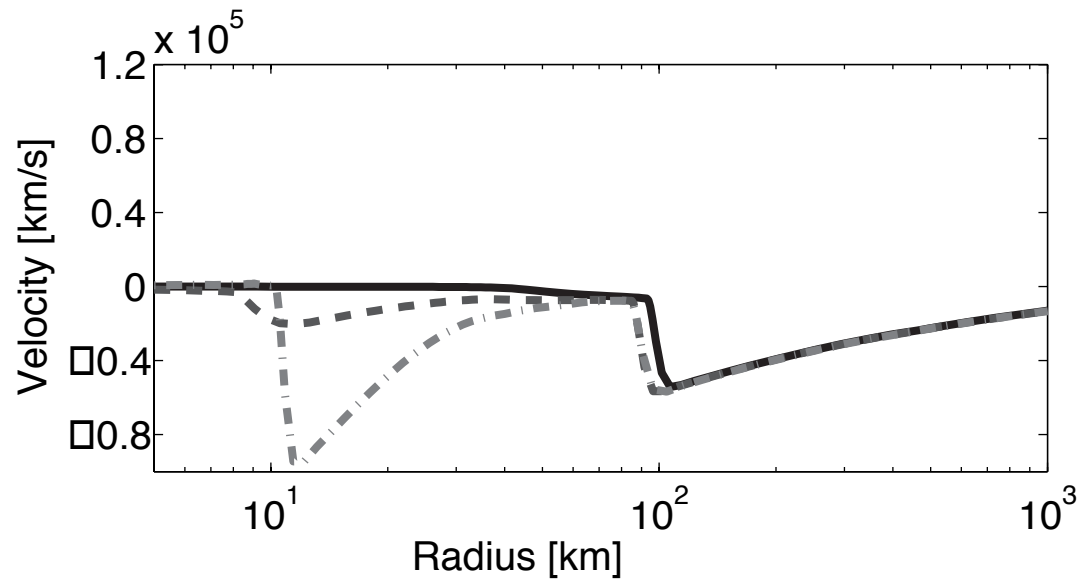
Neutrino signature of phase transition



Shown is a simulation of a 10 Ms star containing quark matter ($B=162$) compared to one with hadronic matter only (black lines)

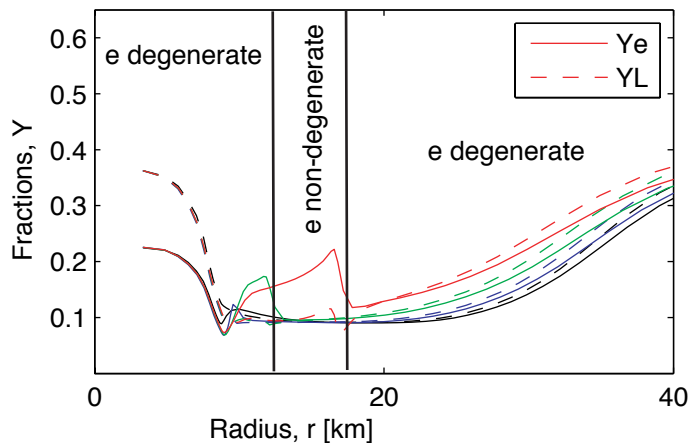
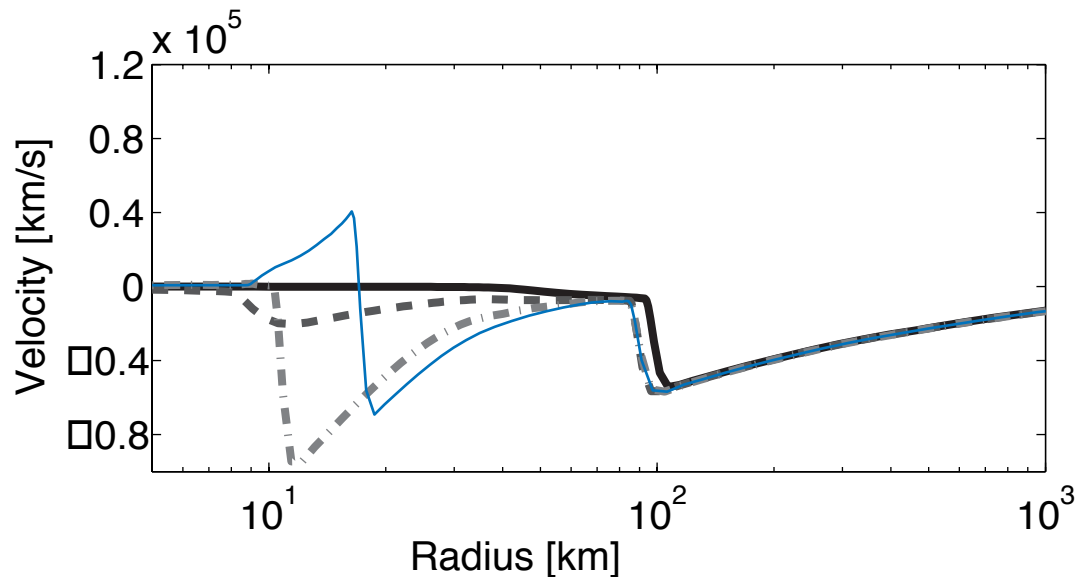
- strong second neutrino burst in all flavours
- electron anti-neutrinos dominate
- step up in neutrino rms energies

Different dynamical stages



- collapse
- conversion to quark phase from inside out
- shrinking mixed phase
- second accretion shock propagating outward
- shock propagates with mixed-hadronic phase boundary

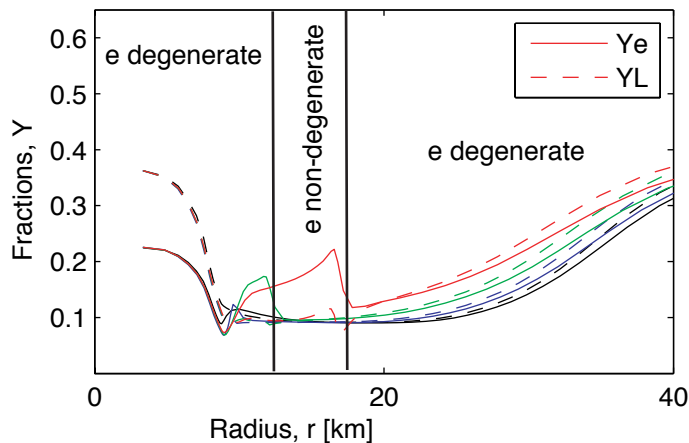
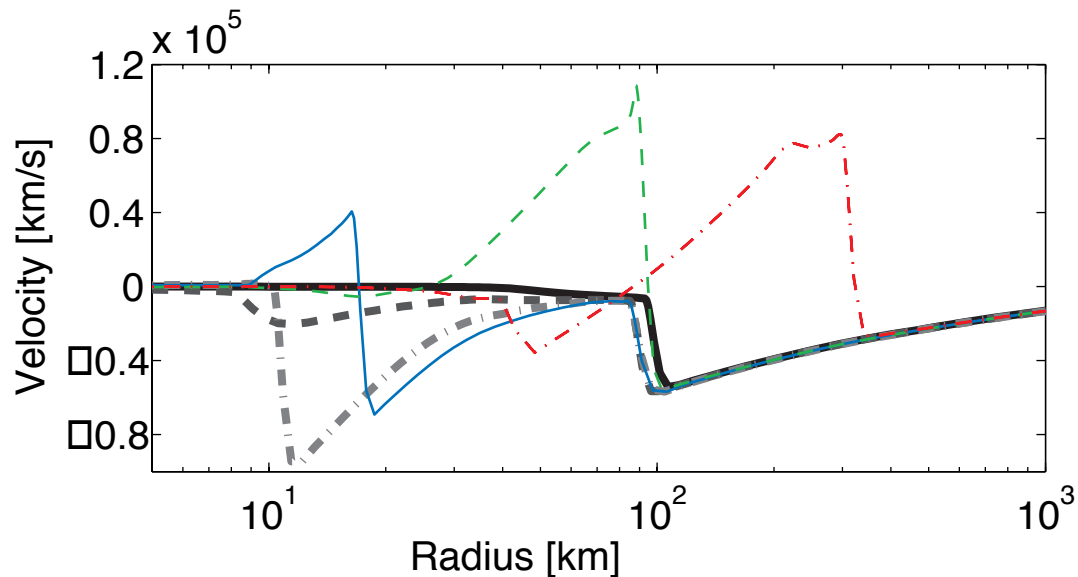
Different dynamical stages



- de-leptonised matter becomes **non-degenerate**
- weak equilibrium steps to **larger Y_e**
- pressure increases

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- conversion to quark phase from inside out
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- accr. shock detaches from phase boundary to reach \square -spheres in the hadronic phase

Different dynamical stages

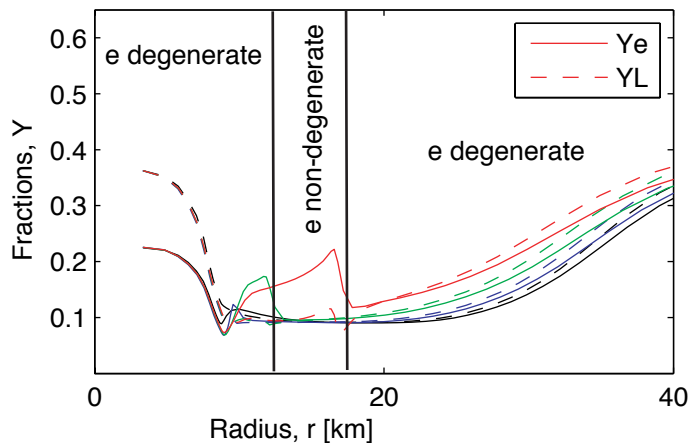
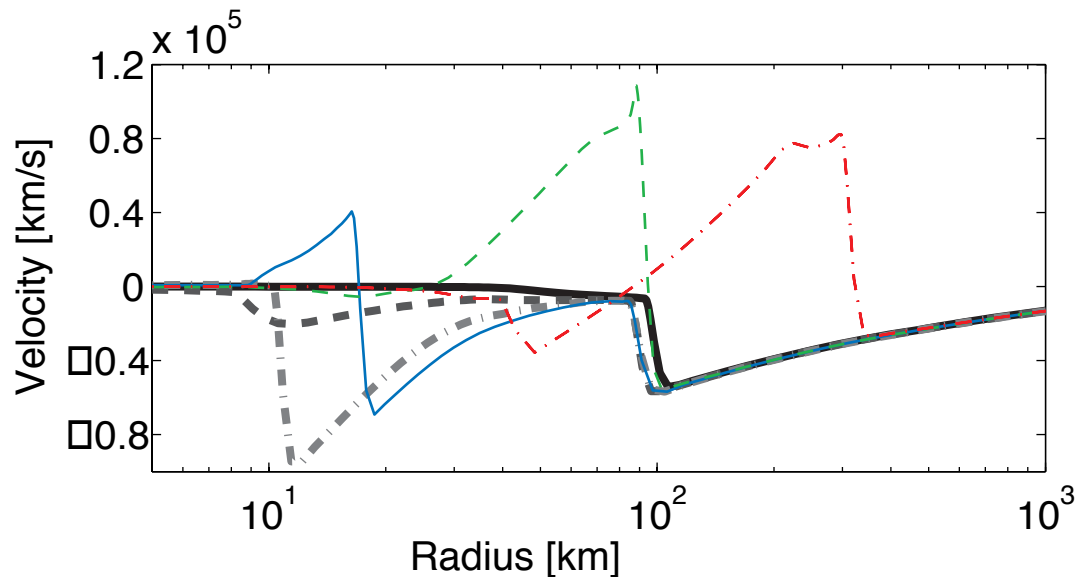


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- shocked matter accelerates and triggers explosion

(Sagert et al., Fischer et al. 2008)

Different dynamical stages



- de-leptonised matter becomes **non-degenerate**
- weak equilibrium steps to **larger Y_e**
- pressure increases
- emission of **anti-neutrino** dominates when neutrino spheres are reached

- collapse
- conversion to quark phase from inside out
- shrinking mixed phase
- second accretion shock propagating outward
- shock propagates with mixed-hadronic phase boundary
- accr. shock detaches from phase boundary to reach \square -spheres in the hadronic phase
- shocked matter accelerates and triggers explosion

(Sagert et al., Fischer et al. 2008)

Bag constant & progenitor variations

Larger bag constant

- > longer postbounce accretion time
- > more massive protoneutron star
- > deeper gravitational potential
- > larger peak luminosity in second neutrino burst
- > larger explosion energies

Prog.	B	t_{pb}	M_Q	M_{mixed}	$M_{PN S}$	E_{expl}
[M _⊙]	[MeV]	[ms]	[M _⊙]	[M _⊙]	[M _⊙]	[10 ⁵¹ erg]
10	162	255	0.850	0.508	1.440	0.44
10	165	448	1.198	0.161	1.478	1.64
15	162	209	1.146	0.320	1.608	0.42
15	165	330 ^a	1.496	0.116	1.700	unknown ^b

^a moment of black hole formation

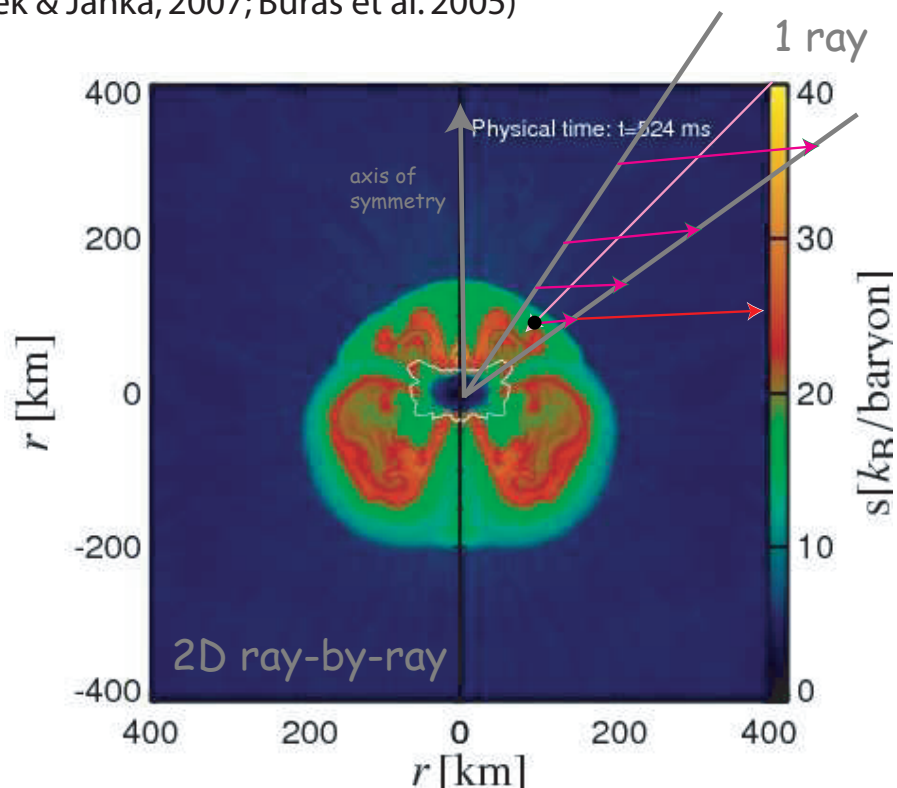
^b black hole formation before positive explosion energy is achieved

- Is tuning of parameters or the model of the quark phase possible to reproduce SN1987A?
- How do more massive progenitors explode?
- Weak Λ -driven explosion followed by phase transition?
- Some models eject low- Y_e matter --> a possible site for the r-process?

Or in combination with...

- Delayed neutrino-driven supernova explosions aided by the standing accretion-shock instability

(Marek & Janka, 2007; Buras et al. 2005)



- Standing accretion shock instability (SASI)

(Blondin & Mezzacappa 2003
Foglizzo et al. 2007)

- Features of the Acoustic Mechanism of Core-Collapse Supernova Explosions

(Burrows et al. 2006)

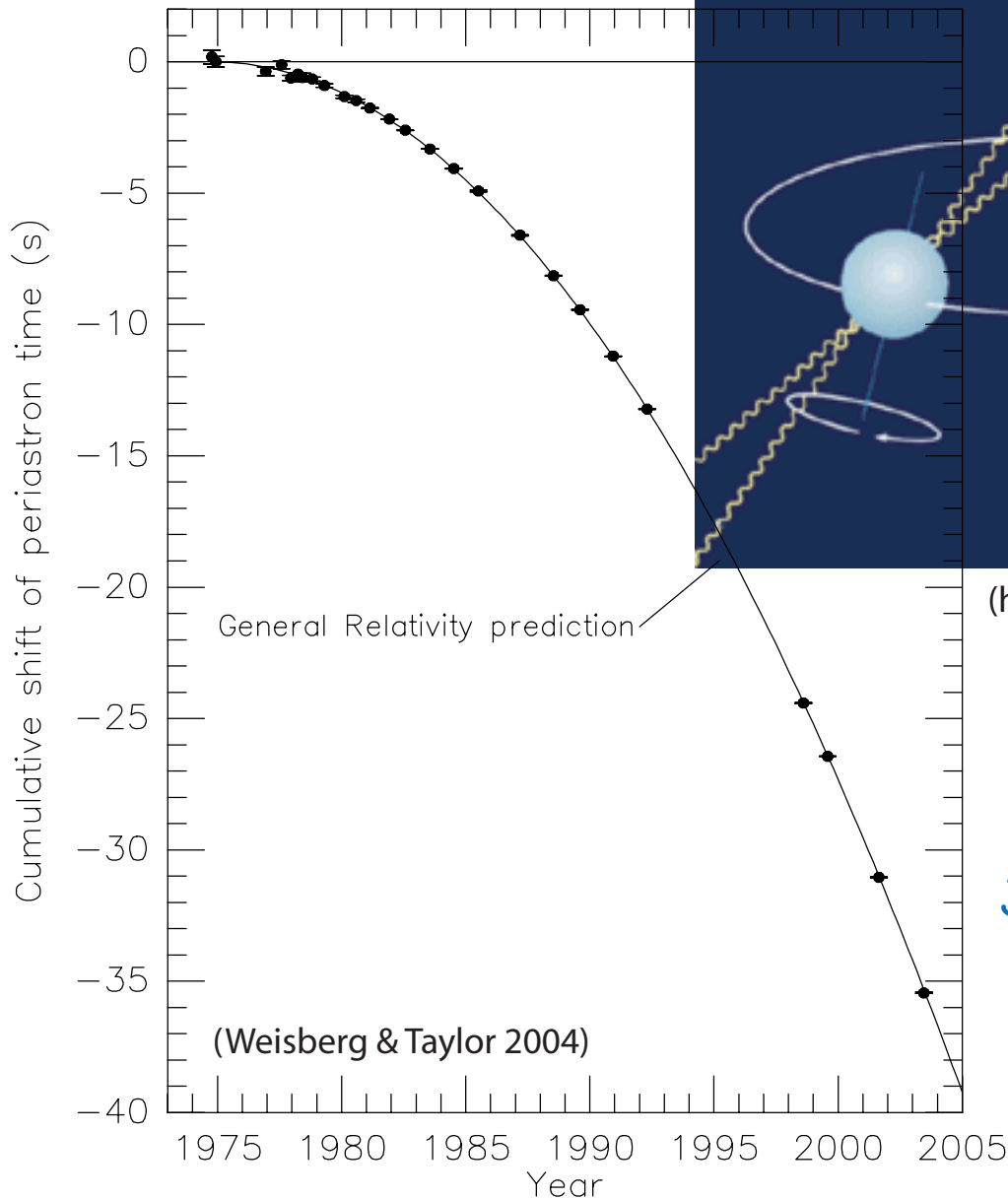
Heating by dissipation of emitted sound waves

- Magneto-rotational explosion mechanisms and collapsar model

(Bisnovatyi-Kogan 1976,
Leblanc & Wilsons 1979,
MacFadyen & Woosley 1999)

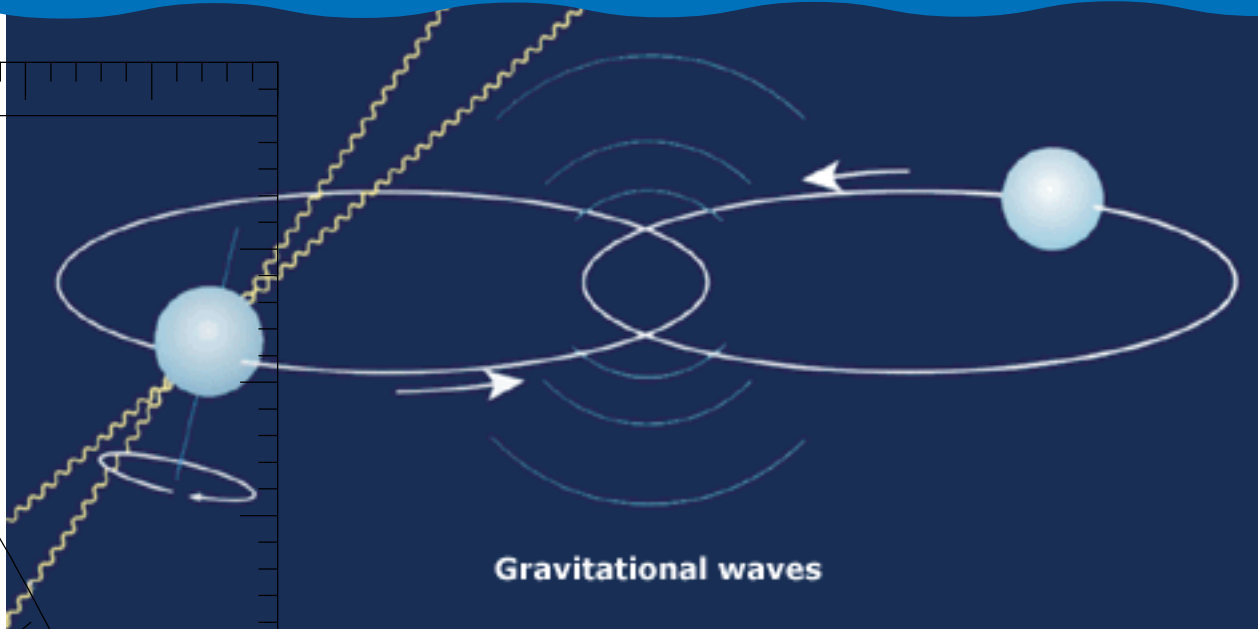
Explosion after black hole formation

Gravitational Waves



(<http://nobelprize.org>)

Nobel prize
1993
Joseph Taylor
Russell Hulse



Pulsar B1913+16

- precision mass measurement
 - slowdown measured
 - compared to GR prediction
- > agreement to within 0.2%!

Prediction of Gravitational Wave Signal



Numerous 3D hydrodynamics simulations of stellar core-collapse in Numerical Relativity community:

- based on simple polytropic equation of state
- neutrino physics neglected
- prediction of *GW* signal: type I-III wave forms

In the mean time improved by using a microscopic equation of state and development of parameterisation scheme for deleptonisation during collapse:

- parameterise Y_e as function of density from 1D
- estimate ds from dY_e
- estimate luminosity and \square -stress from $\text{int}(dY_e)$

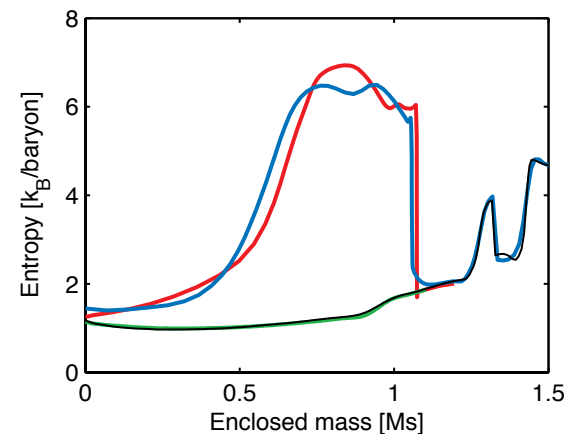
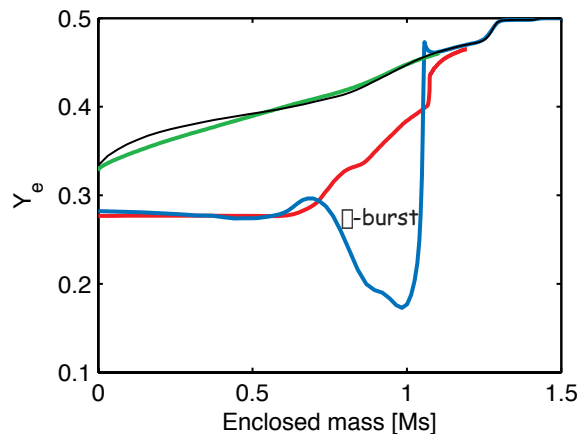
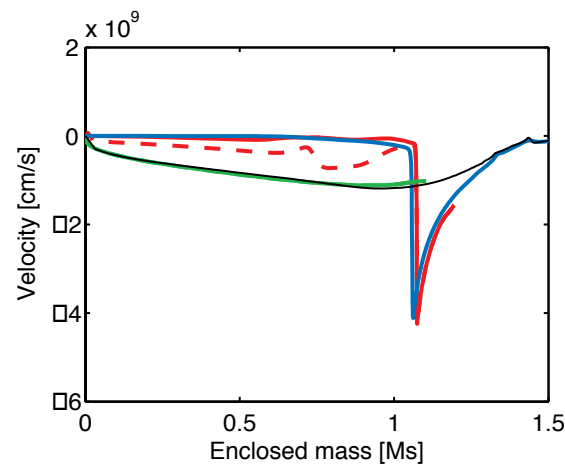
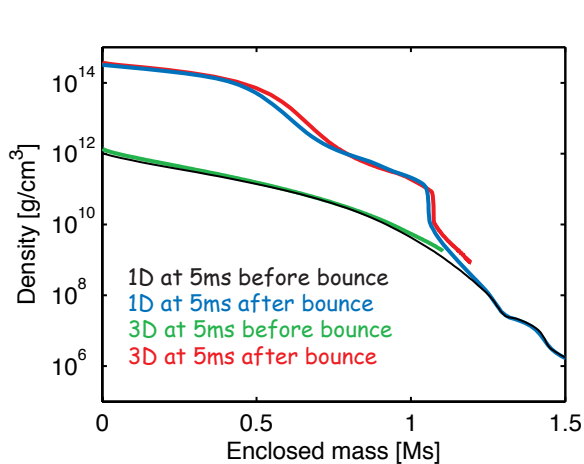
(Liebendörfer 2005)

- Only type I *GW* signals have been found!

(Dimmelmeier et al. 2007,
Ott et al. 2007,
Scheidegger et al. 2007)

3D Magneto-Hydrodynamics

- Parameterization of weak interactions for collapse phase
- Comparison 1D GR Boltzmann \leftrightarrow 3D approximations



3D MHD

(Pen, Arras, Wong 2003)

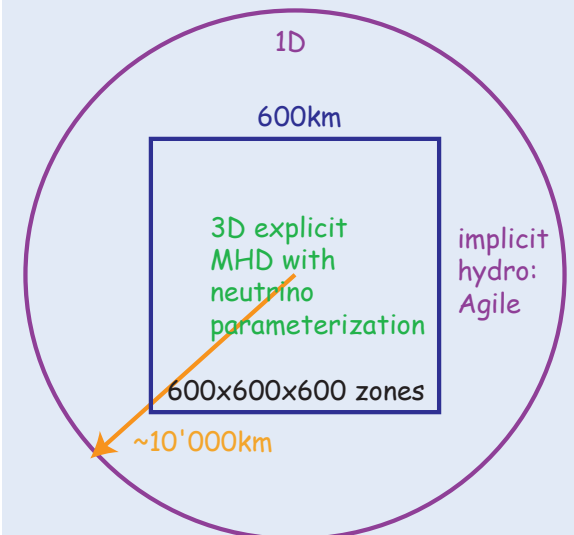
Lattimer-Swesty EOS

(Lattimer & Swesty 1991)

Effective GR potential

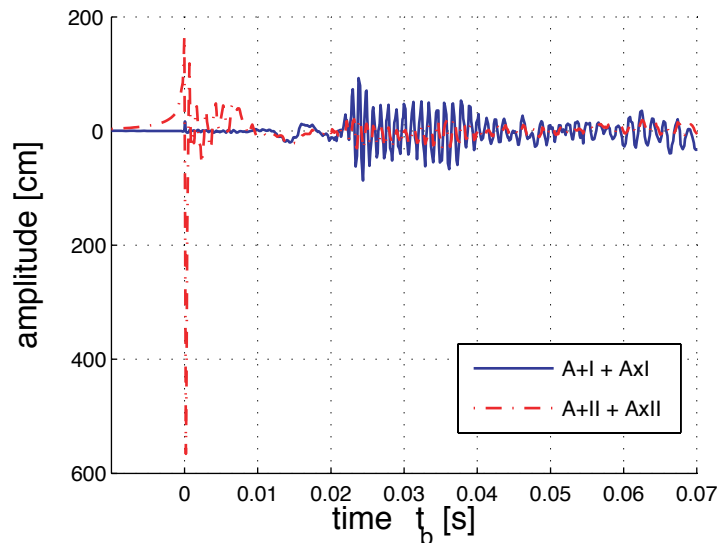
(Marek et al. 2006)

Fully parallelised

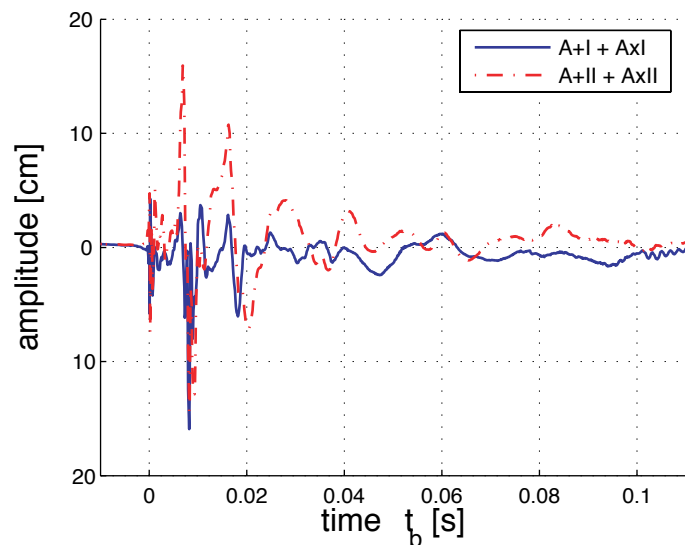


(Liebendörfer, Pen, Thompson, PoS(NIC-IX)132, 2006)

Prediction of Gravitational Wave Signal



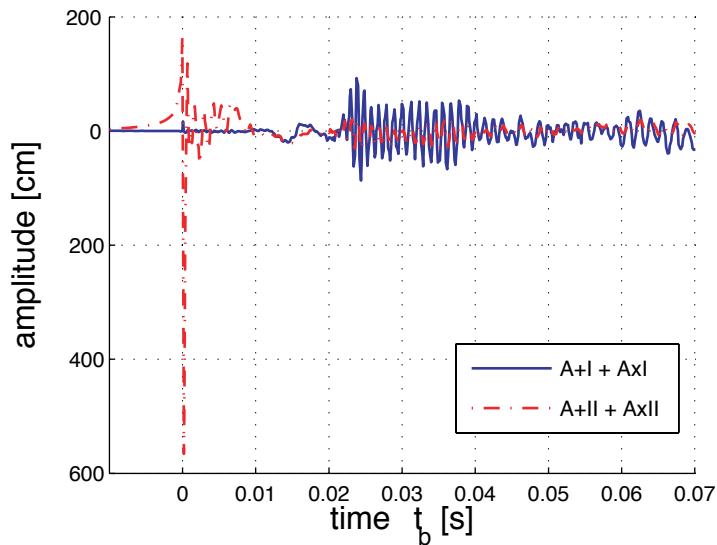
Fast rotating 15Ms
progenitor
 $\Omega \sim 2 \times 10^4$ rad/ps
--> imprint of bounce
and rotation rate
(see Ott et al. 2007)



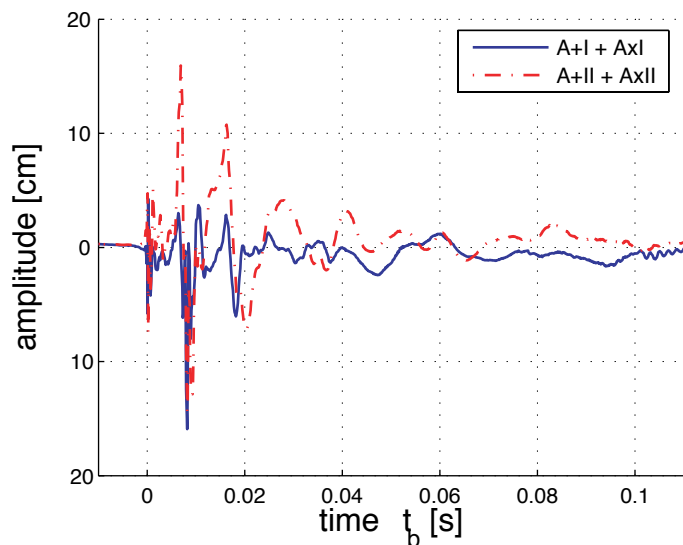
Slowly rotating 15Ms
progenitor according
to
(Heger, Woosley & Spruit 2005)

(Scheidegger et al. 2007/8)

Prediction of Gravitational Wave Signal



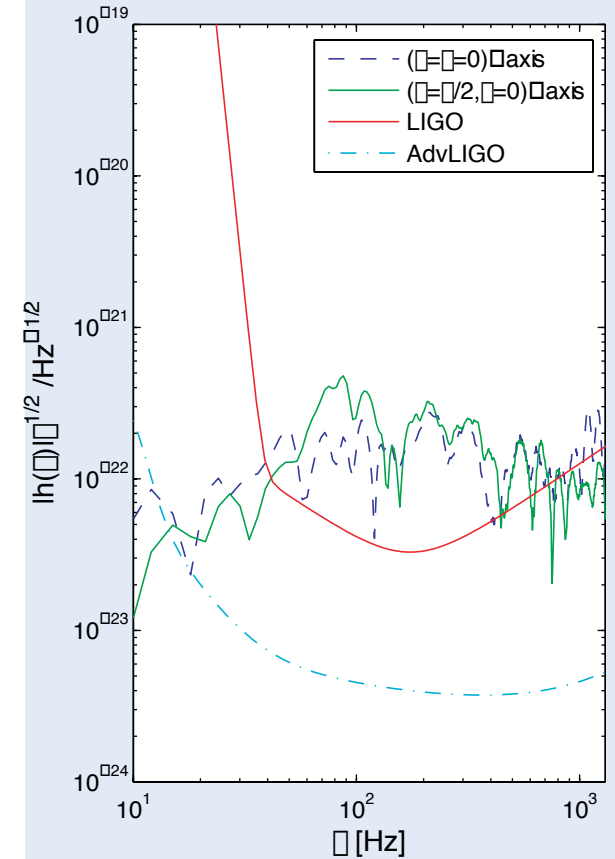
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(Scheidegger et al. 2007/8)

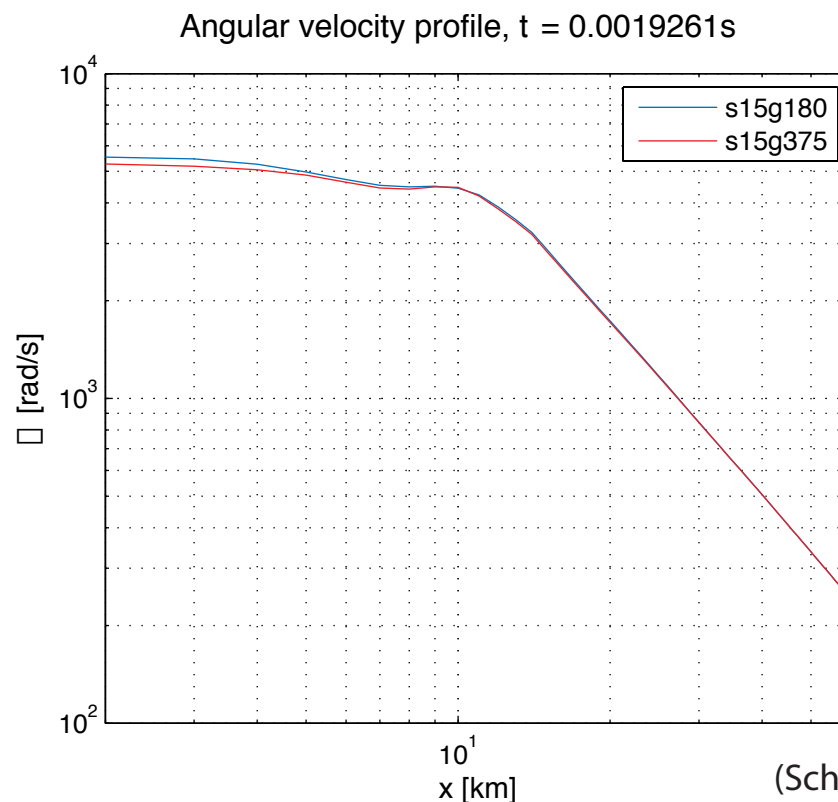


Galactic supernovae
 -- could (LIGO)
 -- should (Adv. LIGO)
 be detectable

Incompressibility & Rotation

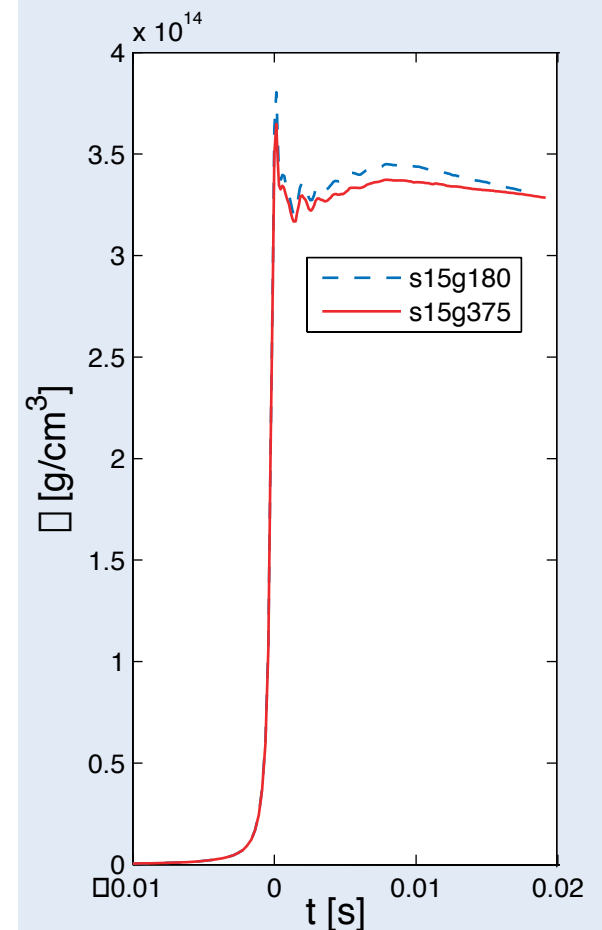
Two runs based on Lattimer-Swesty (1991) EoS
with incompressibility $K=180$ MeV and $K=375$ MeV:

- > stiffer core becomes less dense
- > angular velocity is smaller

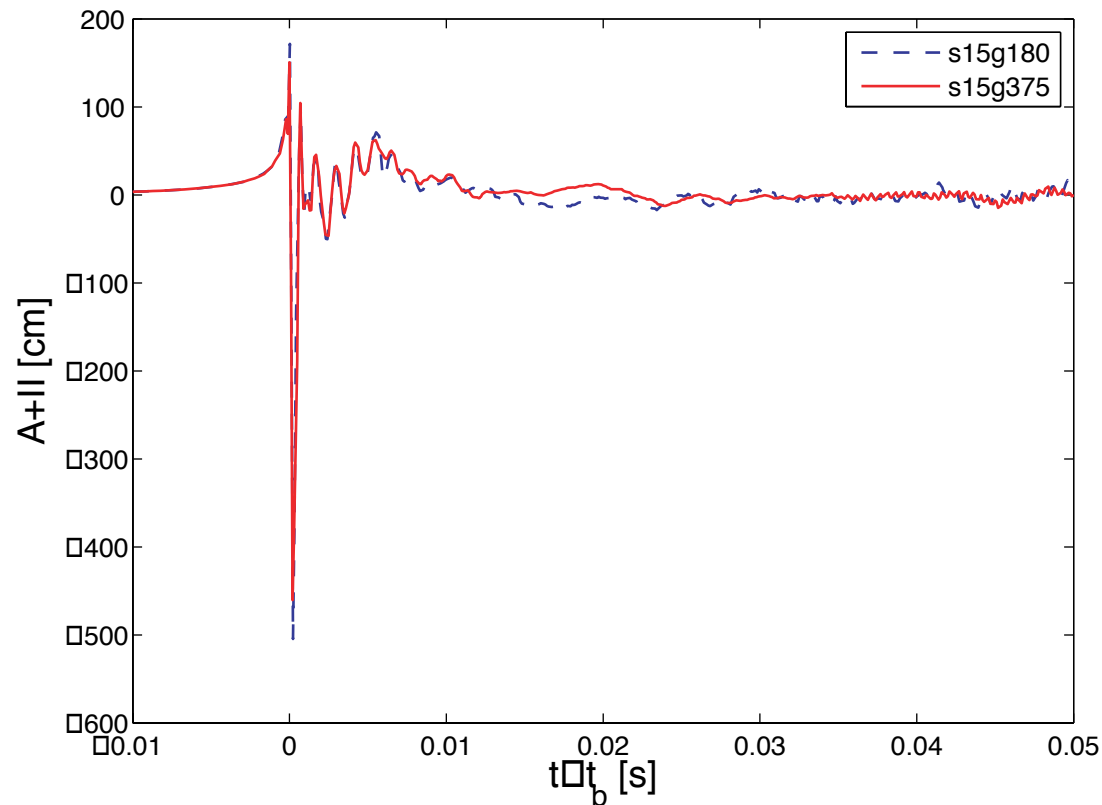


(Scheidegger et al. 200X)

Central density:



Impact of EoS on GW emission



- the direct impact is small!
- Is there an indirect impact on fluid instabilities that produce larger variations in GW emission?

Run1 --> $K=180$ MeV

Run2 --> $K=375$ MeV

Maximum density:

Run1 --> $\rho=3.8E14$ g/cm³

Run2 --> $\rho=3.6E14$ g/cm³

Maximum Amplitude
(A+II at bounce):

Run1 --> $A=506$ cm

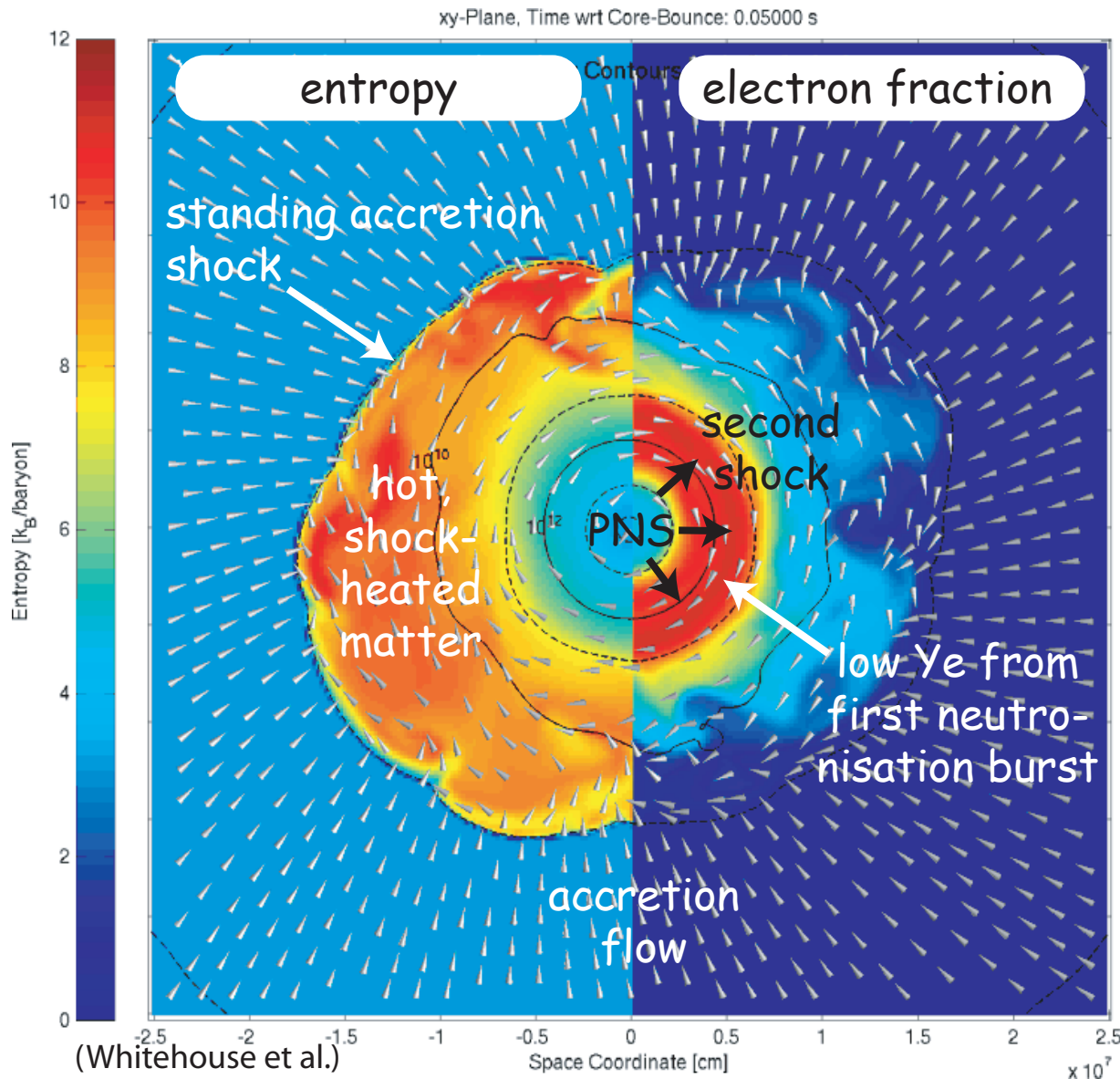
Run2 --> $A=406$ cm

Characteristic
frequency:

Run1 --> $f_c = 657$ Hz

Run2 --> $f_c = 565$ Hz

Conclusions



- Neutrinosignal reflects PNS compressibility and accretion rate, sensitive to
 - > equation of state
 - > PNS thermal profile
 - > weak interaction rates
- Select bag constant for early QCD phase transition to quark matter
 - > second accretion shock
 - > anti-neutrino burst
 - > shift in rms energies
 - > triggers explosion
- 3D models to study EoS compressibility in GW signal
 - > small effect @ bounce
 - > larger @ postbounce fluid instabilities and phase trans.?